

## Exercises for “Quantum Phase Transitions”

Summer 24

DR. L. JANSSEN

## Exercise 2 (03.05.24)

## 1. Tricritical point in an antiferromagnet

(6 points)

An external magnetic field  $h$  applied to an antiferromagnet couples to the total magnetization  $m$  instead of the antiferromagnetic order parameter, the staggered magnetization  $n$ . Assume that the coupling between  $m$  and  $n$  is described phenomenologically by the Landau free energy density

$$f(n, m) = \frac{t}{2}n^2 + \frac{b}{4}n^4 + \frac{v}{2}m^2 + \frac{w}{2}n^2m^2 - hm, \quad (1)$$

where  $t = (T - T_0)a$ , and  $a, b, v, w$  are positive constants.

- (a) Show that this model features a paramagnetic phase with magnetization  $m_0$ . Derive a temperature-independent relation  $m = m(n^2)$  between the magnetization  $m$  and the staggered magnetization  $n$  in the antiferromagnetic phase, i.e., for  $n^2 > 0$ .

*Hint:* Equilibrium states are given by minima of  $f(n, m)$ .

- (b) Consider the antiferromagnetic phase near the phase transition, i.e., for small values of  $n^2$ . Write  $m = m_0 + \delta m$ , expand  $m(n^2)$  for small  $n^2$ , and derive a relation between  $\delta m$  and  $n^2$ .
- (c) Show that the effective free energy density for the staggered magnetization  $g(n) = f(n, m_0 + \delta m) - f(0, m_0)$  can be written as

$$g(n) = \frac{\bar{a}}{2}n^2 + \frac{\bar{b}}{4}n^4 + \frac{\bar{c}}{6}n^6 + \mathcal{O}(n^8), \quad (2)$$

with to-be-determined temperature- and field-dependent coefficients  $\bar{a}$ ,  $\bar{b}$ , and  $\bar{c}$ .

- (d) Show that the model (1) features a tricritical point at temperature  $T_t$  and field  $h_t$ , where

$$T_t = T_0 - \frac{bv}{2aw}, \quad h_t^2 = \frac{bv^3}{2w^2}. \quad (3)$$

*Hint:* Here we define a tricritical point as a point where first- and second-order transition lines meet. Depending on the sign of the coefficient  $\bar{b}$  in  $g(n)$ , the transition between the paramagnetic and antiferromagnetic phases are first or second order, cf. Problem 1 on Exercise 1.

- (e) Show that the second-order phase transition occurs for  $h < h_t$  at

$$T_c = T_0 - \frac{wh^2}{av^2}, \quad (4)$$

and the first-order transition occurs for  $h > h_t$  at

$$T_c = T_0 - \frac{3wh^2}{4av^2} - \frac{bv}{4aw} + \frac{b^2v^4}{16aw^3h^2}. \quad (5)$$

- (f) Sketch the phase diagram in the  $(T, H)$  plane.

*please turn over!*

## 2. Static scaling hypothesis

(4 points)

Consider the static scaling hypothesis for the free energy density

$$f_s(t, h) = b^{-d} f_s(b^{y_t} t, b^{y_h} h) \quad (6)$$

with scaling exponents  $y_t$  for the reduced temperature  $t$  and  $y_h$  for the external field  $h$ .

(a) Use the static scaling hypothesis to derive the relation

$$\delta = \frac{d + 2 - \eta}{d - 2 + \eta} \quad (7)$$

between the critical-isotherm exponent  $\delta$  and the anomalous dimension  $\eta$ .

*Hint:* Use the relation  $y_t = 1/\nu$  and Fisher's law  $\gamma = \nu(2 - \eta)$  derived in class.

(b) In principle, critical exponents above and below a transition could differ from each other. Show for the example of the correlation-length exponents  $\nu$  and  $\nu'$  above and below the transition, respectively, that the static scaling hypothesis implies that they are equal,  $\nu(T > T_c) = \nu'(T < T_c)$ .

*Hint:* For fixed  $h \neq 0$ ,  $f_s(t, h)$  should be a smooth function of  $t$ , because the only singularity that we expect is at  $t = h = 0$ . Show that  $f_s(t, h)$  can be written in the form

$$f_s(t, h) = h^{d/y_h} F_f^\pm \left( \frac{|t|}{h^{1/(\bar{\nu} y_h)}} \right), \quad (8)$$

and explain how the smoothness assumption mentioned above constrains the analytic form of the function  $F_f^\pm(x)$ .