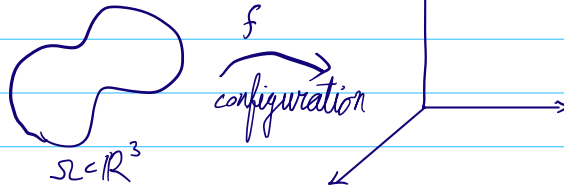


Introduction & motivation

Standard hyperelasticity



Elastic energy: $E(f) = \int_{\Omega} W(Df_x) dx$

- $W: \mathbb{R}^{3 \times 3} \rightarrow [0, \infty)$, $\{W=0\} = SO(3)$ (single well structure)
- $W(RA) = W(A) \quad \forall A \in \mathbb{R}^{3 \times 3}, R \in SO(3)$ (frame indifference)
- $W(A) \geq c \text{dist}^2(A, SO(3))$ (coercivity)
- \vdots

Isotropy group of W $G_w = \{Q \mid W(AQ) = W(A) \quad \forall A\} \subset SO(3)$

Isotropic material $G_w = SO(3)$

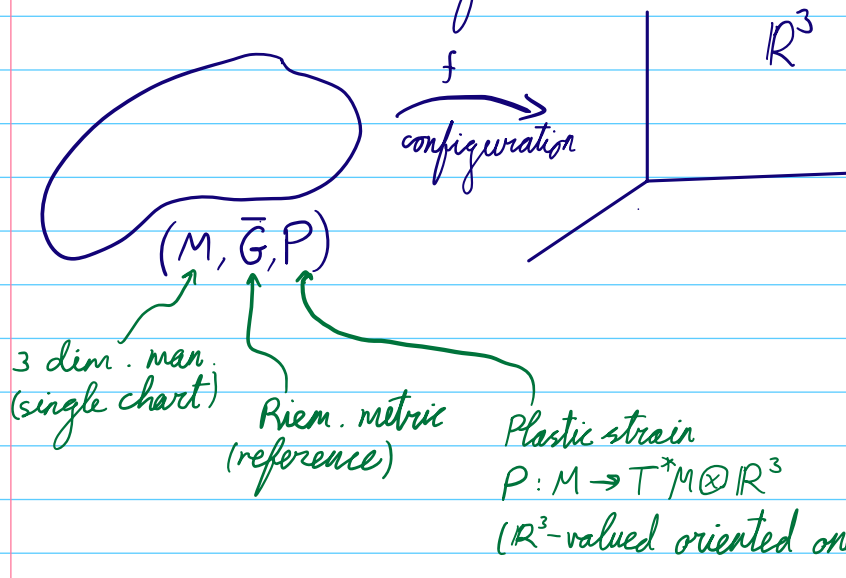
Example: $W(A) = \text{dist}^2(A, SO(3))$

The well assumption $\{W=0\} = SO(3)$ implies the existence of zero energy (reference) configurations: $f(x) = Rx + b$.

Main focus of this course:

Problems in which no reference configurations exist

I Non-Euclidean elasticity



Elastic energy: $E_M(f) = \int_M W(Df_x \circ P_x^{-1}) dV_{\bar{G}}$

$\underset{E_{M,P}(f)}{\parallel}$

↑ "implant map"

Volume form of \bar{G}
 $|\det \bar{G}|^{\frac{1}{2}} dx$

Compatibility of P & \bar{G} : $\bar{G}_x(v, w) = \langle P_x v, P_x w \rangle$, hence $P \rightsquigarrow \bar{G}$.

I.e. $P_x \in SO(\bar{G}_x) := \{A: T_x M \rightarrow \mathbb{R}^3 \mid A \text{ is orientation-preserving isometry}\}$

Example: In a chart, $\sqrt{\bar{G}}$ is a compatible strain.

If P, P' are two plastic strains compatible with \bar{G} , then $P'_x = Q_x P_x$ for some $Q_x \in SO(3)$.

$$\forall_x Q_x \in G_W \implies E_{M,P} = E_{M,P'}$$

In particular, if W is isotropic, then

$$E_M(f) = \int W(Df_x \circ \bar{G}_x^{-\frac{1}{2}}) |\det \bar{G}_x|^{\frac{1}{2}} dx$$

Corollary: $E_M(f) = 0 \iff Df_x \in SO(\bar{G}_x)$ a.e.

Pf: $W(Df_x \circ P_x^{-1}) = 0 \iff Df_x \circ P_x^{-1} \in SO(3)$

$$\iff Df_x \in SO(3) P_x = SO(\bar{G}_x)$$

Laplace Beltrami
 $\frac{1}{|\bar{G}|^{1/2}} \partial_j (\bar{G}^{jk} |\bar{G}|^{1/2} \partial_k)$

Prop: $Df_x \in SO(\bar{G}_x)$ a.e. $\Rightarrow \Delta_{\bar{G}} f = 0$.

In particular, f is smooth.

This implies that $f: M \rightarrow \mathbb{R}^3$ is an isometric imm., hence \bar{G} is flat - $\mathcal{R}^{\bar{G}} = 0$.

Pf: $A \in SO(\bar{G}_x) \Leftrightarrow \det_{\bar{G}} A = 1, \text{ cof}_{\bar{G}} A = A$

$\det_{\bar{G}} A := |\bar{G}|^{-1/2} \det A$
 $\text{cof}_{\bar{G}} A := |\bar{G}|^{-1/2} \text{cof} A \cdot \bar{G}$

As in Euclidean case, $\text{div}_{\bar{G}} (\text{cof}_{\bar{G}} Df)^\# = 0 \quad \forall f$

$A^\# = A \bar{G}^{-1}$
 $\text{div}_{\bar{G}} V = |\bar{G}|^{-1/2} \partial_j (|\bar{G}|^{1/2} V^j)$

Under our assumption $\text{cof}_{\bar{G}} Df = Df$, hence

$\Delta_{\bar{G}} f = \text{div}_{\bar{G}} (Df)^\# = 0.$

(can also be written using a weak for.)

The flatness follows from Riemann's theorem.

Thm (LP'11, KMS'19): $\inf E_M = 0 \Rightarrow \mathcal{R}^{\bar{G}} = 0$. The converse holds if M is simply-connected.

In KMS'19 also for non-Euc. target

These works show a stronger claim: $\inf E_M = 0 \Leftrightarrow \exists f \in C^\infty(M; \mathbb{R}^n)$ iso. imm.

Pf: \Rightarrow We will deduce it from a following estimate.

\Leftarrow flatness + simply-conn. $\leadsto \exists f: M \rightarrow \mathbb{R}^3$ iso. imm. $\Rightarrow E_M(f) = 0$.

So for M simply-conn. we get $\inf E_M > 0 \Leftrightarrow \mathcal{R}^{\bar{G}} \neq 0$. Can we get more quantitative?
 Let's "localize" the question:

Thm (MS'19): Let $p \in M$, then $\inf_{KM'21} E_{B_r(p)} = r^4 |\mathcal{R}^{\bar{G}}(p)|^2 + o(r^4)$.

some norm, incl. of p for isotropic W .

Pf: For simplicity, choose $W(A) = \text{dist}^2(A, SO(n))$.

We can choose normal coordinates on $B_r(p)$ (for small enough $r > 0$), and then

$B_r(p) \xrightarrow{z} B_r(0) \subset \mathbb{R}^n, \quad \bar{G}(z) = \delta_{ij} + \frac{1}{6} \mathcal{R}_{kijl}(p) x^k x^l + O(|x|^3)$

hence $\bar{G}^{-1/2}(z)_i^j = \delta_i^j - \frac{1}{6} \mathcal{R}_{kijl}(p) x^k x^l + O(|x|^3)$

and thus, for the map $\iota(x) = x$ we have

our energy scaling

$E_{B_r(p)}(\iota) = \int_{B_r(0)} \text{dist}^2(I - \frac{1}{6} \mathcal{R}_{kijl}(p) x^k x^l + O(|x|^3)) \cdot (1 + O(|x|^2)) dx \sim r^4$

Lower bound: Assume $E_{B_r(0)}(f_r) \leq Cr^4$. Then, since $|\bar{G}^{-\frac{1}{2}} - I| \leq C|x|^2$,

$$\begin{aligned} \text{dist}(f_r, SO(n)) &\leq \text{dist}(f_r \circ \bar{G}^{-\frac{1}{2}}, SO(n)) + C|f_r||x|^2 \\ &\leq C(\text{dist}(f_r \circ \bar{G}^{-\frac{1}{2}}, SO(n)) + |x|^2) \end{aligned}$$

hence $\int_{B_r(0)} \text{dist}^2(f_r, SO(n)) dx = O(r^4)$.

From FJM we have $R \in SO(n)$ s.t. $\int_{B_r(0)} |Df_r - R|^2 \leq Cr^4$.

$\exists c_r$ s.t. $\bar{f}_r = R^T f_r - c_r$ sat. $\int |\bar{f}_r - x|^2 + |D\bar{f}_r - I|^2 \leq Cr^4$

Define now $y = \frac{x}{r} \in B_1(0)$, $v_r(y) = \bar{f}_r(ry) - ry$, then

$$\frac{1}{r^3} Dv_r \xrightarrow{L^2} Df \text{ in } L^2(B_1(0), \mathbb{R}^n)$$

$$D\bar{f}_r(ry) \circ \bar{G}^{-\frac{1}{2}}(ry) = I + \underbrace{\frac{1}{r^3} Dv_r(y)}_{\xrightarrow{L^2} Df} + \underbrace{\frac{\bar{G}^{-\frac{1}{2}}(ry) - I}{r^2}}_{\xrightarrow{L^\infty} -T(y)} + \underbrace{\frac{1}{r^3} Dv_r(y)(\bar{G}^{-\frac{1}{2}}(ry) - I)}_{\xrightarrow{L^2} 0}$$

Hence:

$$\liminf \frac{1}{r^4} E_{B_r(0)}(f_r) = \liminf \frac{1}{r^4} E_{B_r(0)}(\bar{f}_r)$$

$$= \liminf \frac{1}{r^4} \int_{B_r(0)} \text{dist}^2(Df_r(x) \circ \bar{G}^{-\frac{1}{2}}(x)) (1 + O(|x|^2)) dx$$

$$\begin{aligned} \text{dist}(I+A, SO(n)) \\ = |\text{sym} A| + O(|A|^2) \end{aligned}$$

$$\geq \int_{B_1(0)} |\text{sym} Df - T|^2 dy \geq \min_{S \in W^{1,2}(B_1(0))} \int |\text{sym} Df - T|^2 dy =: I(\mathcal{R}(p))$$

Now, $\mathcal{R} \mapsto (I(\mathcal{R}))^{\frac{1}{2}}$ is a norm:

homogeneity + triangle ineq. are easy.

Positivity follows from Saint-Venant compatibility condition:

$$T = \text{sym} Df \Leftrightarrow \underbrace{\partial_{ij} T_{kl} + \partial_{kl} T_{ij} - \partial_{ik} T_{jl} - \partial_{jl} T_{ik}}_{R_{ijkl}(p)} = 0$$

$R_{ijkl}(p)$

Upper bound: Let f be a min. of $\int_{B_r^0} |\text{sym} Df - T|^2 dy$

and define $f_r(x) = x + r^3 f(x/r)$.

The above calculation yields $\frac{1}{r^4} E_{B_r^0}(f_r) \rightarrow I(\mathcal{R})$ □

Comment: If we take another isotropic W , we get $\int Q(\text{sym} Df - T)$,

where $Q(A) = \frac{1}{2} D_A^2 W(A, A)$ is positive def. on sym. matrices.

If W is not isotropic, then, writing $P_x = \underbrace{Q_x}_{\text{SO}(n)} \bar{G}^{-\frac{1}{2}}$, we get

$$\int_{B_r^0} Q(Q_p^T (\text{sym} Df - T) Q_p) dy$$