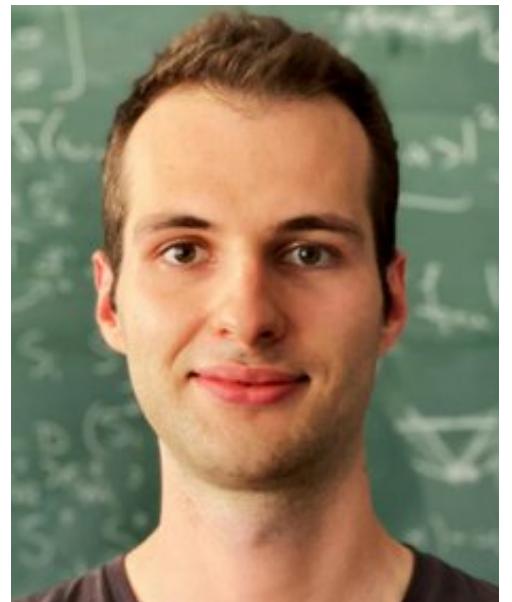


# Fractionalized fermionic quantum criticality in spin-orbital Mott insulators

Lukas Janssen  
(TU Dresden)



Urban Seifert



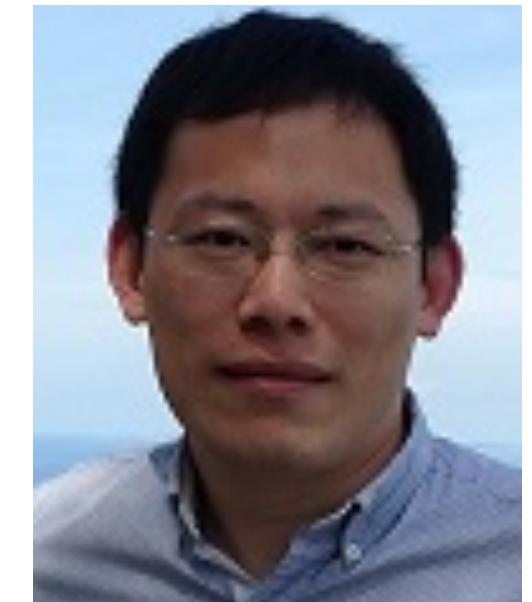
Xiao-Yu Dong



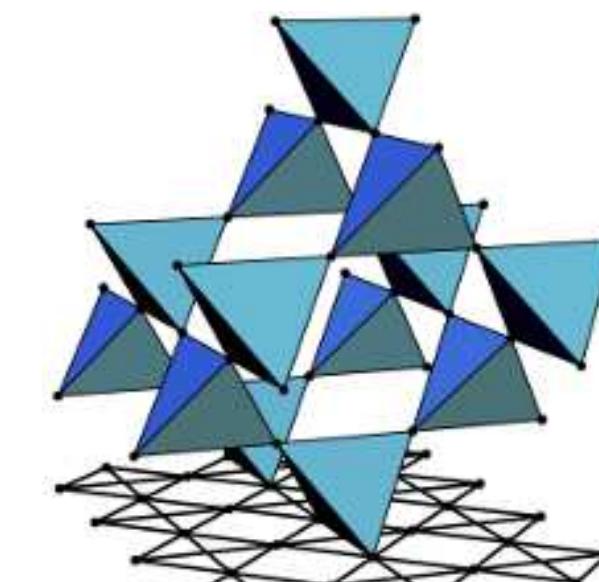
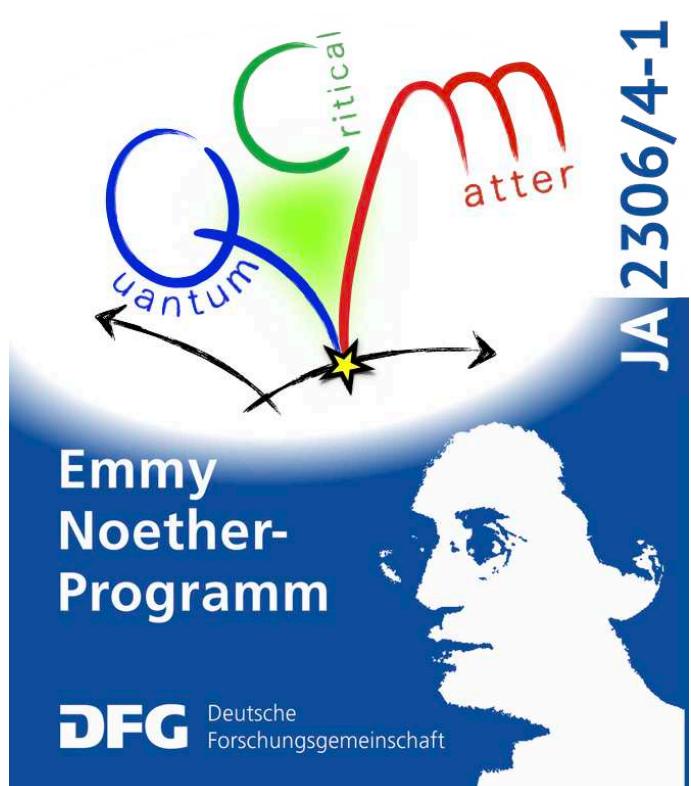
Sreejith Chulliparambil



Matthias Vojta



Hong-Hao Tu

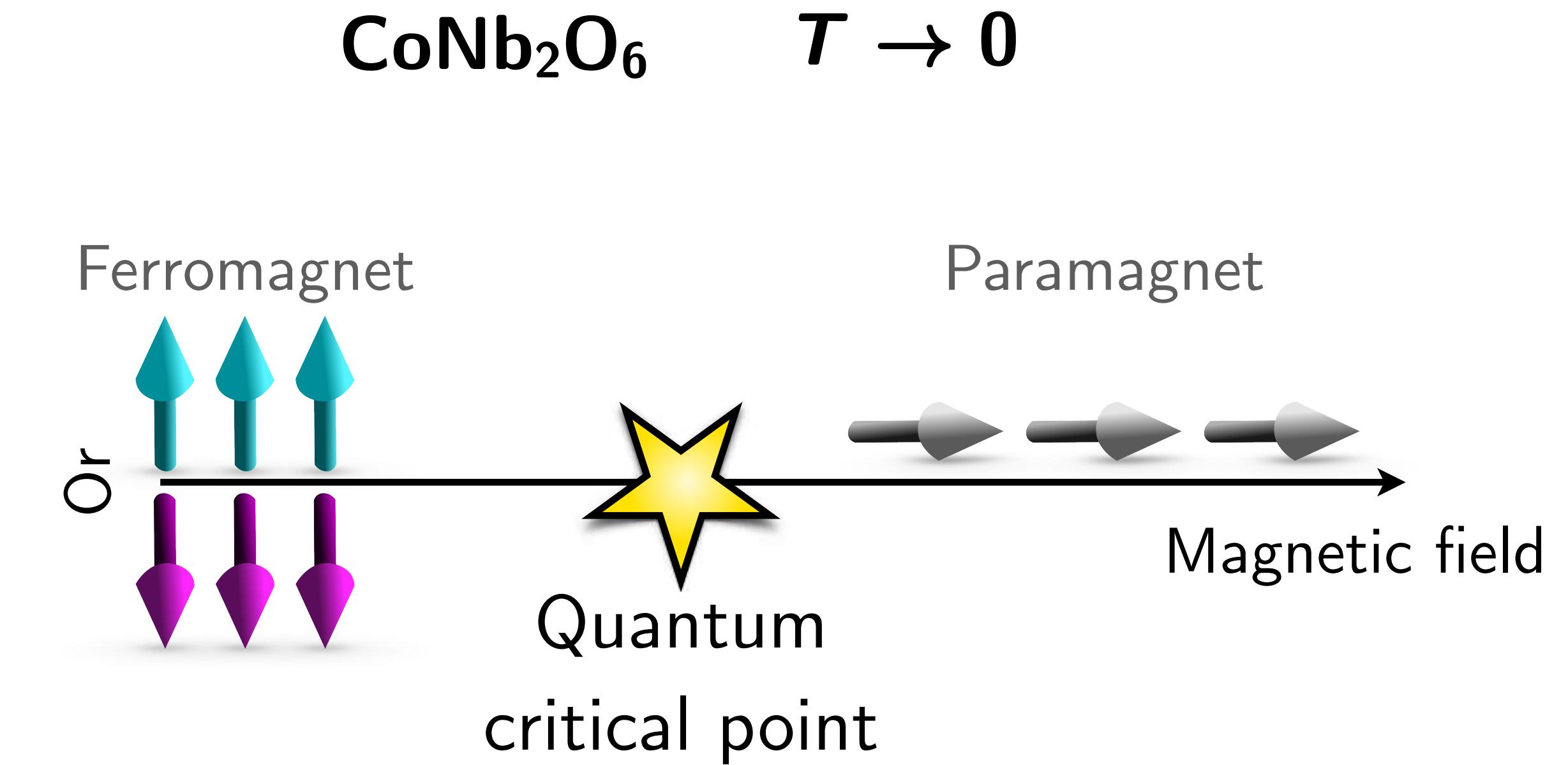
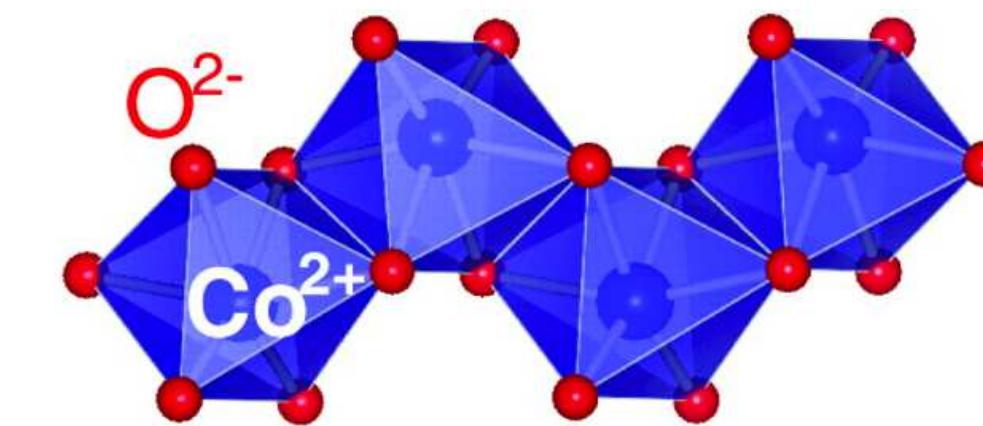
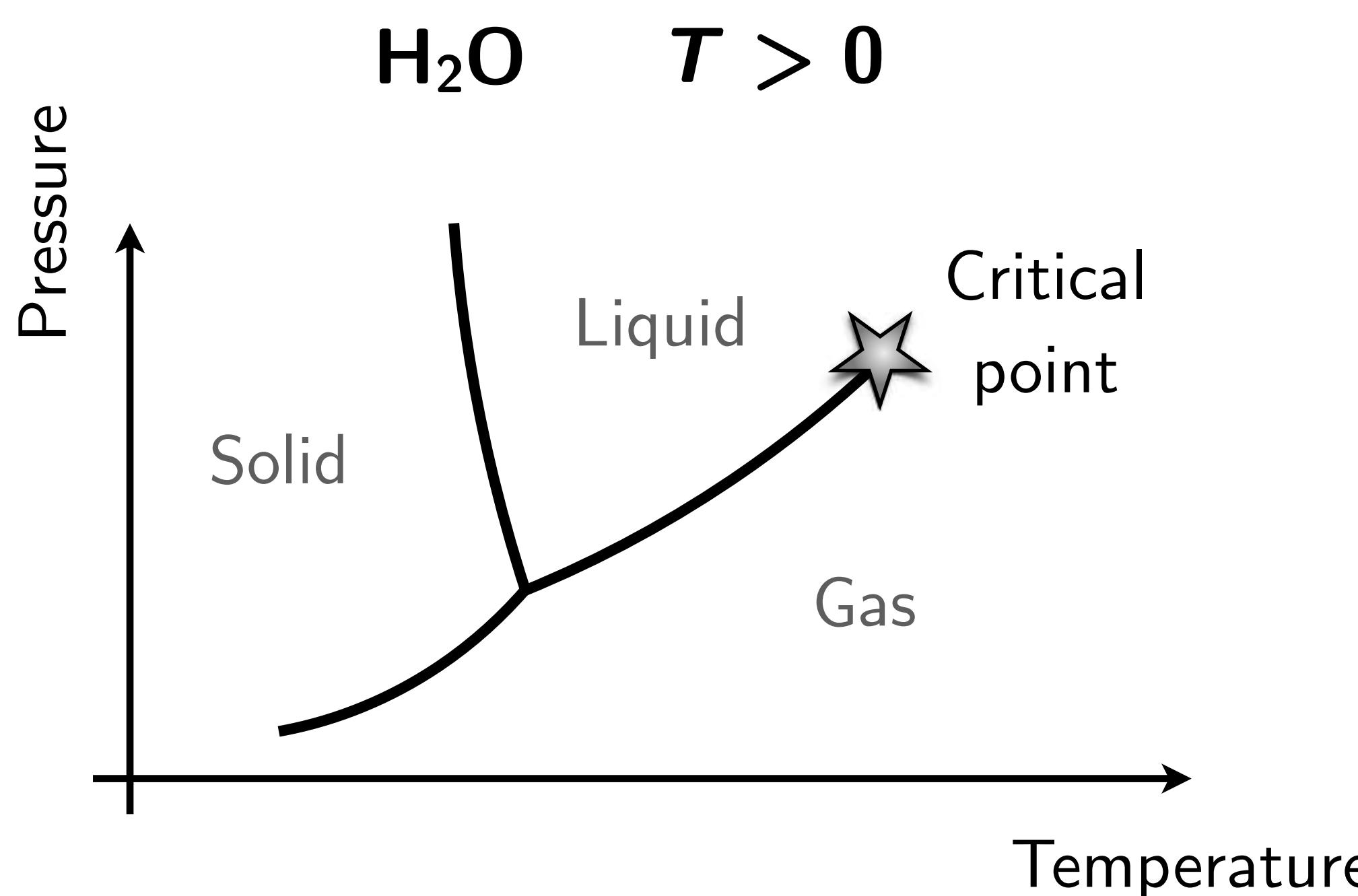


SFB 1143

# Outline

- (1) Introduction: “*Beyond-Landau*” quantum criticality
- (2) Spin-1/2: *Field-induced criticality in Kitaev materials*
- (3) Spin-3/2: *Fractionalized fermionic criticality in spin-orbital models*
- (4) Conclusions

# Classical vs quantum criticality

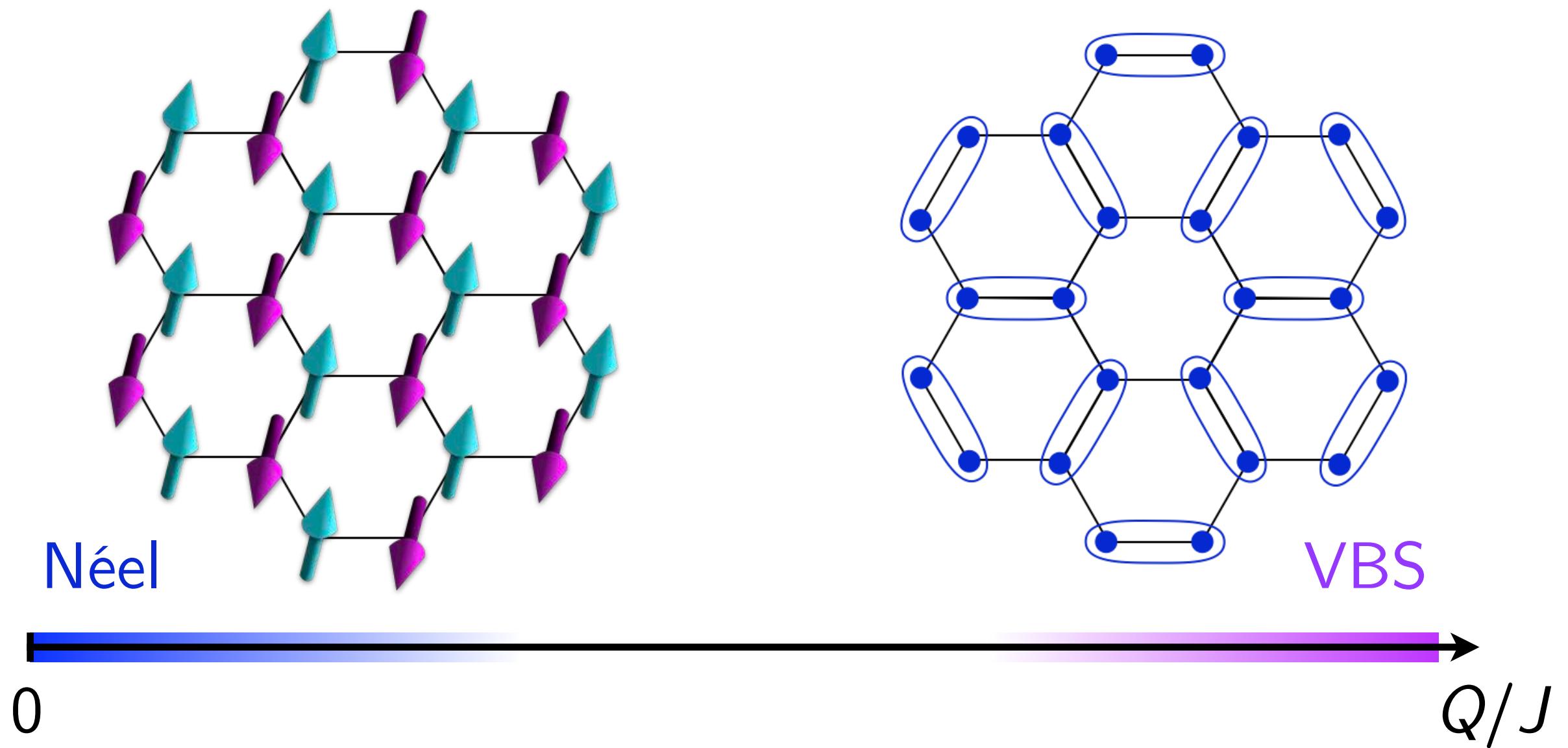


[Coldea et al., Science '10]  
[Kinross et al., PRX '14]

...

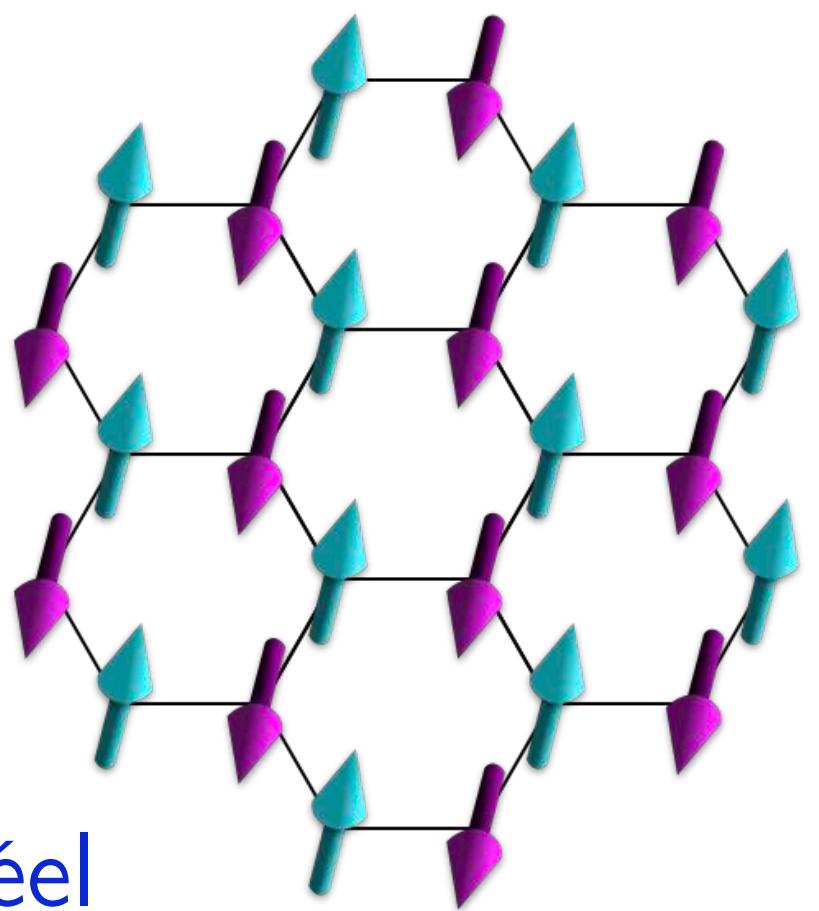
# Deconfined quantum criticality

$$\text{O} = (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) / \sqrt{2}$$

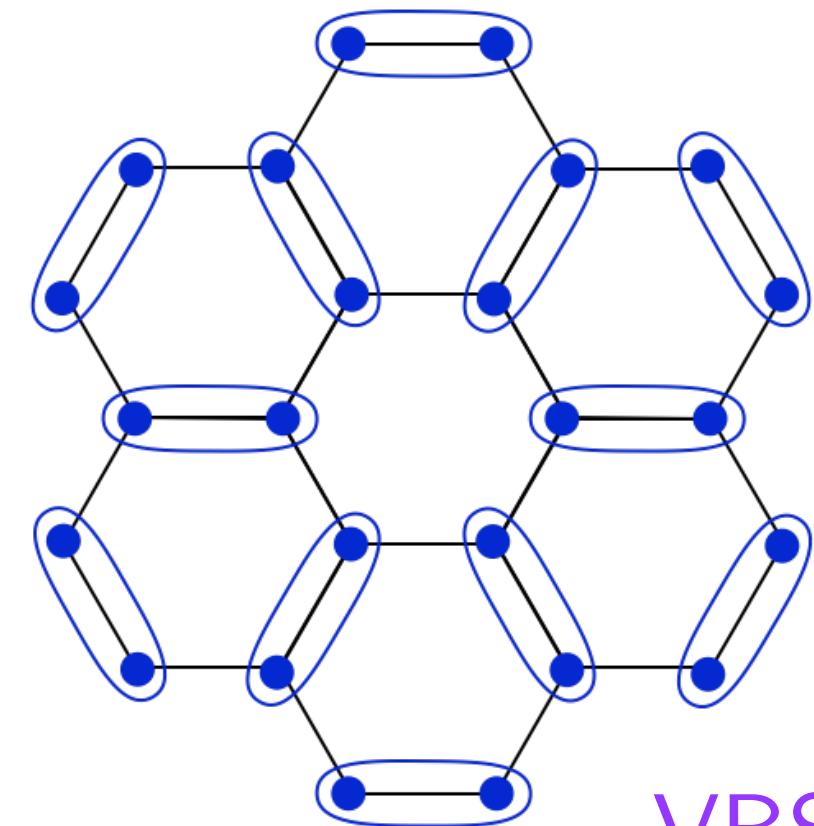


# Deconfined quantum criticality

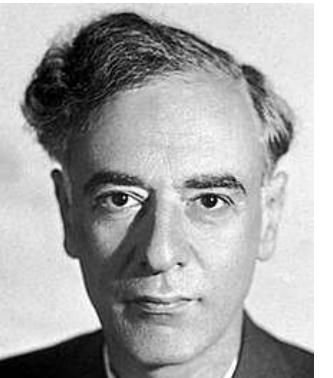
$$\text{Or} = (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) / \sqrt{2}$$



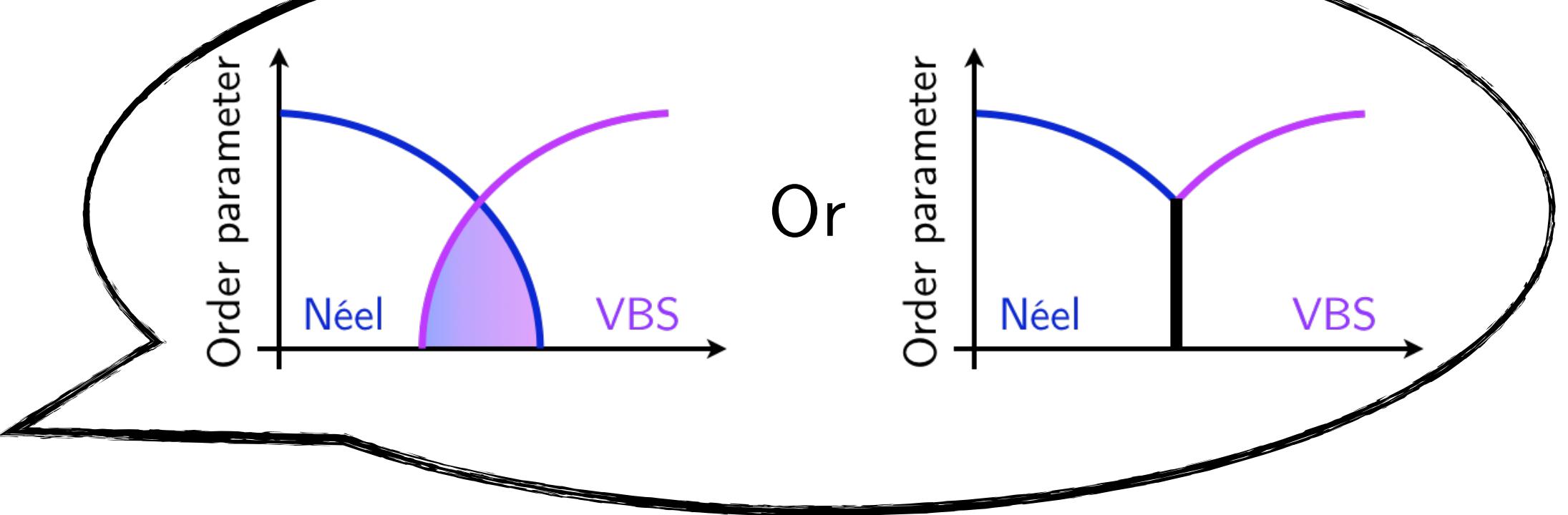
Néel



VBS

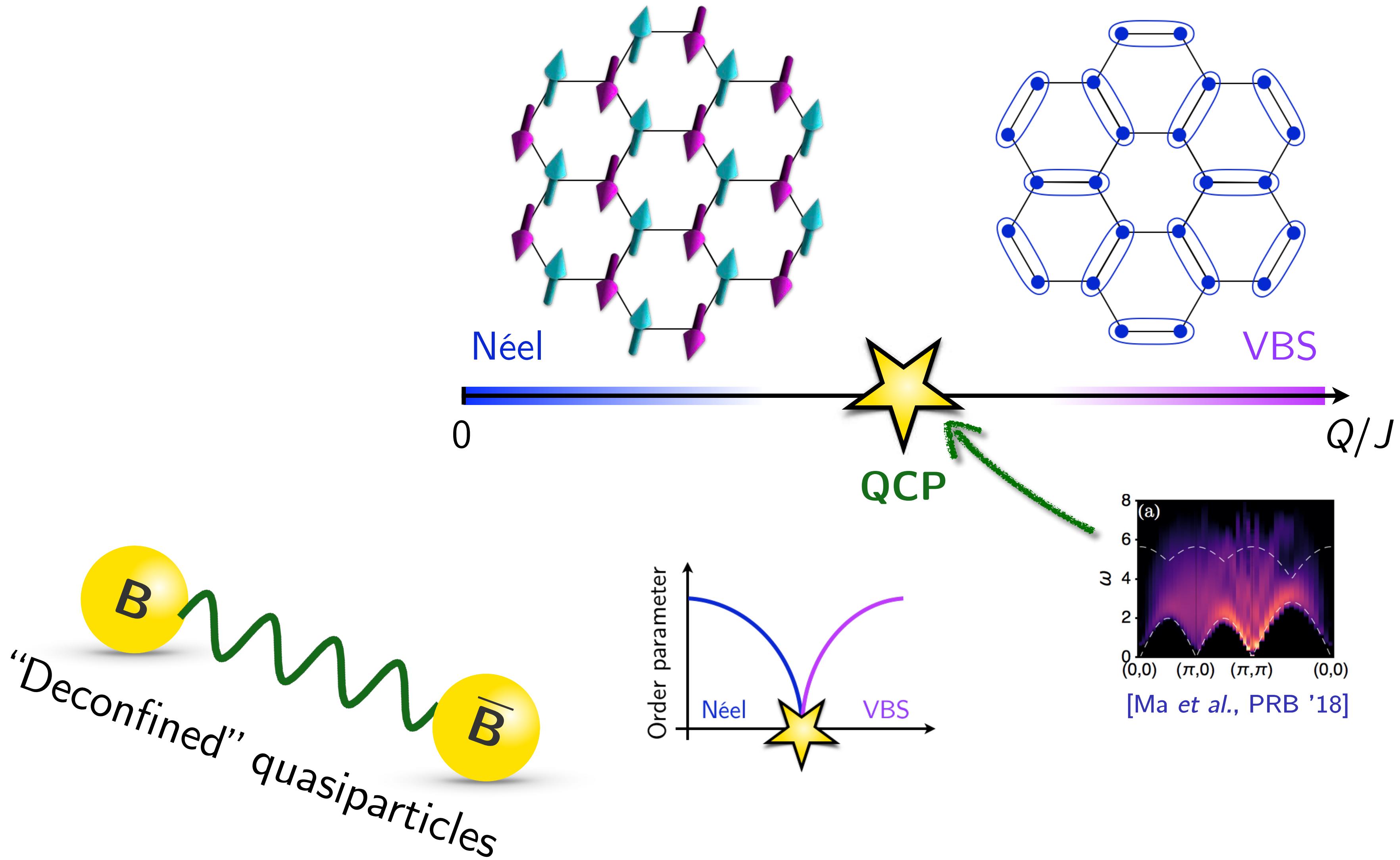


Landau



# Deconfined quantum criticality

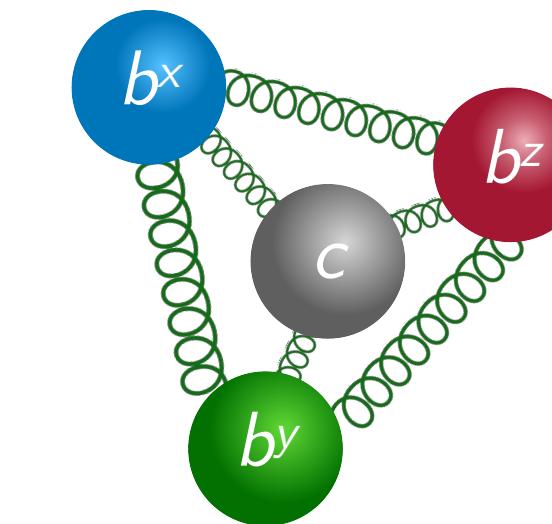
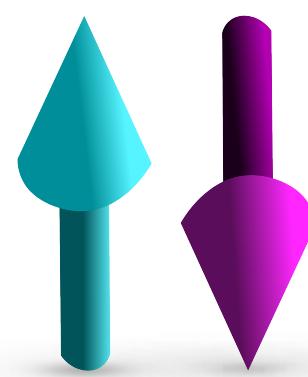
$$\text{O} = (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) / \sqrt{2}$$



- [Senthil et al., Science '04]
- [Pujari, Damle, Alet, PRL '13]
- [Block, Melko, Kaul, PRL '13]
- [Shao, Guo, Sandvik, Science '16]

...

# Spin-liquid transitions



Ordered state

0

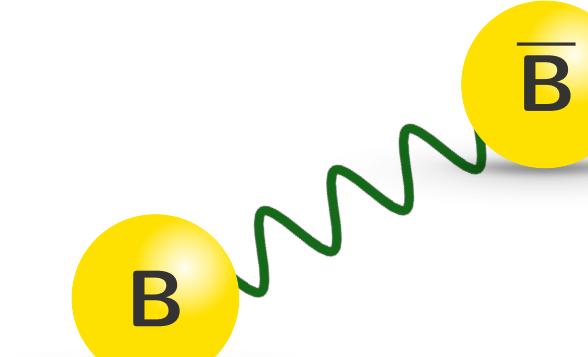
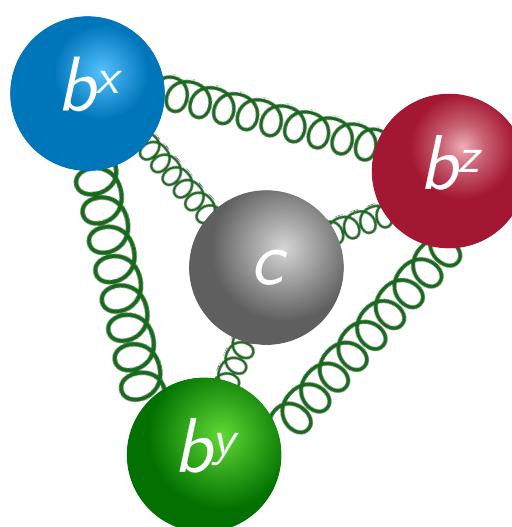
Quantum paramagnet

$J$

[Assaad & Grover, PRX '16]

[LJ, Wang, Scherer, Meng, Xu, PRB '20]

...



Quantum paramagnet #1

0

Quantum paramagnet #2

$J$

[Metlitski, Mross, Sachdev, Senthil, PRB '15]

[LJ & He, PRB '17]

[Boyack, Lin, Zerf, Rayyan, Maciejko, PRB '18]

...

# Outline

- (1) Introduction: “*Beyond-Landau*” quantum criticality
- (2) Spin-1/2: *Field-induced criticality in Kitaev materials*
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- (4) Conclusions

# Frustrated magnets

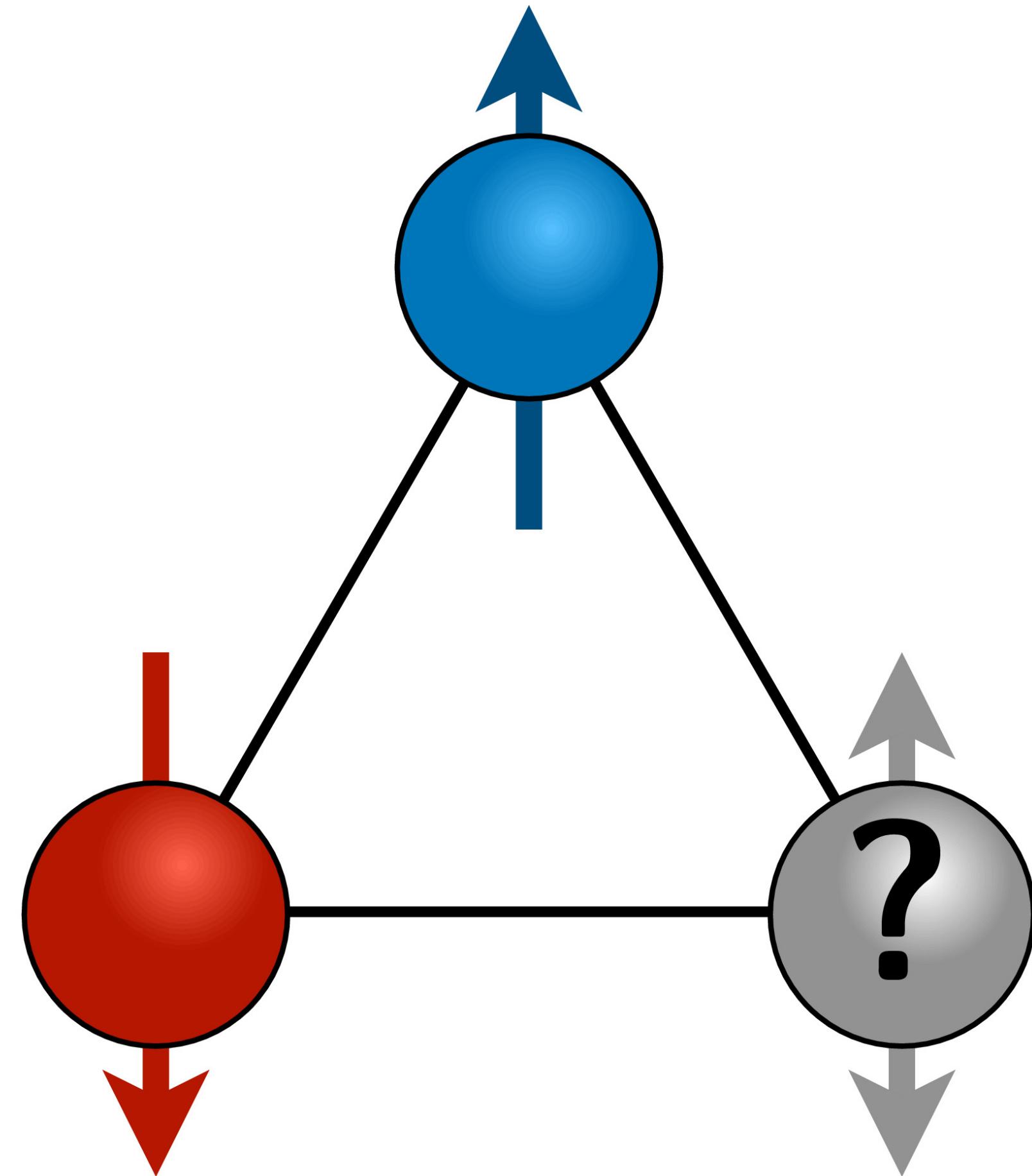
Frustration:

Not all local constraints can be simultaneously **satisfied**

Consequences:

Classical: Exponentially large ground-state manifold

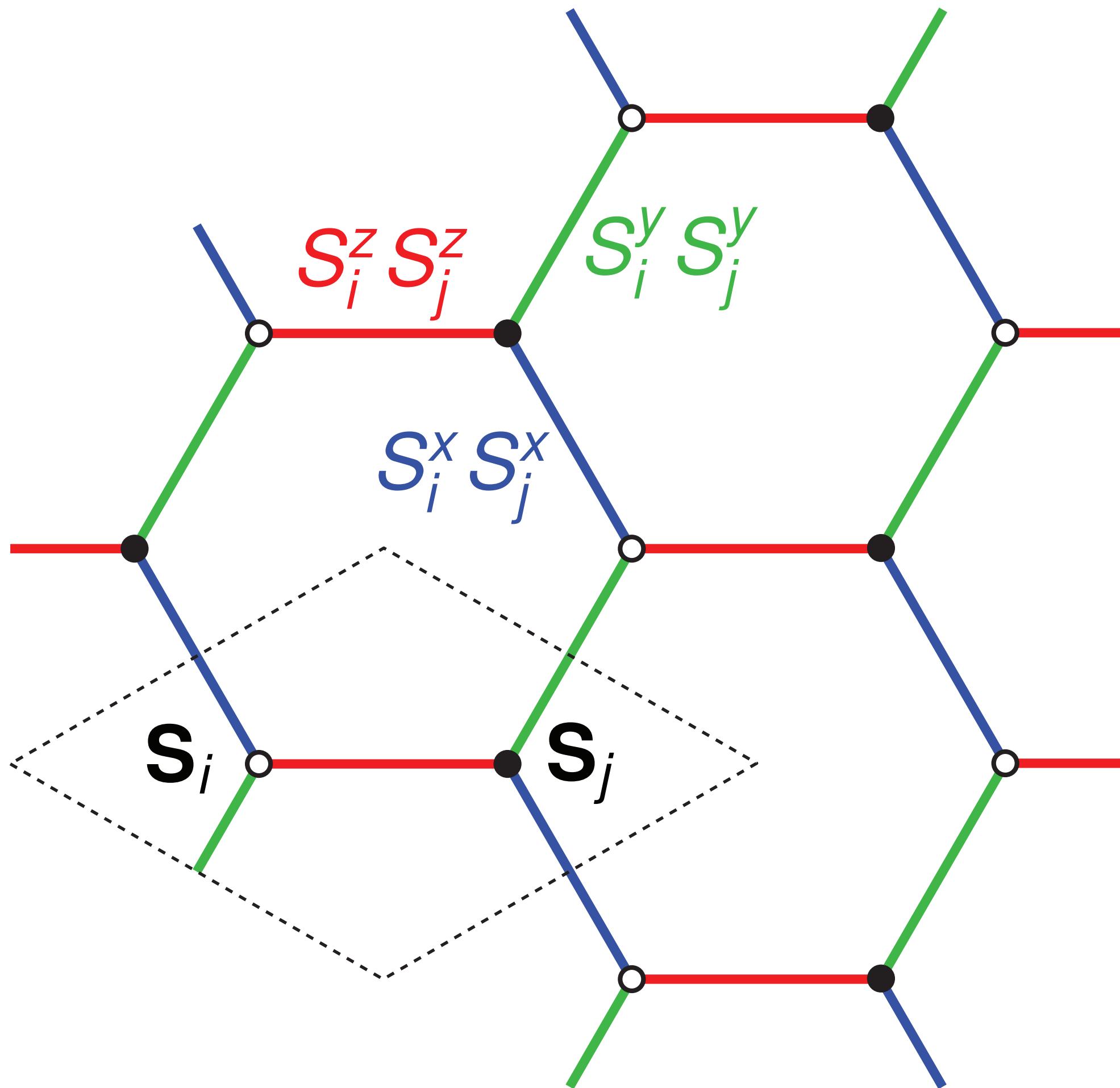
Quantum: New phases of matter?



Antiferromagnetic coupling of 3 Ising spins

# Kitaev honeycomb model

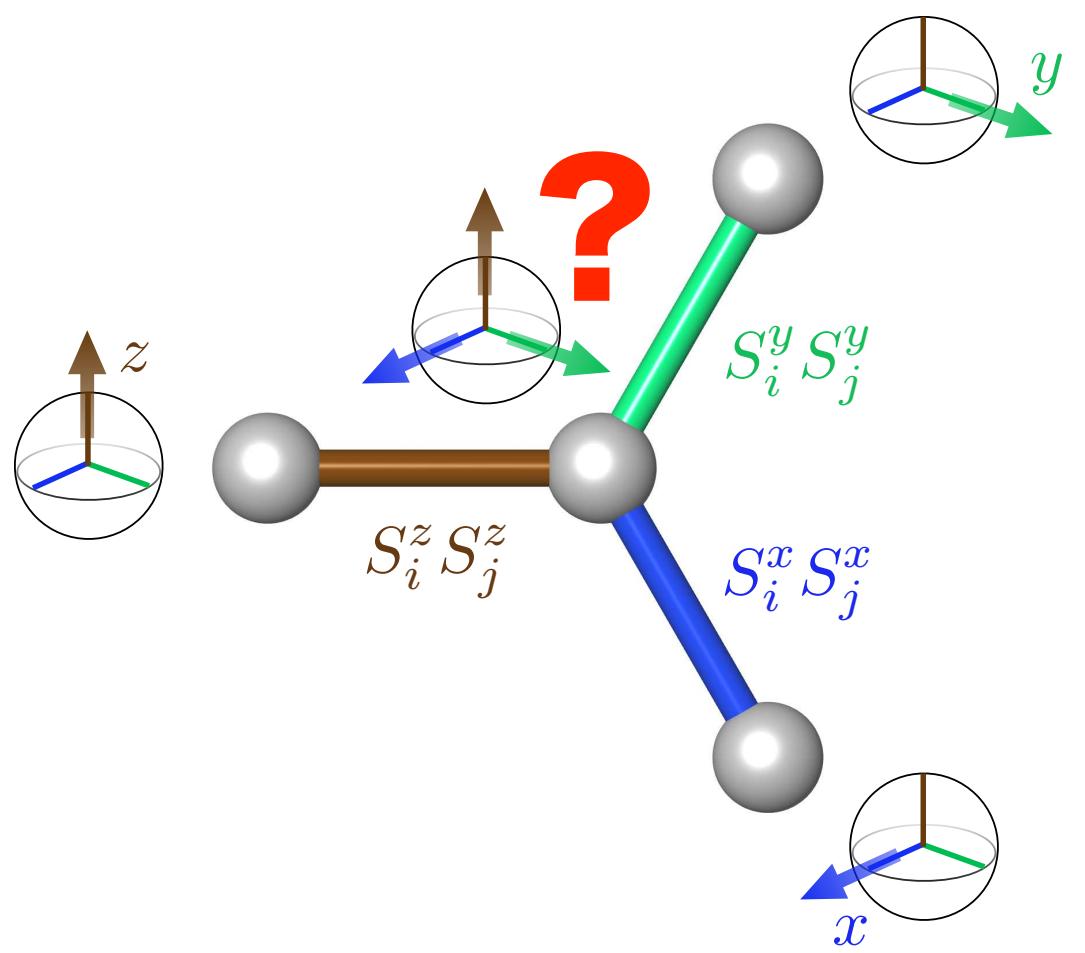
Spin-1/2 on honeycomb lattice:



Hamiltonian:

$$H = -K_x \sum_{\text{blue links}} \sigma_i^x \sigma_j^x - K_y \sum_{\text{green links}} \sigma_i^y \sigma_j^y - K_z \sum_{\text{red links}} \sigma_i^z \sigma_j^z$$

[Kitaev, Ann. Phys. '06]



Exchange frustration

Review: [Trebst, arXiv:1701.07056]



Alexei Kitaev

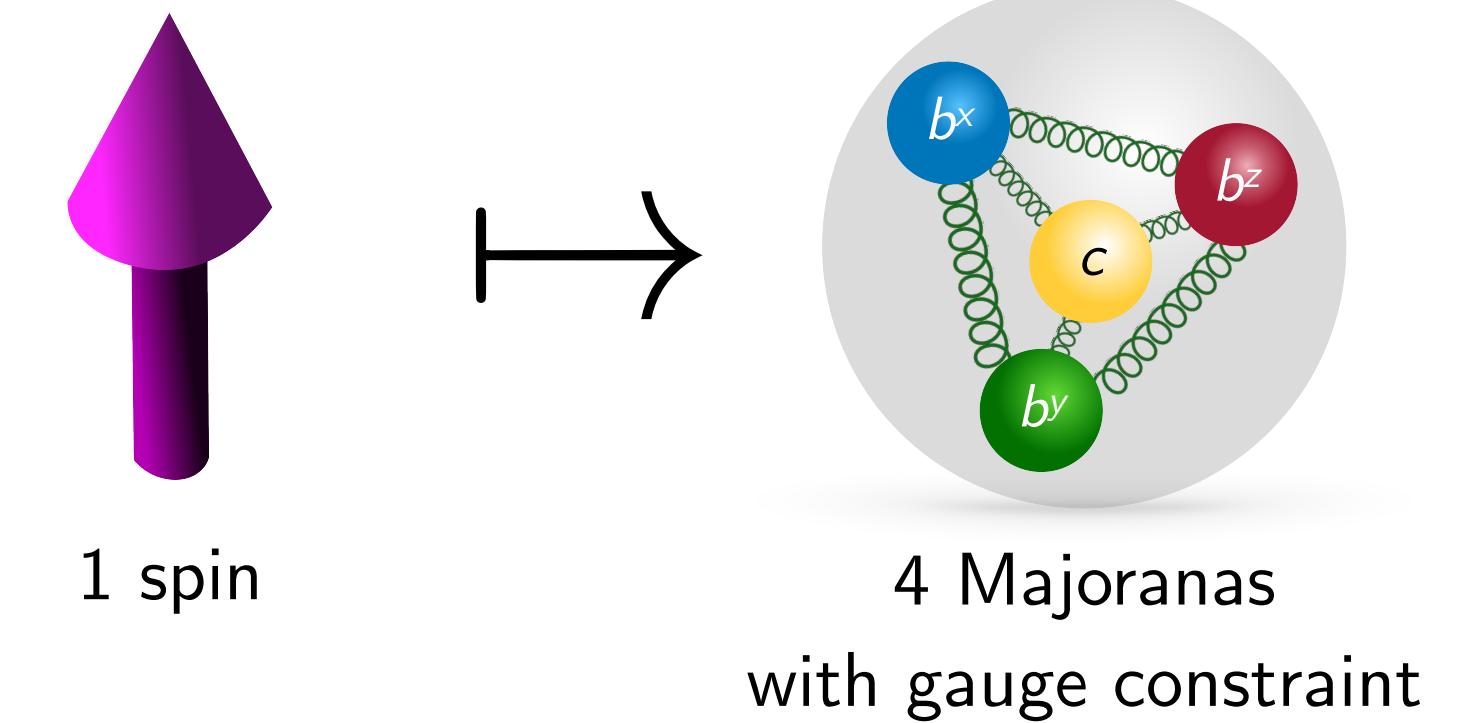
# Parton construction

Majorana representation:

$$\sigma^x \mapsto \tilde{\sigma}^x = i b^x c$$

$$\sigma^y \mapsto \tilde{\sigma}^y = i b^y c$$

$$\sigma^z \mapsto \tilde{\sigma}^z = i b^z c$$



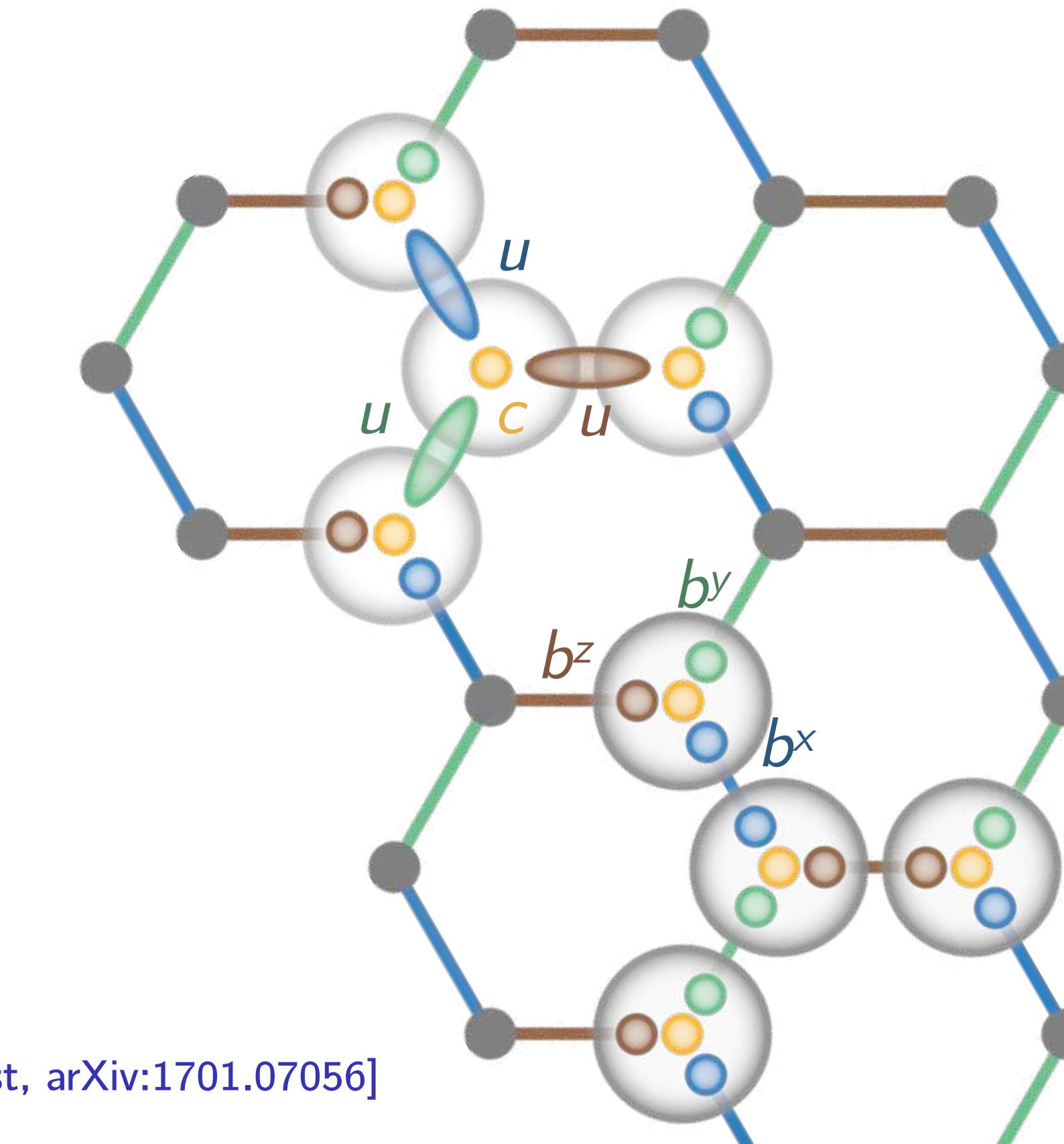
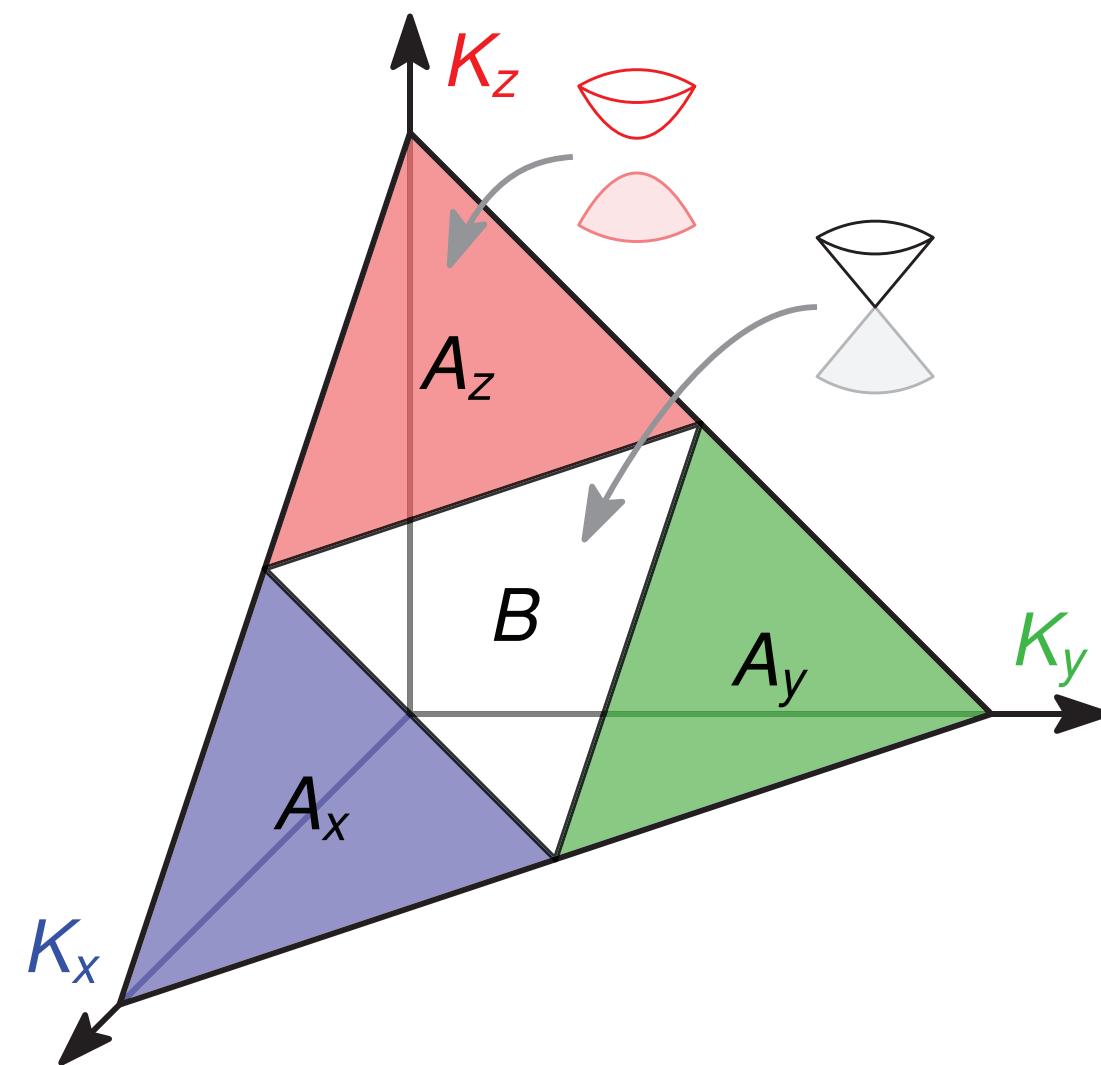
Fractionalization:

$$H \mapsto \tilde{H} = -i \sum_{\langle ij \rangle_\gamma} K_\gamma \underbrace{(ib_i^\gamma b_j^\gamma)}_{\equiv \hat{u}_{ij} = \hat{u}_{ij}^\dagger} c_i c_j$$

static!

Ground-state flux pattern:  $u = 1$   
[Lieb, PRL '94]

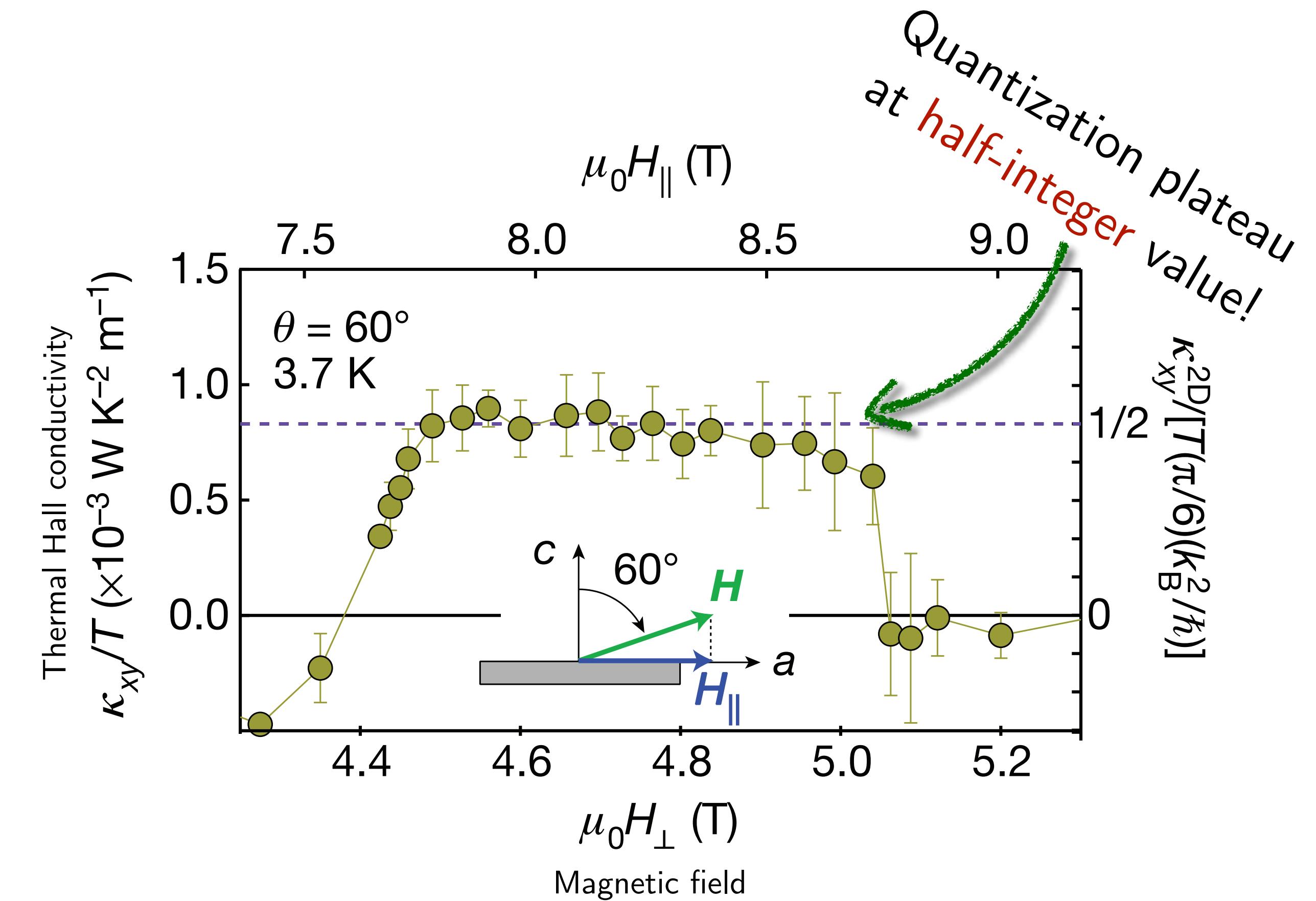
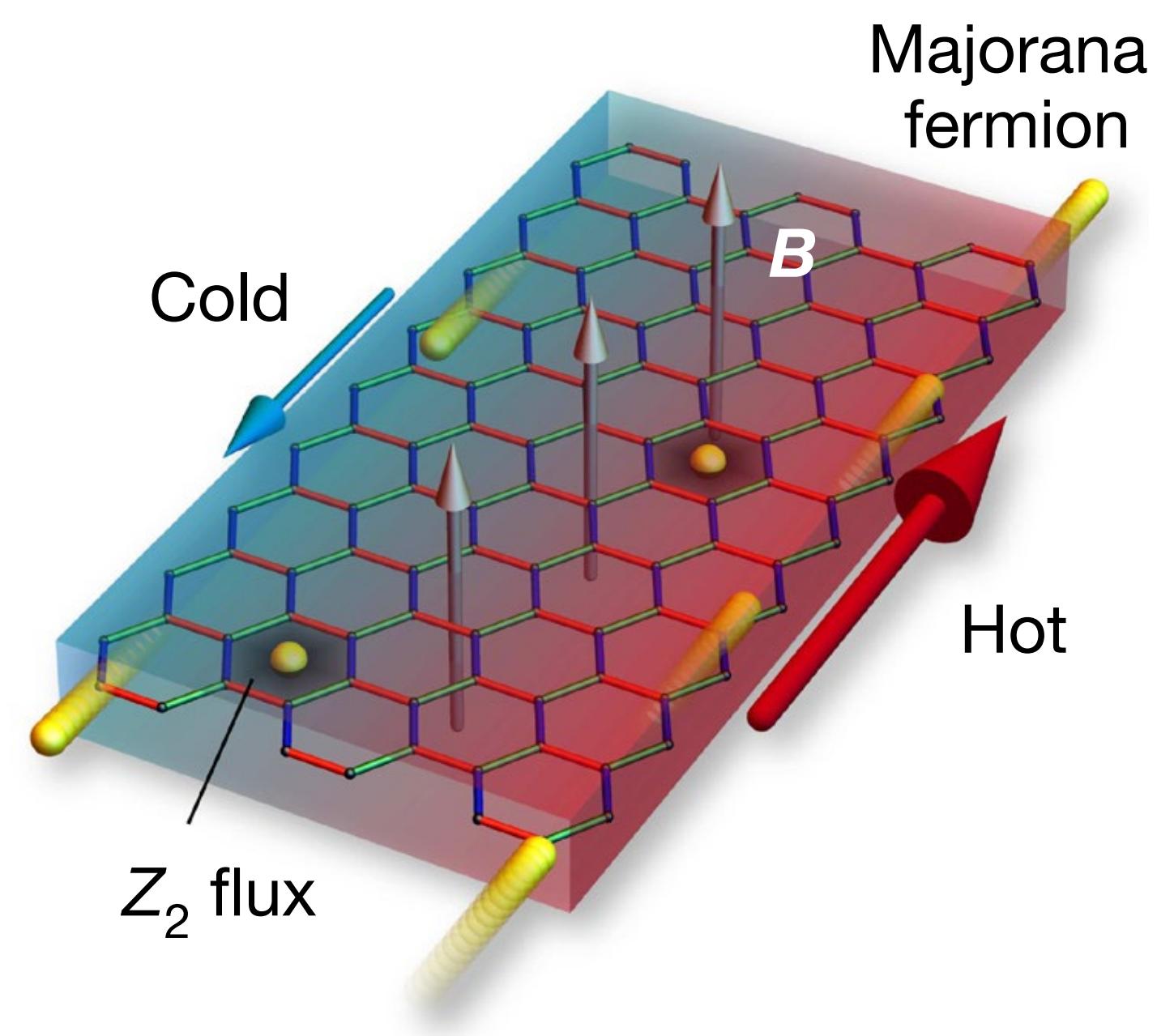
Fermion spectrum:



Review: [Trebst, arXiv:1701.07056]

# Experimental search: $\alpha$ -RuCl<sub>3</sub> in field

Half-integer thermal Quantum Hall effect:



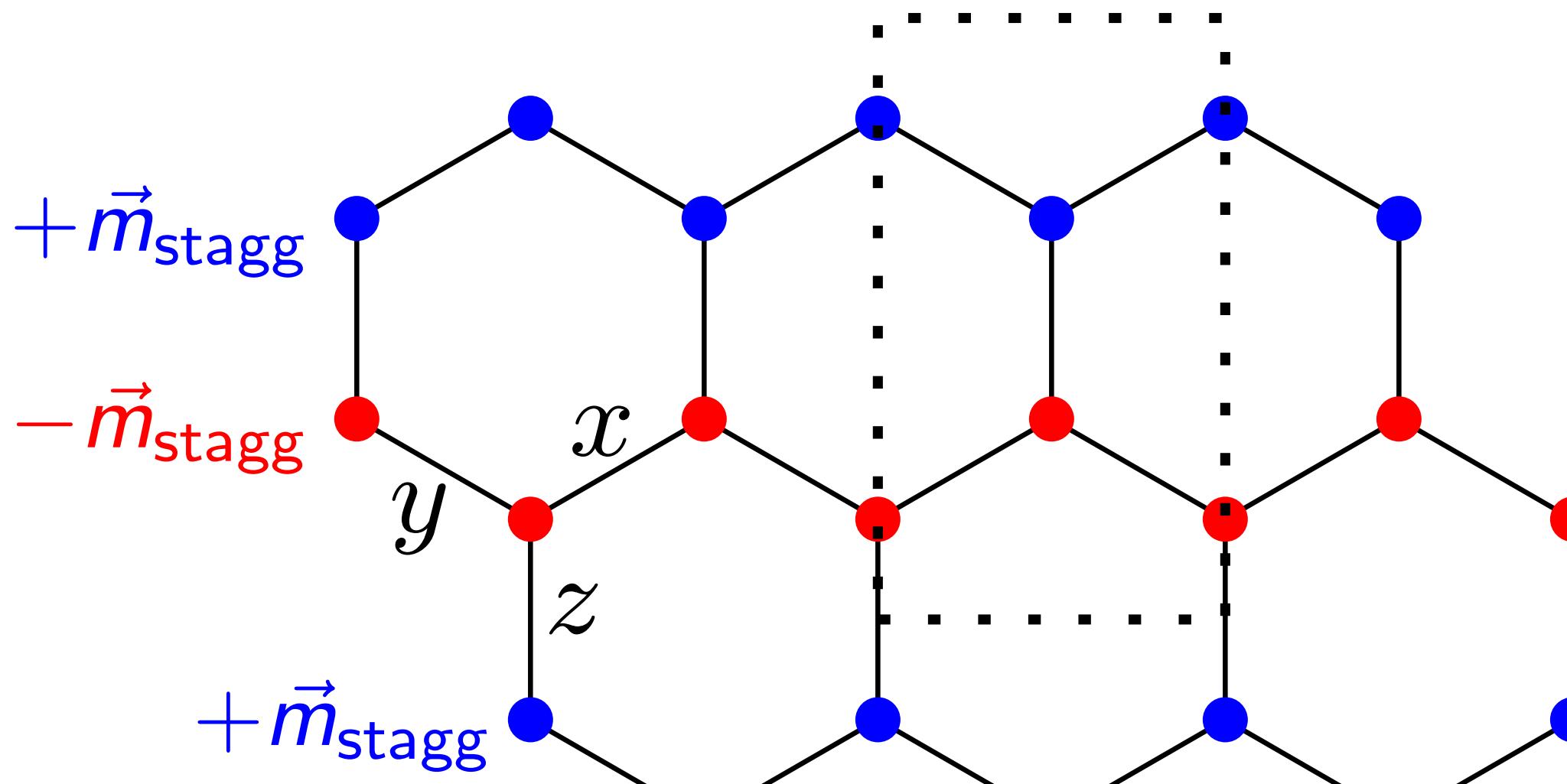
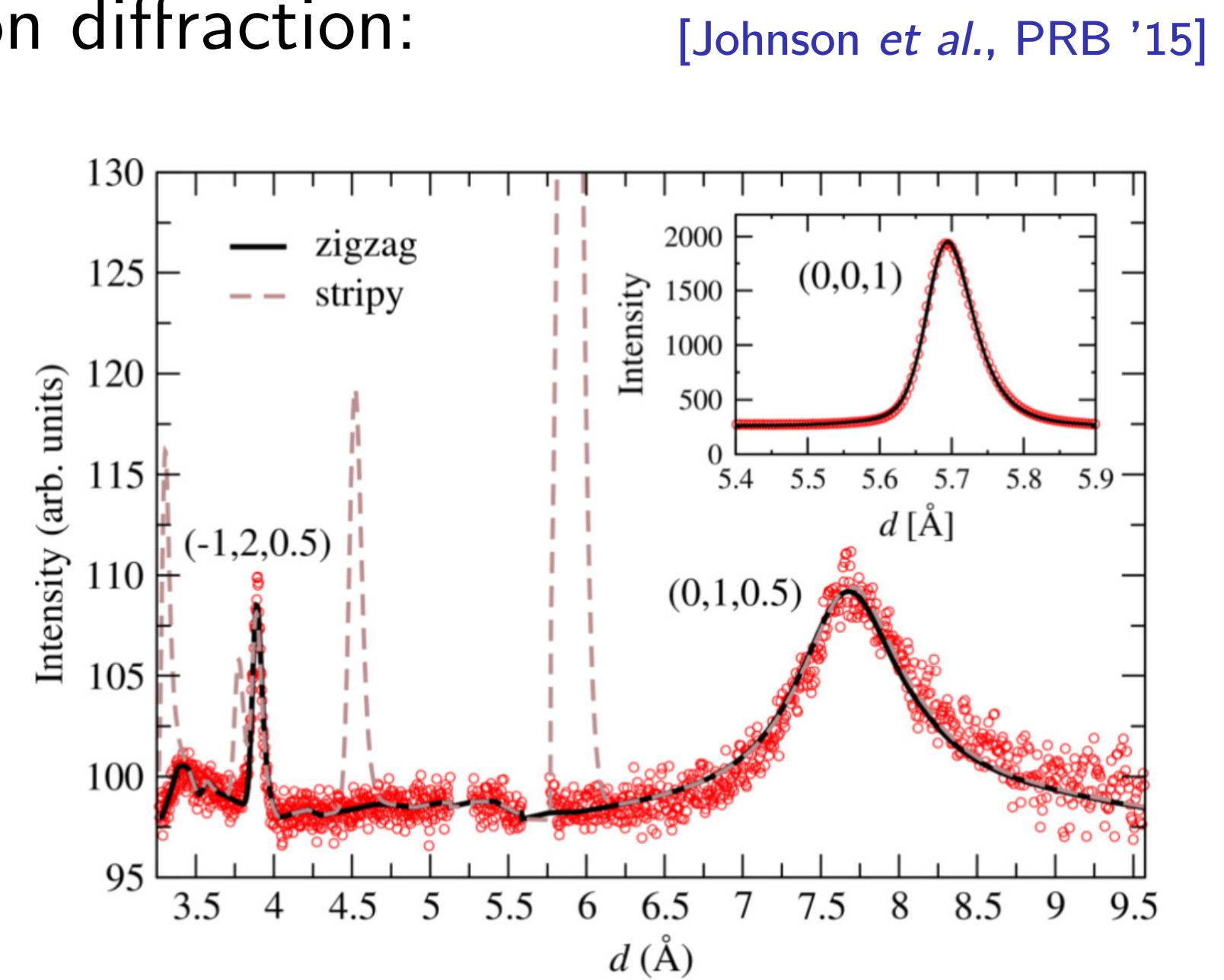
[Kasahara et al., Nature '18]

Topical Review: [LJ & Vojta, JPCM '19]

Smoking-gun signature of Majorana edge states?

# $\alpha$ -RuCl<sub>3</sub> in zero field: Zigzag antiferromagnet

Neutron diffraction:



Extended Heisenberg-Kitaev models:

$$\mathcal{H} = \sum_{\langle ij \rangle} \left[ J_1 \vec{S}_i \cdot \vec{S}_j + K_1 S_i^\gamma S_j^\gamma + \Gamma_1 (S_i^\alpha S_j^\beta + S_i^\beta S_j^\alpha) \right] + \sum_{\langle\langle ij \rangle\rangle} J_3 \vec{S}_i \cdot \vec{S}_j + \dots$$

[Winter *et al.*, PRB '16]

[LJ, Andrade, Vojta, PRB '17]

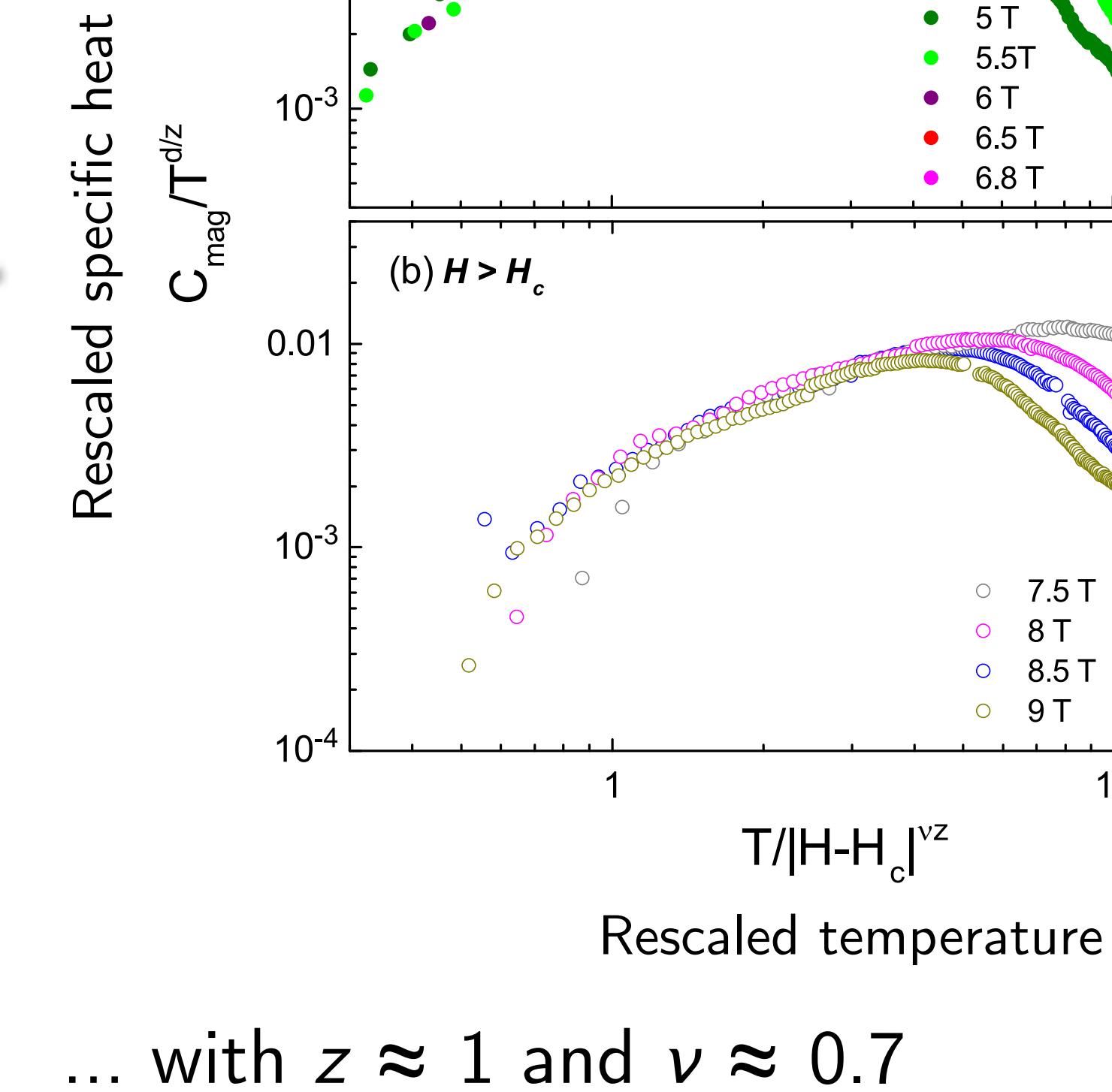
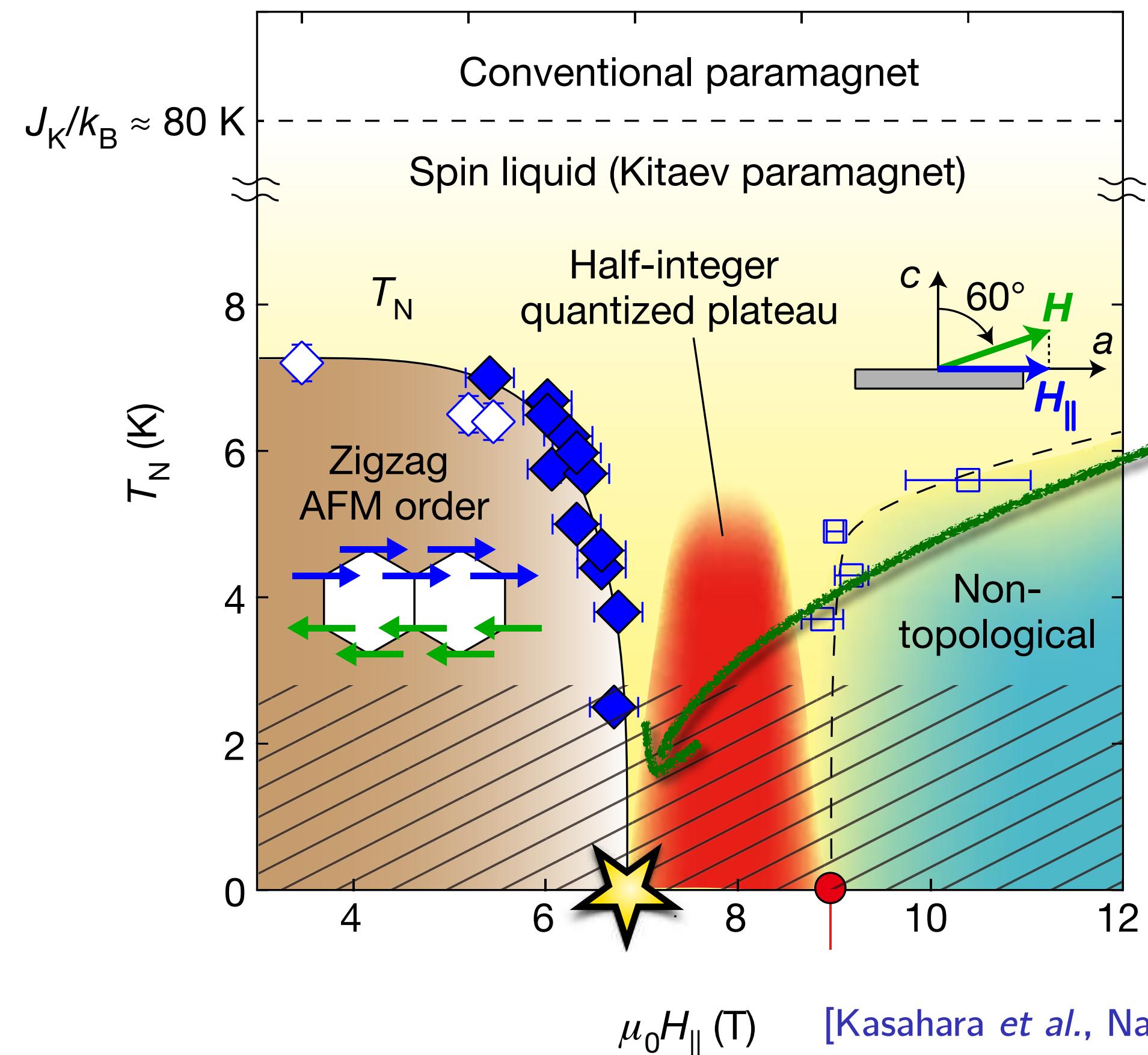
[Kaib, Winter, Valentí, PRB '19]

[LJ, Koch, Vojta, PRB '20]

[Maksimov & Chernyshev, PRR '20]

...

# $\alpha$ -RuCl<sub>3</sub>: Field-induced quantum criticality



[Wolter, Corredor, LJ, *et al.*, PRB '17]

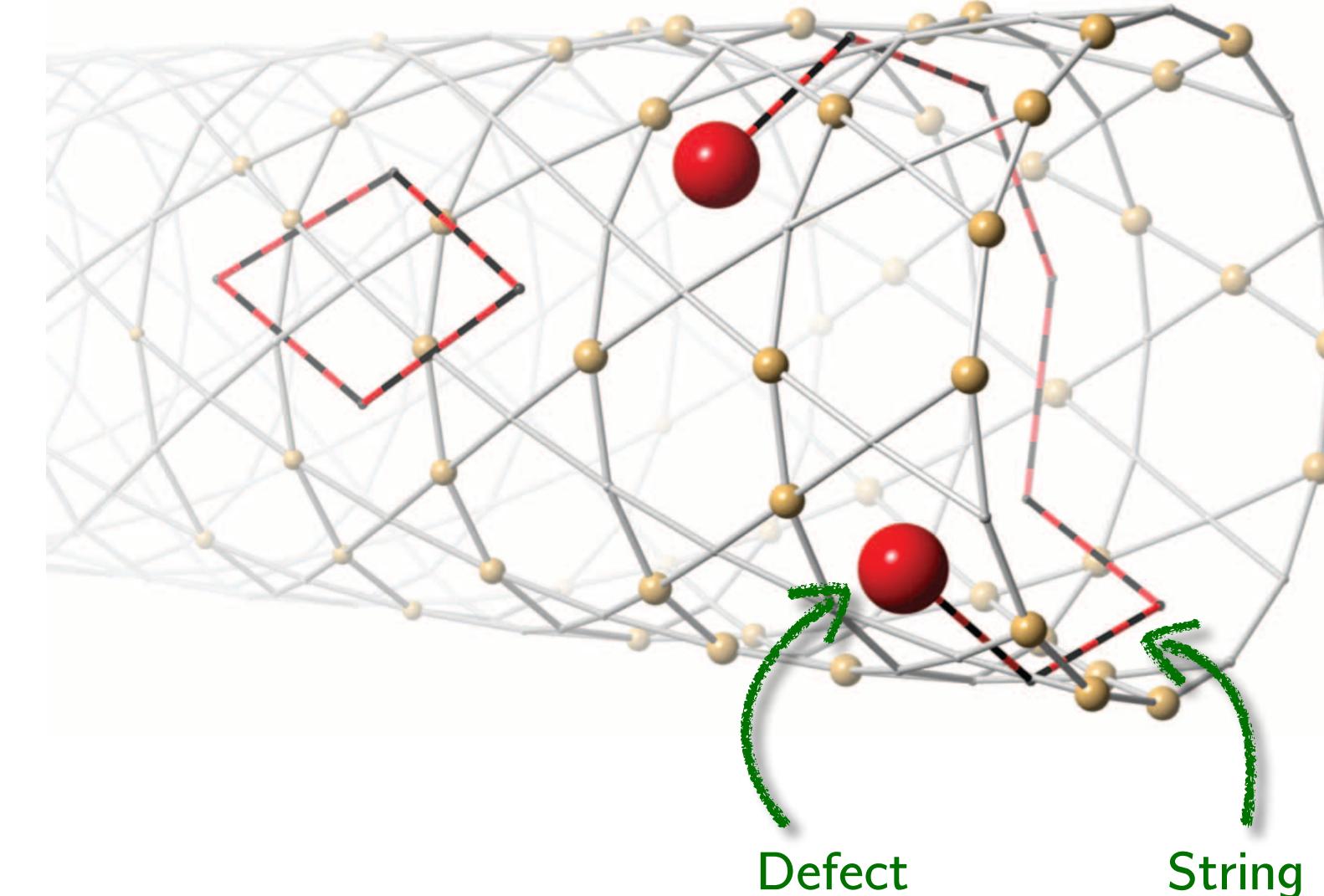
Fractionalized transition?

# Fractionalized quantum criticality: XY\*

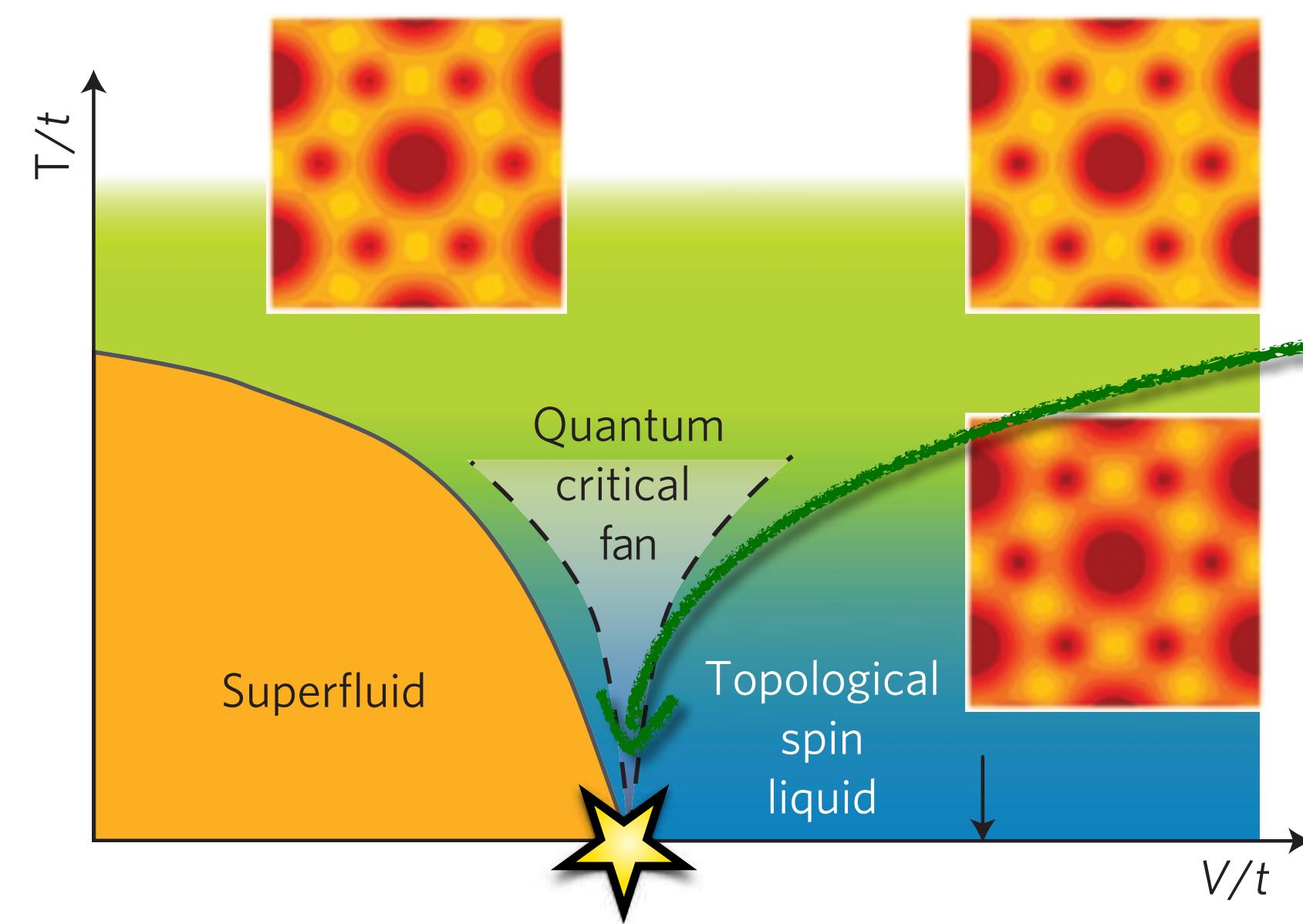
Bose-Hubbard-like model (kagome lattice):

$$\mathcal{H} = -t \sum_{\langle ij \rangle} [b_i^\dagger b_j + b_i b_j^\dagger] + V \sum_{\text{○}} (n_{\text{○}})^2$$

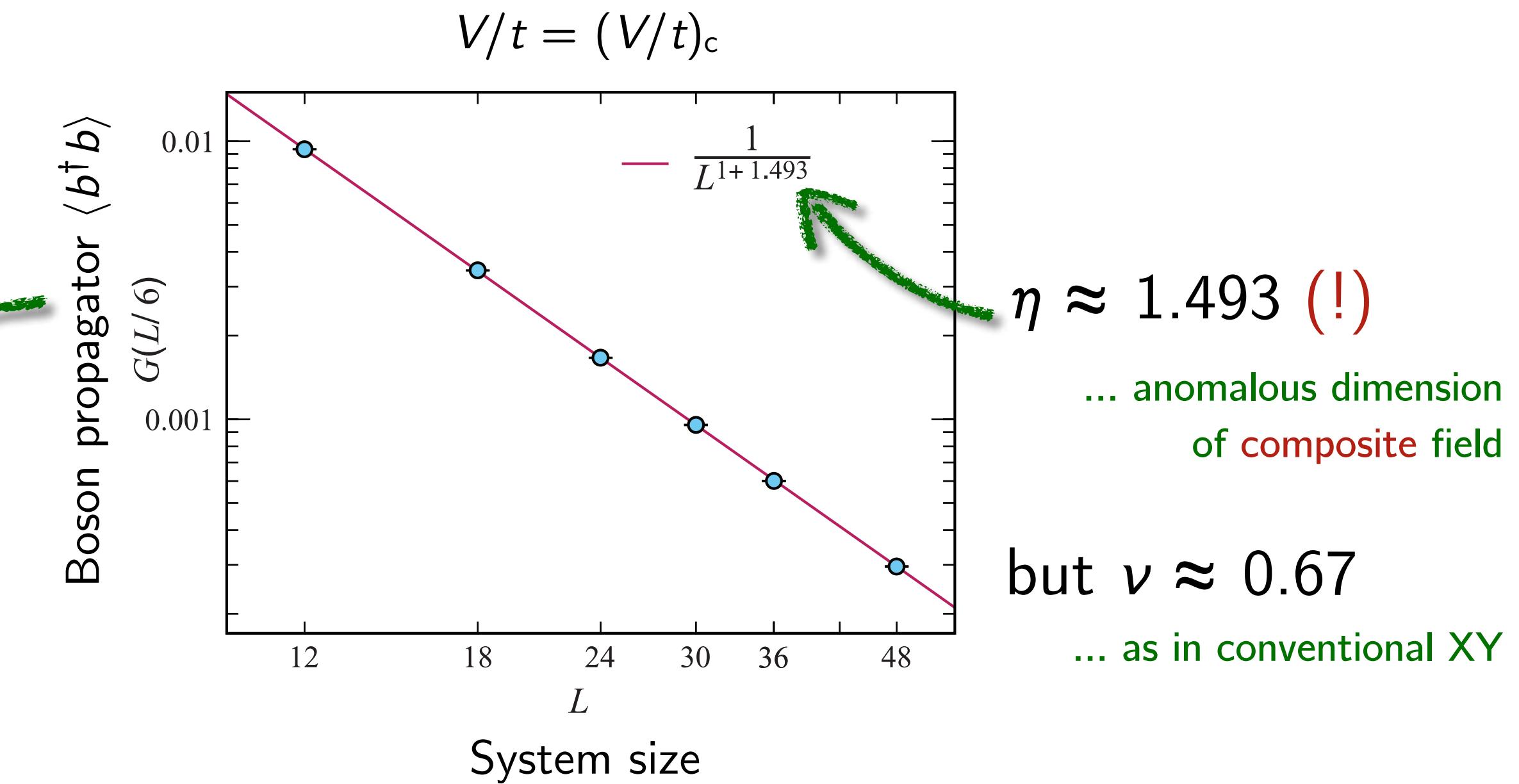
Hopping bosons      Boson density in plaquette



Phase diagram:

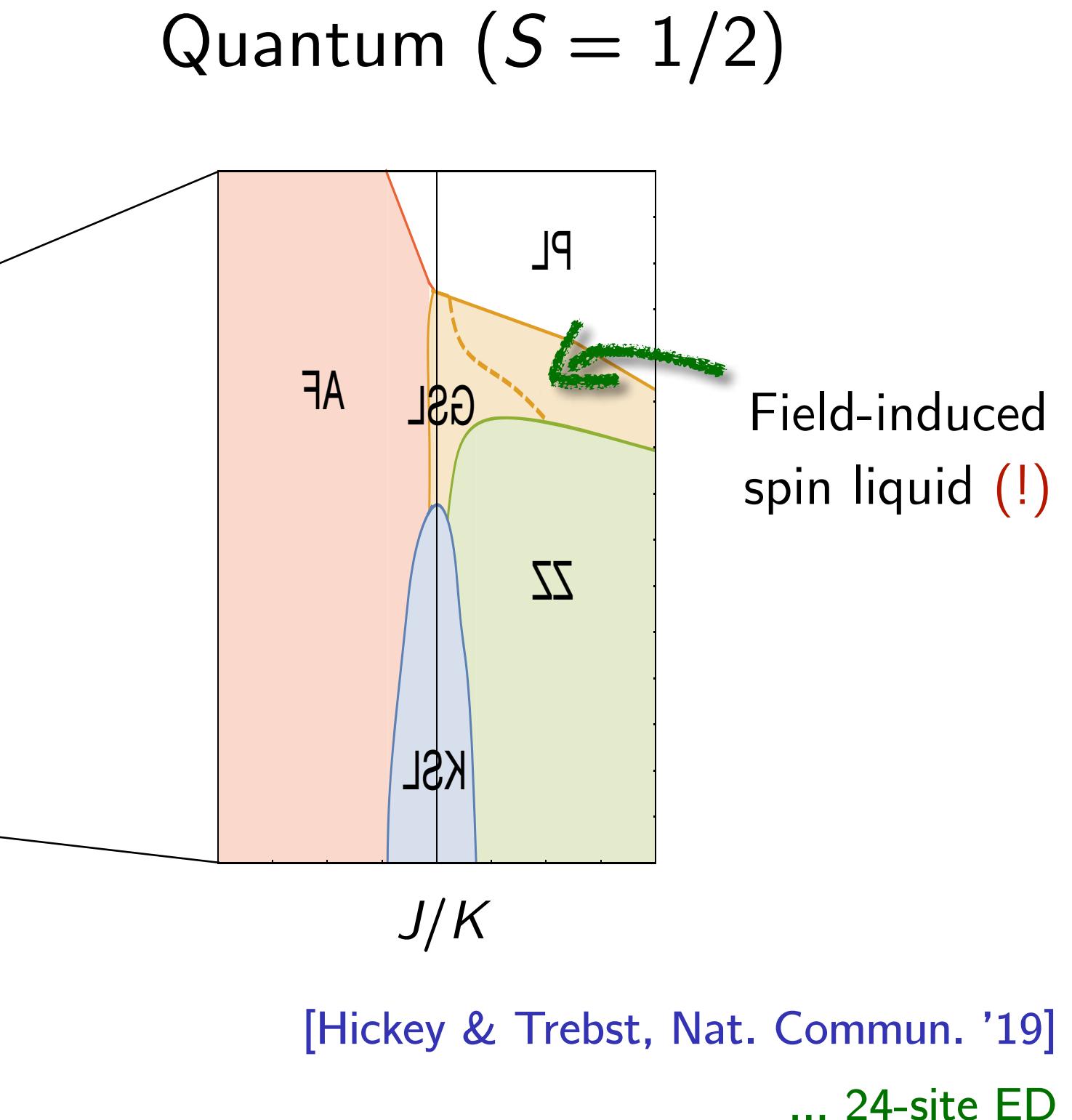
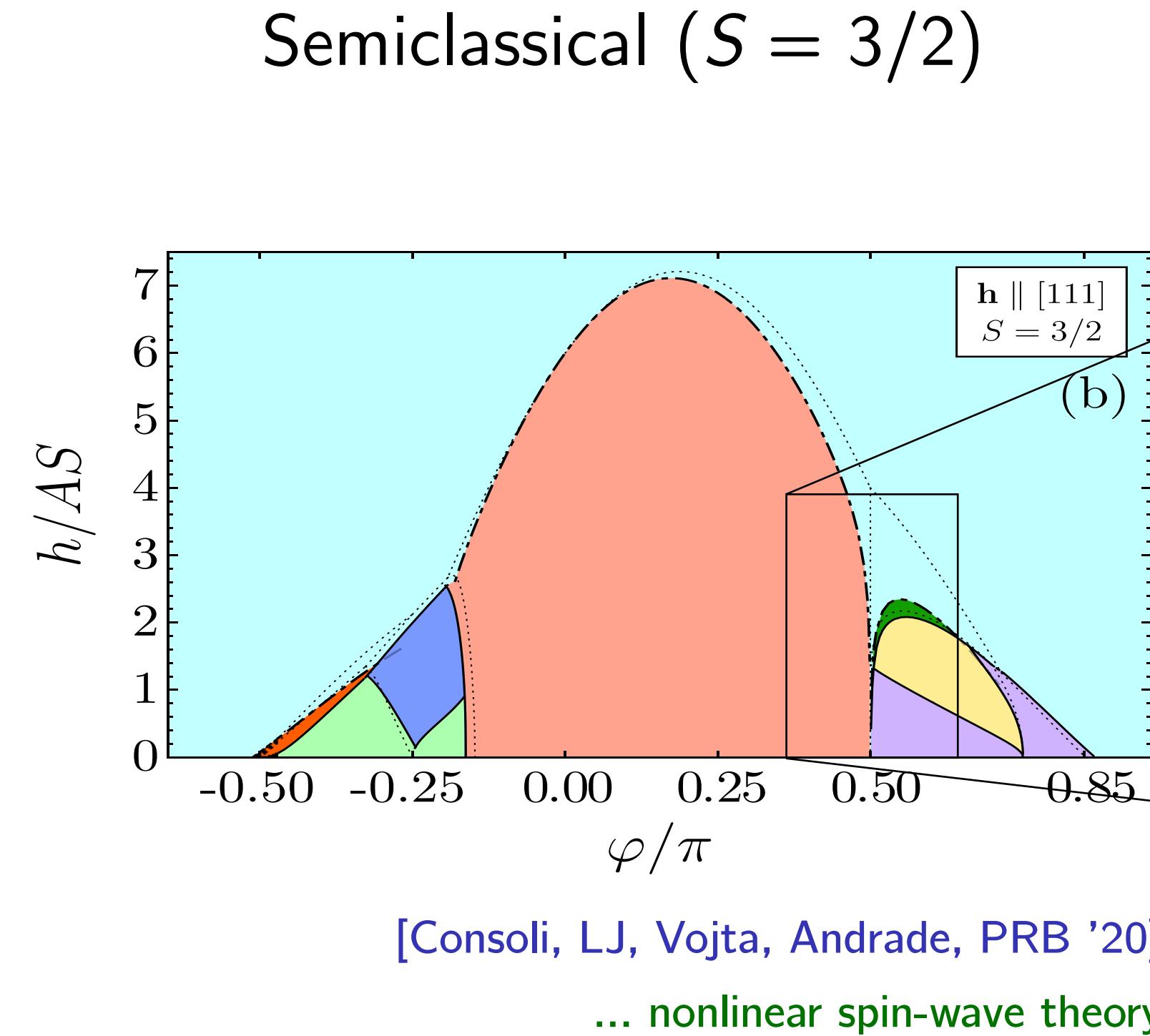
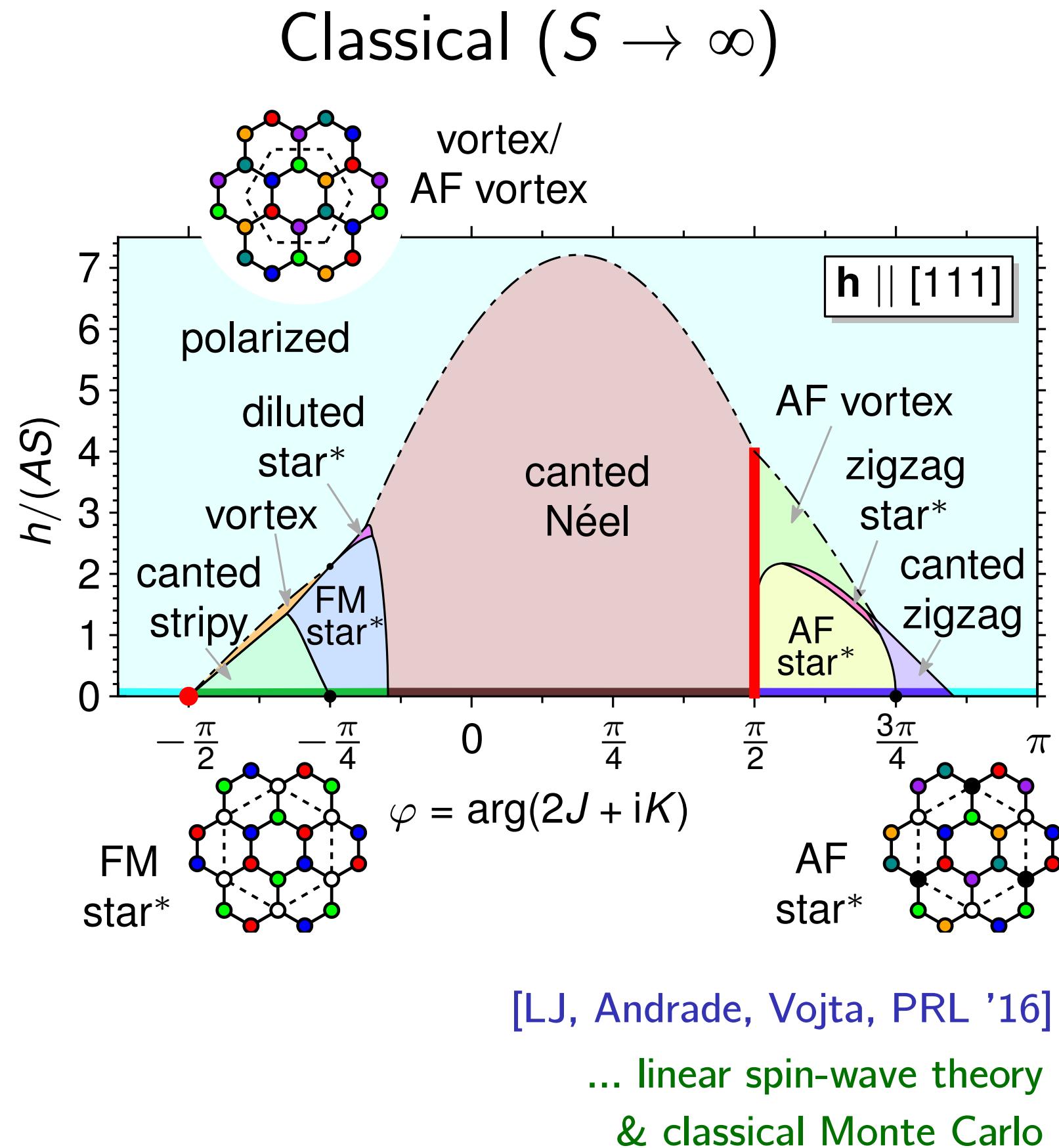


[Isakov, Hastings, Melko, Nat. Phys. '11]



[Isakov, Melko, Hastings, Science '12]  
[Chubukov, Senthil, Sachdev, PRL '94]

# Heisenberg-Kitaev models in field



Technical challenge: Dynamical  $\mathbb{Z}_2$  gauge field!

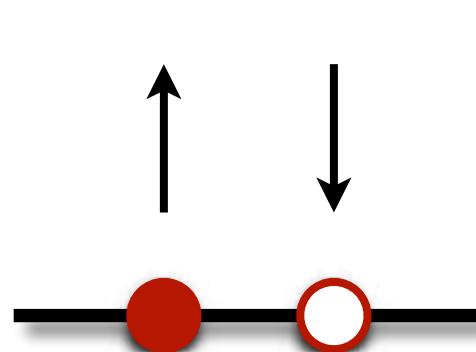
... no sign-problem-free QMC available

# Outline

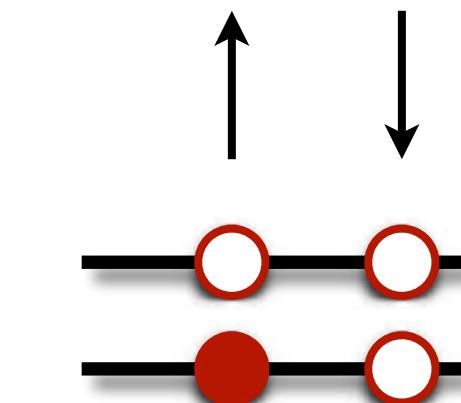
- (1) Introduction: “*Beyond-Landau*” quantum criticality
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# Generalizations of Kitaev model: Spin-orbital liquids

Spin + orbital + ... degrees of freedom:



$$\sigma^\alpha \quad 2 \times 2$$



$$\sigma^\alpha \otimes \tau^\beta = \gamma^i \quad 4 \times 4$$



...

$$\Gamma^\mu \quad 8 \times 8$$

... can realize all 16  $\mathbb{Z}_2$  topological superconductors  
[Chulliparambil, ..., LJ, Tu, arXiv:2005.13683]

Example #1 (square lattice):

$$H_K = -K \sum_{\langle ij \rangle_\gamma} (\sigma_i^x \sigma_j^x + \sigma_i^y \sigma_j^y) \otimes \tau_i^\gamma \tau_j^\gamma$$

XY spin  
↓  
Kitaev orbital  
↓

Majorana representation:

$$\sigma^y \otimes \tau^x = i b^1 c^x$$

$$\sigma^y \otimes \tau^y = i b^2 c^x$$

$$\sigma^y \otimes \tau^z = i b^3 c^x$$

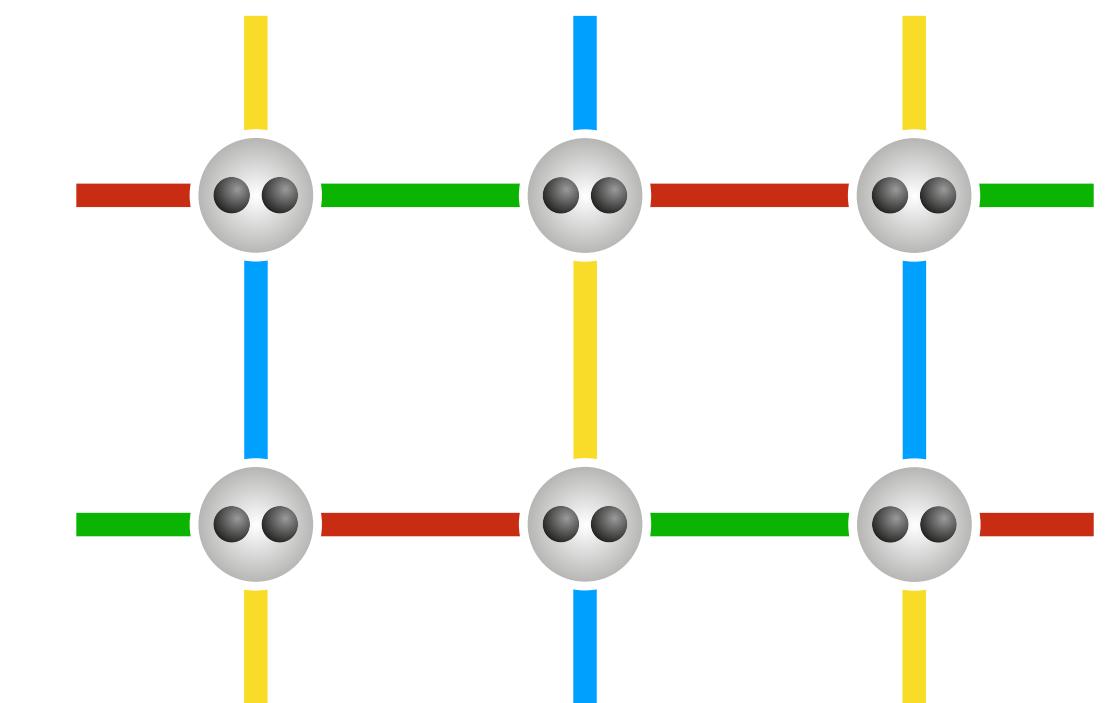
$$\sigma^x \otimes \mathbb{1} = i b^4 c^x$$

$$\sigma^z \otimes \mathbb{1} = i c^y c^x$$

... recover known model for  $j = 3/2$  spin liquid:

[Yao, Zhang, Kivelson, PRL '09]

[Nakai, Ryu, Furusaki, PRB '12]



2 itinerant fermions

# Static perturbations

“Heisenberg-Kitaev” spin-orbital model:

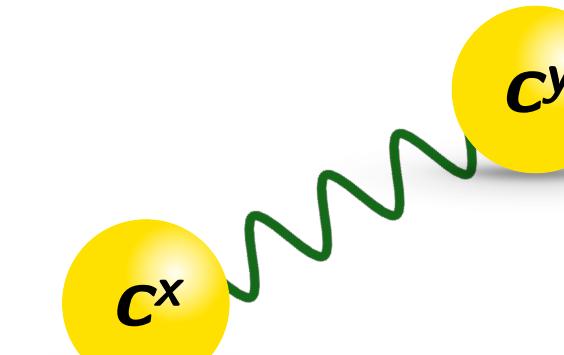
$$H = H_K + J^z \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z \otimes \mathbf{1}_i \mathbf{1}_j$$



# Static perturbations

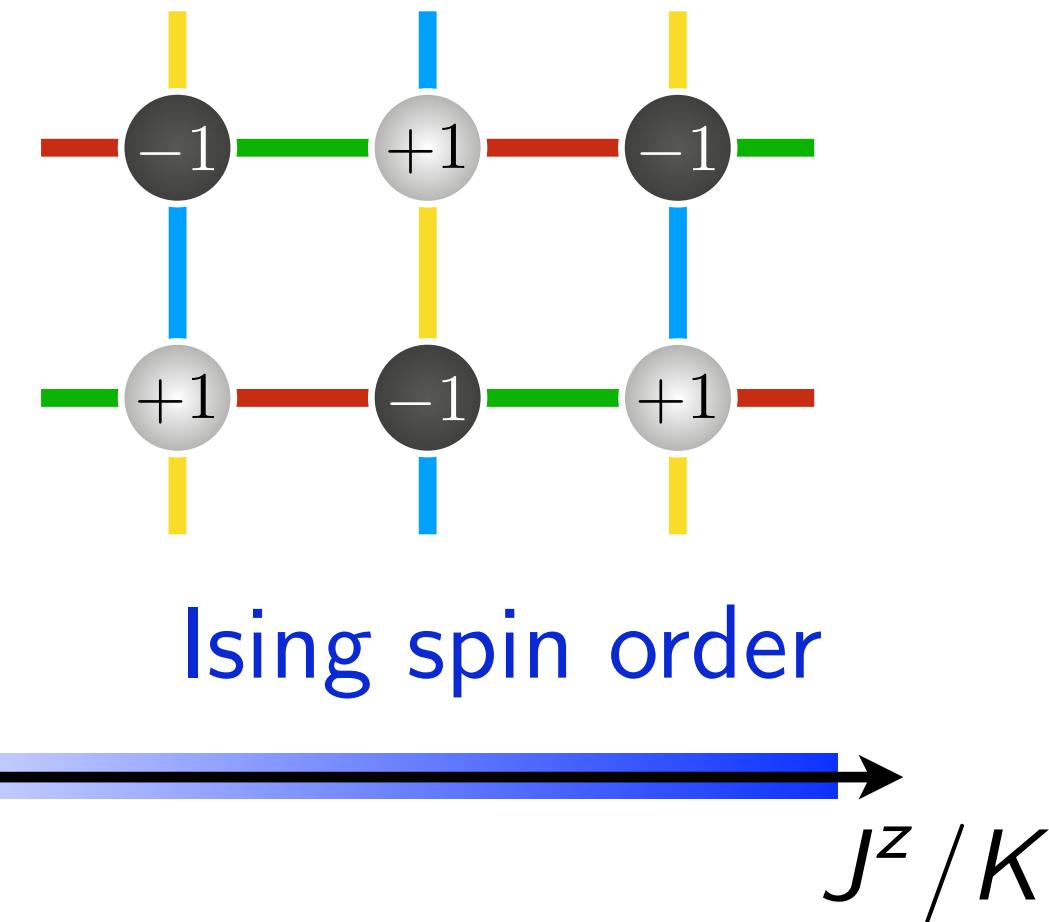
“Heisenberg-Kitaev” spin-orbital model:

$$H = H_K + J^z \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z \otimes \mathbf{1}_i \mathbf{1}_j$$



“Kitaev” spin-orbital liquid

0



Parton representation:

$$H \mapsto \sum_{\langle ij \rangle} \left[ 2K u_{ij} (f_i^\dagger f_j + f_j^\dagger f_i) + 4J^z (n_i - \frac{1}{2})(n_j - \frac{1}{2}) \right]$$

hopping parameter  $t = 2K$

$\pi$  flux

nearest-neighbor repulsion  $V = 4J^z$

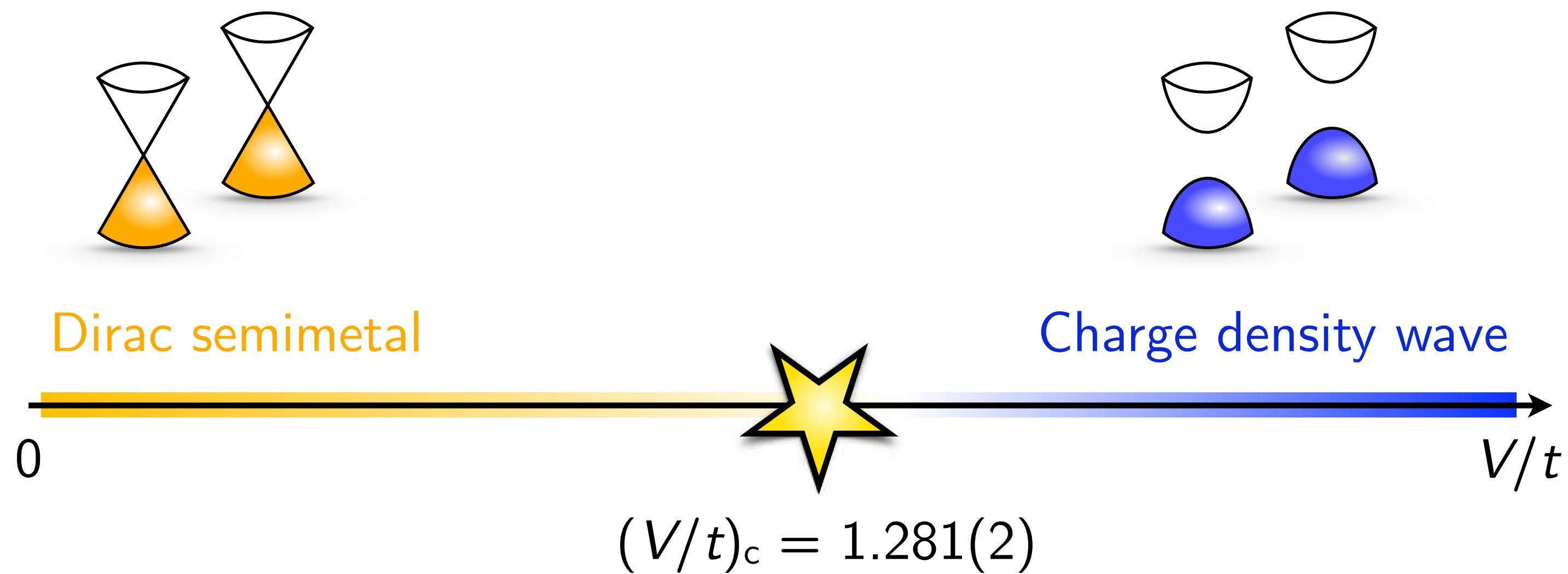
$f = \frac{1}{2}(c^x + i c^y)$

electron density  $f^\dagger f$

Ground-state flux pattern:  
[Lieb, PRL '94]

Spin-orbital model  $\mapsto$  interacting fermions on  $\pi$ -flux lattice

# Spinless fermions on $\pi$ -flux lattice: QMC



Gross-Neveu- $\mathbb{Z}_2$  universality:

$$1/\nu = 1.12(1), \quad \eta_\phi = 0.51(3), \quad \eta_\psi \approx 0.1$$

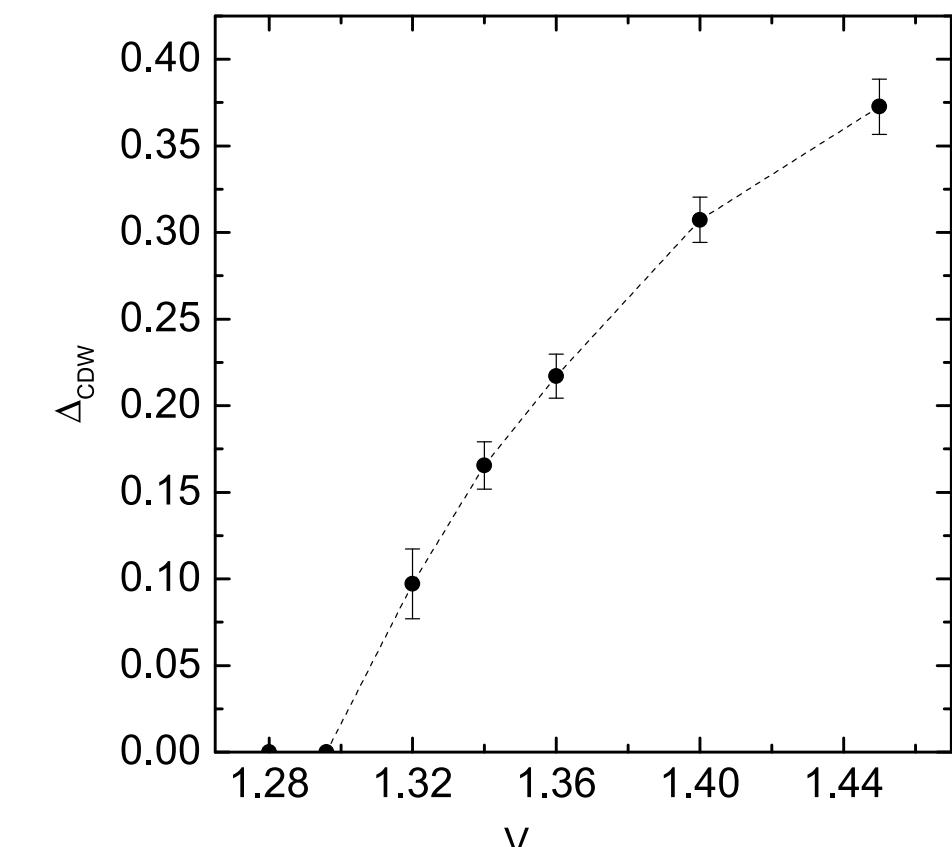
[Gracey, IJMP '94]

[LJ & Herbut, PRB '14]

[Iliesiu *et al.*, JHEP '18]

[Ihrig, Mihaila, Scherer, PRB '18]

...

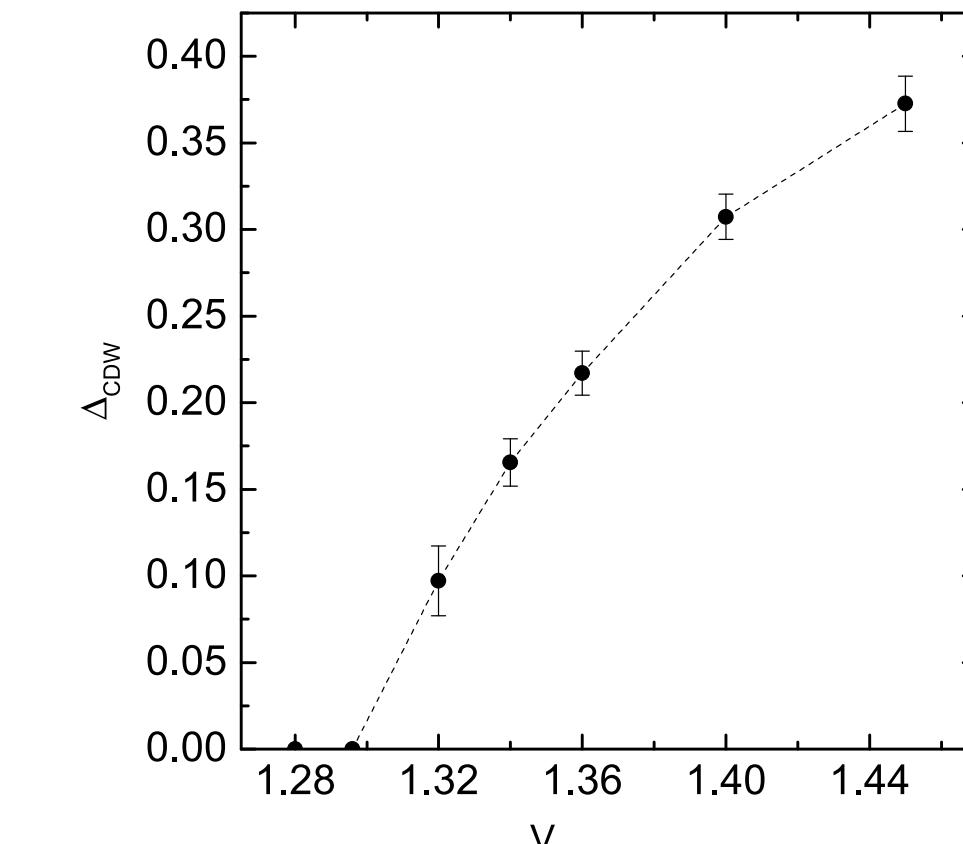
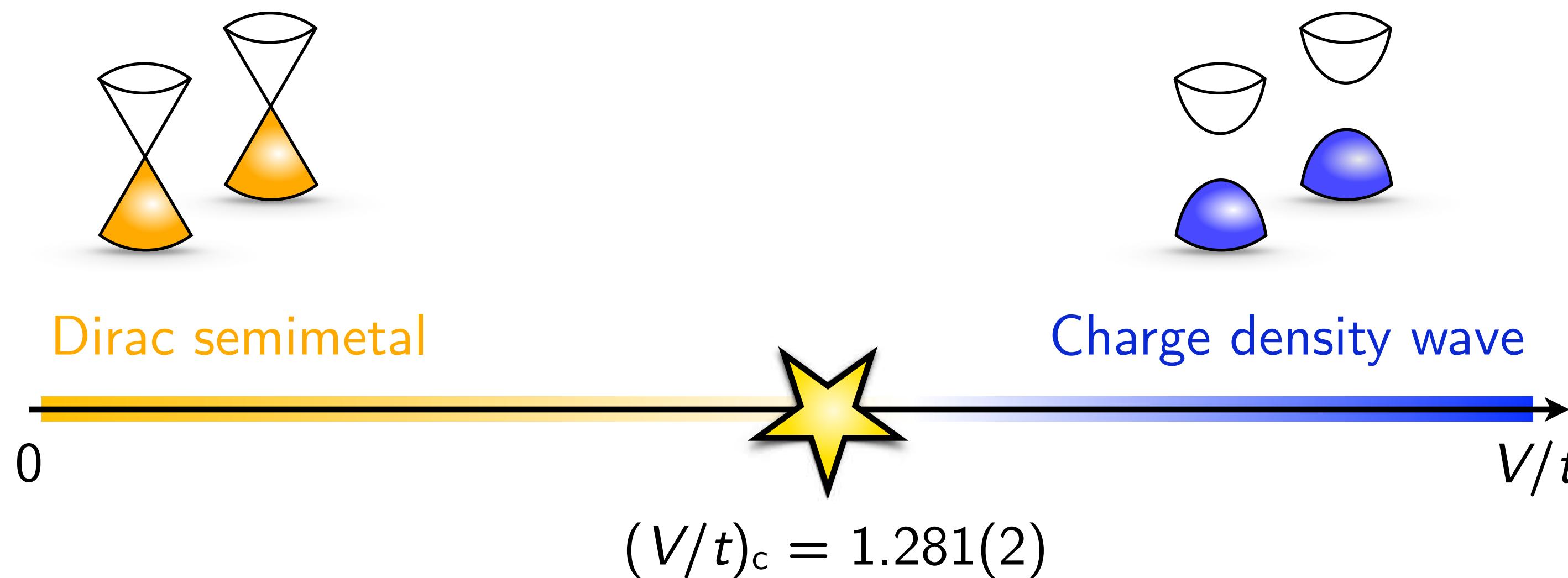


[Wang, Corboz, Troyer, NJP '14]

[Li, Jiang, Yao, NJP '15]

[Huffman & Chandrasekharan, PRD '17; PRD '20]

# Spinless fermions on $\pi$ -flux lattice: QMC



[Wang, Corboz, Troyer, NJP '14]

[Li, Jiang, Yao, NJP '15]

[Huffman & Chandrasekharan, PRD '17; PRD '20]

Gross-Neveu- $\mathbb{Z}_2$  universality:  $1/\nu = 1.12(1)$ ,  $\eta_\phi = 0.51(3)$ ,  $\eta_\psi \approx 0.1$

[Gracey, IJMP '94]

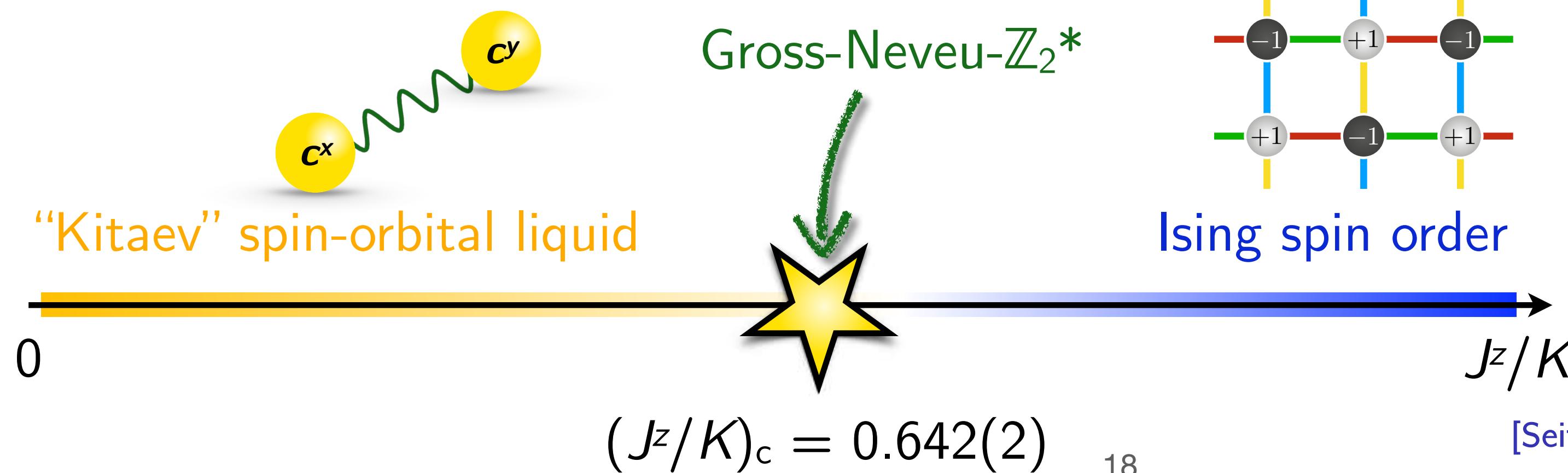
[LJ & Herbut, PRB '14]

[Iliesiu *et al.*, JHEP '18]

[Ihrig, Mihaila, Scherer, PRB '18]

...

Spin-orbital model:

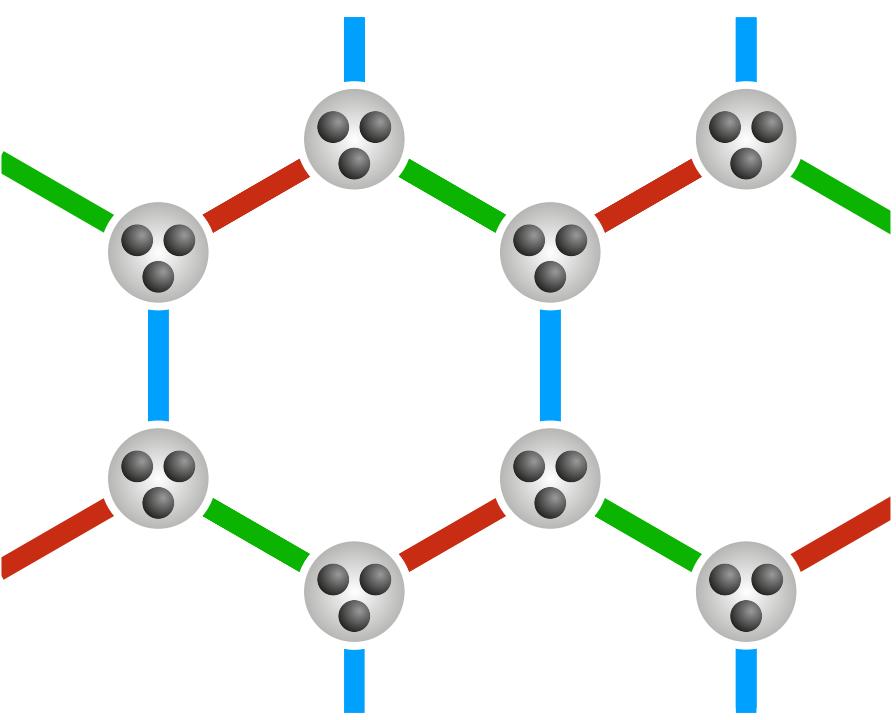


[Seifert, Dong, Chulliparambil, Vojta, Tu, LJ, arXiv:2009.05051]

# Fractionalized fermionic universality classes

Example #2 (honeycomb lattice):

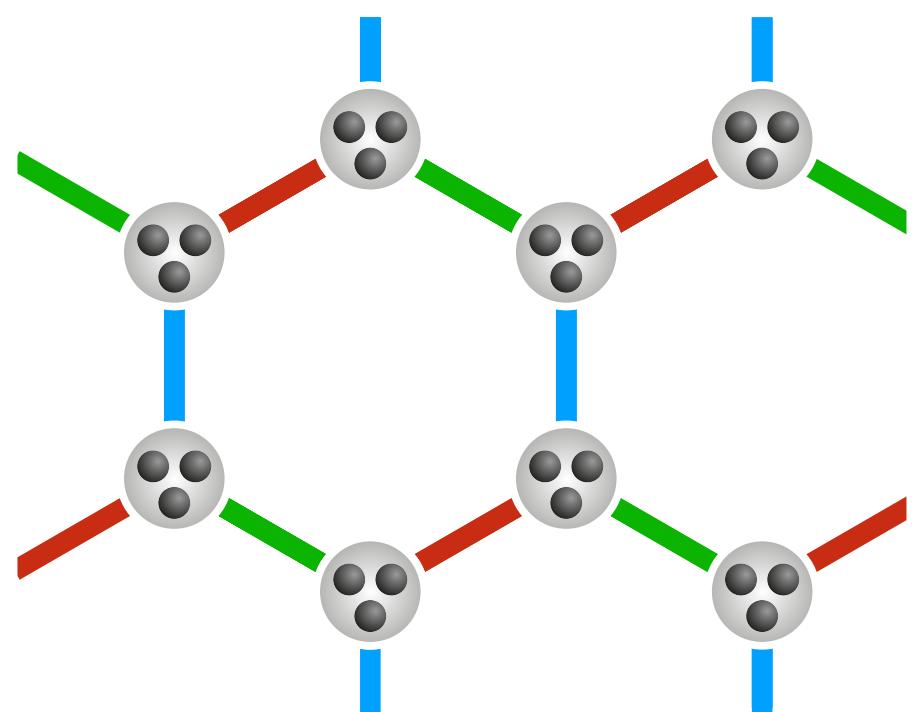
$$H = -K \sum_{\langle ij \rangle_\gamma} \vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \tau_i^\gamma \tau_j^\gamma + J \sum_{\langle ij \rangle} \vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \mathbb{1}_i \mathbb{1}_j$$



# Fractionalized fermionic universality classes

Example #2 (honeycomb lattice):

$$H = -K \sum_{\langle ij \rangle_\gamma} \vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \tau_i^\gamma \tau_j^\gamma + J \sum_{\langle ij \rangle} \vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \mathbb{1}_i \mathbb{1}_j$$



3 itinerant fermions

Majorana representation:

$$H \mapsto \sum_{\langle ij \rangle} \left[ K i u_{ij} c_i^\top c_j + \frac{J}{4} (\vec{c}_i^\top \vec{L} c_i) \cdot (\vec{c}_j^\top \vec{L} c_j) \right]$$

$c \equiv \begin{pmatrix} c^x \\ c^y \\ c^z \end{pmatrix}$

*spin-1 matrices*

*0 flux*

Ordered state:

$$\langle c_{iA}^\top \vec{L} c_{iA} \rangle \neq \langle c_{jB}^\top \vec{L} c_{jB} \rangle$$

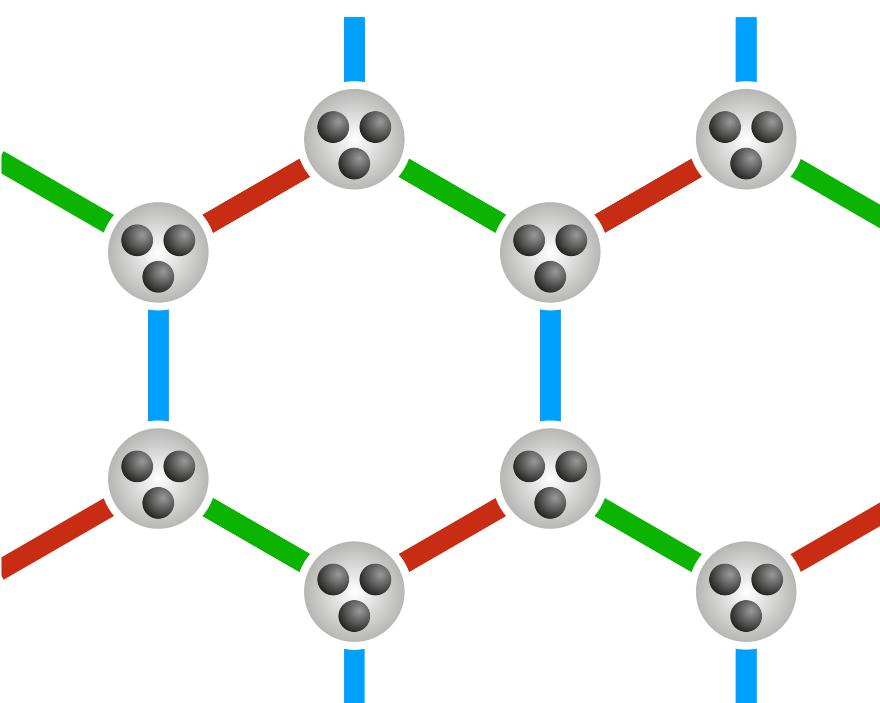
spin density wave



# Fractionalized fermionic universality classes

Example #2 (honeycomb lattice):

$$H = -K \sum_{\langle ij \rangle_\gamma} \vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \tau_i^\gamma \tau_j^\gamma + J \sum_{\langle ij \rangle} \vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \mathbb{1}_i \mathbb{1}_j$$



3 itinerant fermions

Majorana representation:

$$H \mapsto \sum_{\langle ij \rangle} \left[ K i u_{ij} c_i^\top c_j + \frac{J}{4} (\vec{c}_i^\top \vec{L} c_i) \cdot (\vec{c}_j^\top \vec{L} c_j) \right]$$

*c*  $\equiv \begin{pmatrix} c^x \\ c^y \\ c^z \end{pmatrix}$

*spin-1 matrices*

*0 flux*

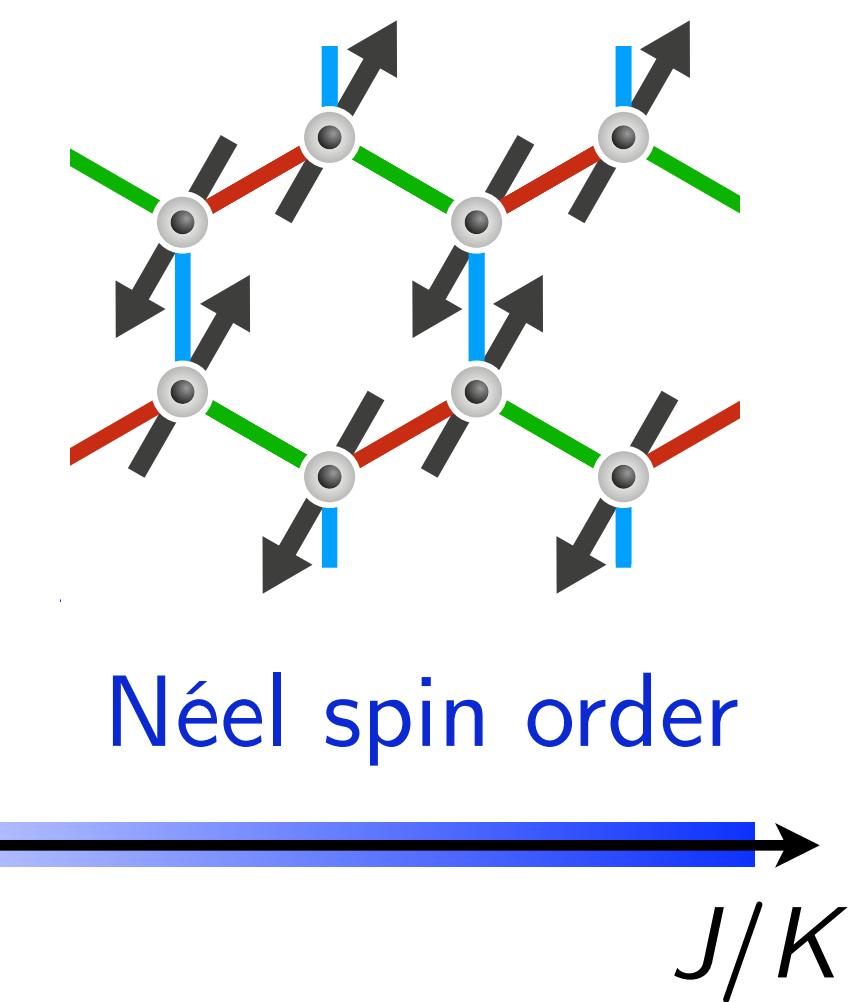
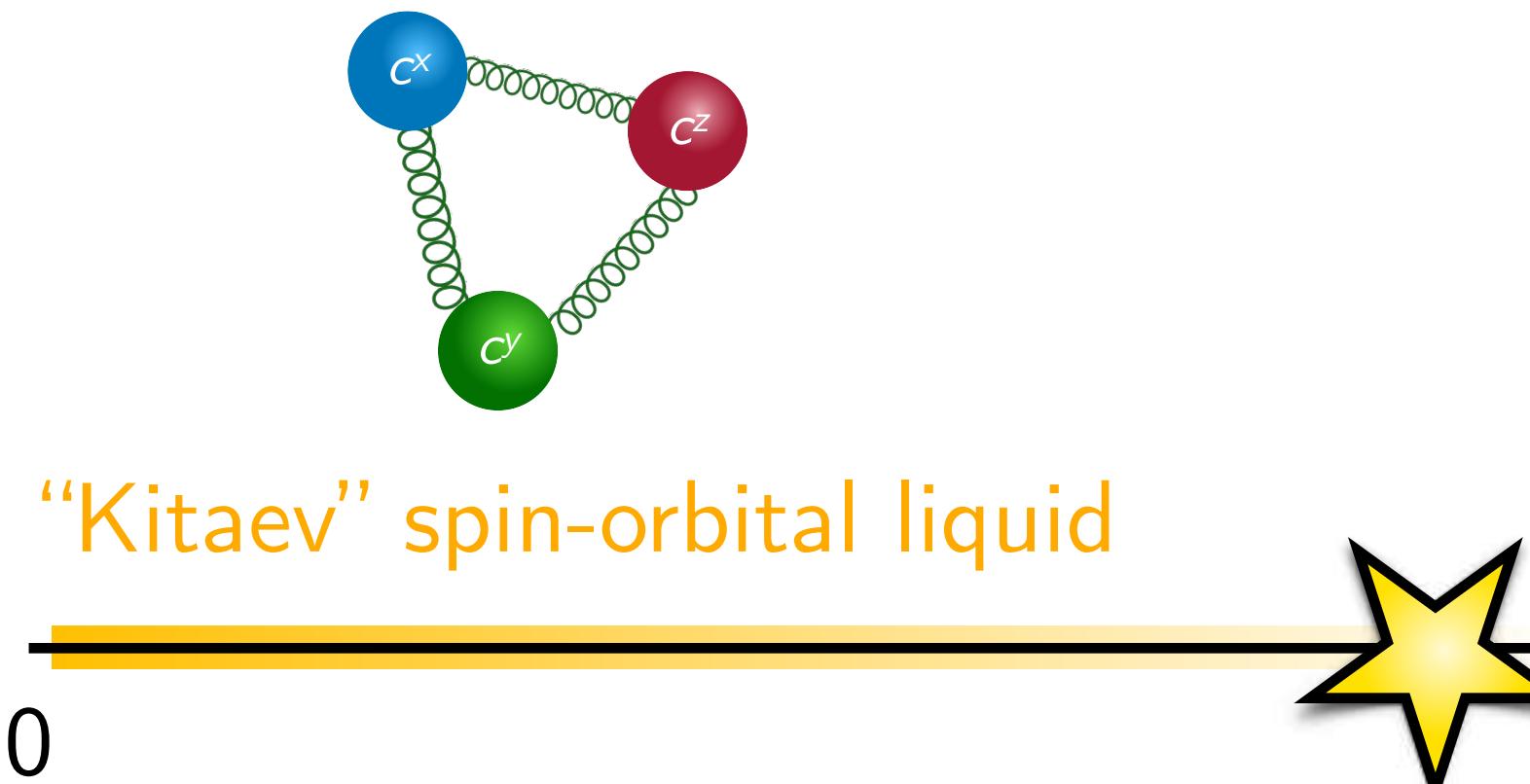
Ordered state:

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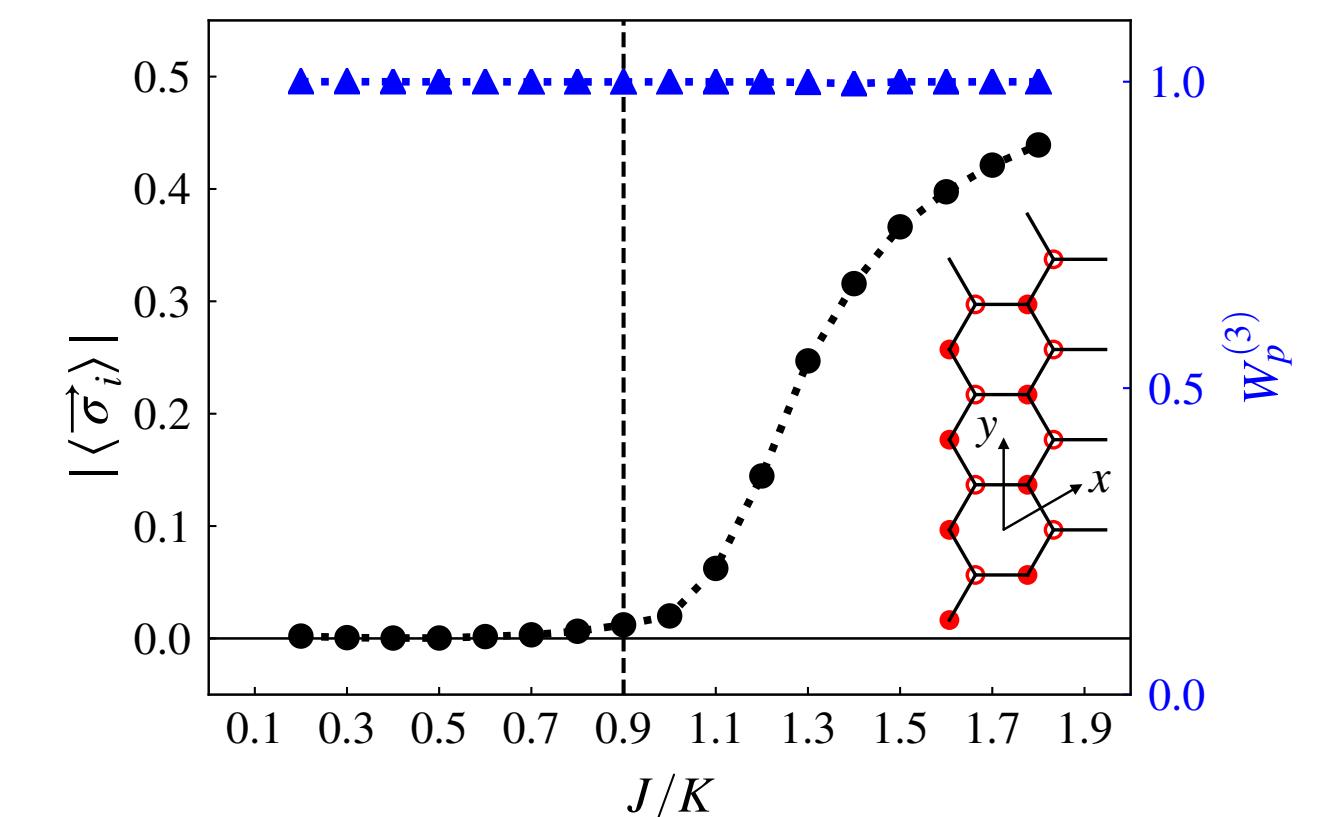
spin density wave



Phase diagram:



DMRG:

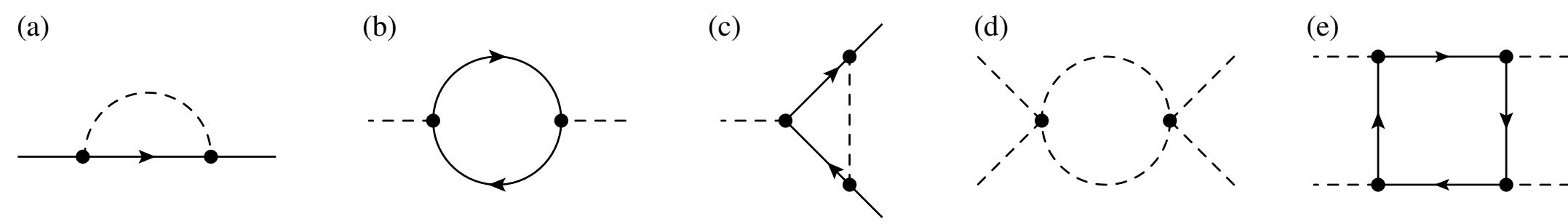


# Effective field theory: Gross-Neveu-SO(3)

Action:

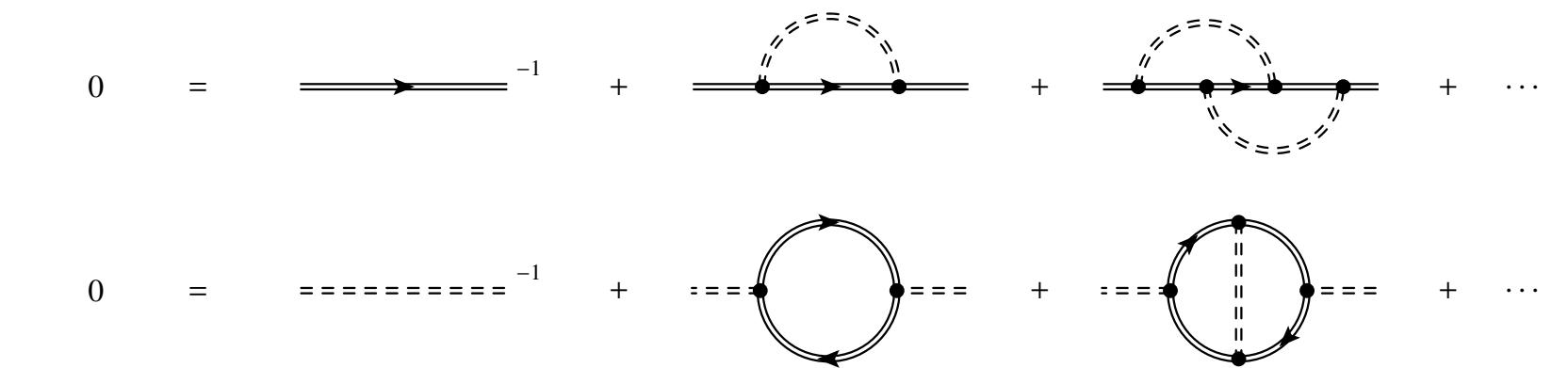
$$\mathcal{S} = \int d^2\vec{x}d\tau \left[ \bar{\psi}\gamma^\mu\partial_\mu\psi + g\vec{\varphi}\cdot\bar{\psi}(1_2 \otimes \vec{L})\psi + \frac{1}{2}\vec{\varphi}\cdot(-\partial_\mu^2 + m^2)\vec{\varphi} + \lambda(\vec{\varphi}\cdot\vec{\varphi})^2 \right]$$

$4 - \varepsilon$  expansion:



$$1/\nu \approx 1.1, \quad \eta_\phi \approx 0.33, \quad \eta_\psi \approx 0.17$$

$1/N$  expansion:



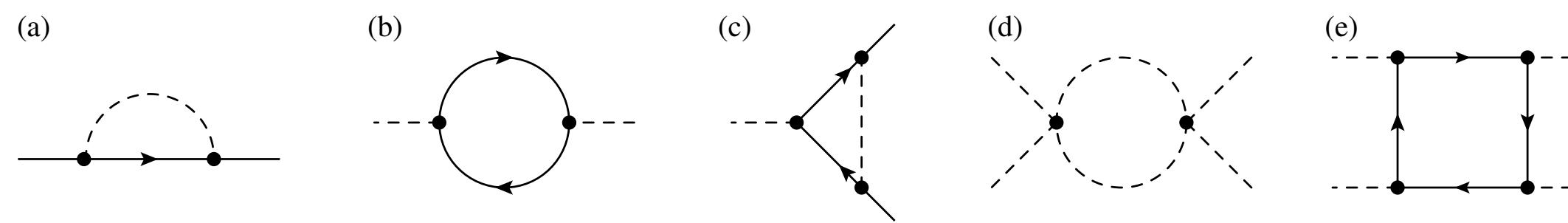
$$1/\nu \approx 0.5, \quad \eta_\phi \approx 0.32, \quad \eta_\psi \approx 0.14$$

# Effective field theory: Gross-Neveu-SO(3)

Action:

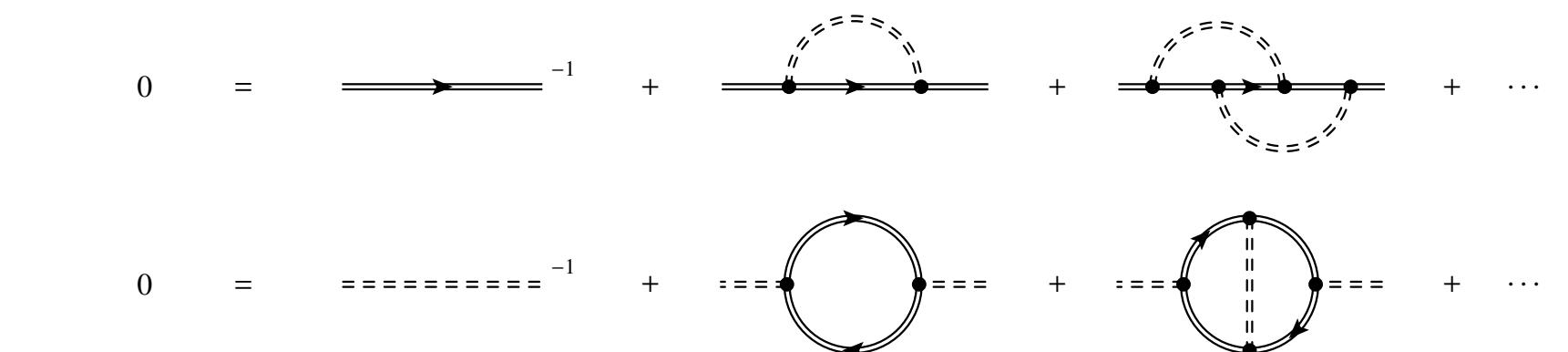
$$\mathcal{S} = \int d^2\vec{x}d\tau \left[ \bar{\psi}\gamma^\mu\partial_\mu\psi + g\vec{\varphi}\cdot\bar{\psi}(1_2 \otimes \vec{L})\psi + \frac{1}{2}\vec{\varphi}\cdot(-\partial_\mu^2 + m^2)\vec{\varphi} + \lambda(\vec{\varphi}\cdot\vec{\varphi})^2 \right]$$

$4 - \varepsilon$  expansion:

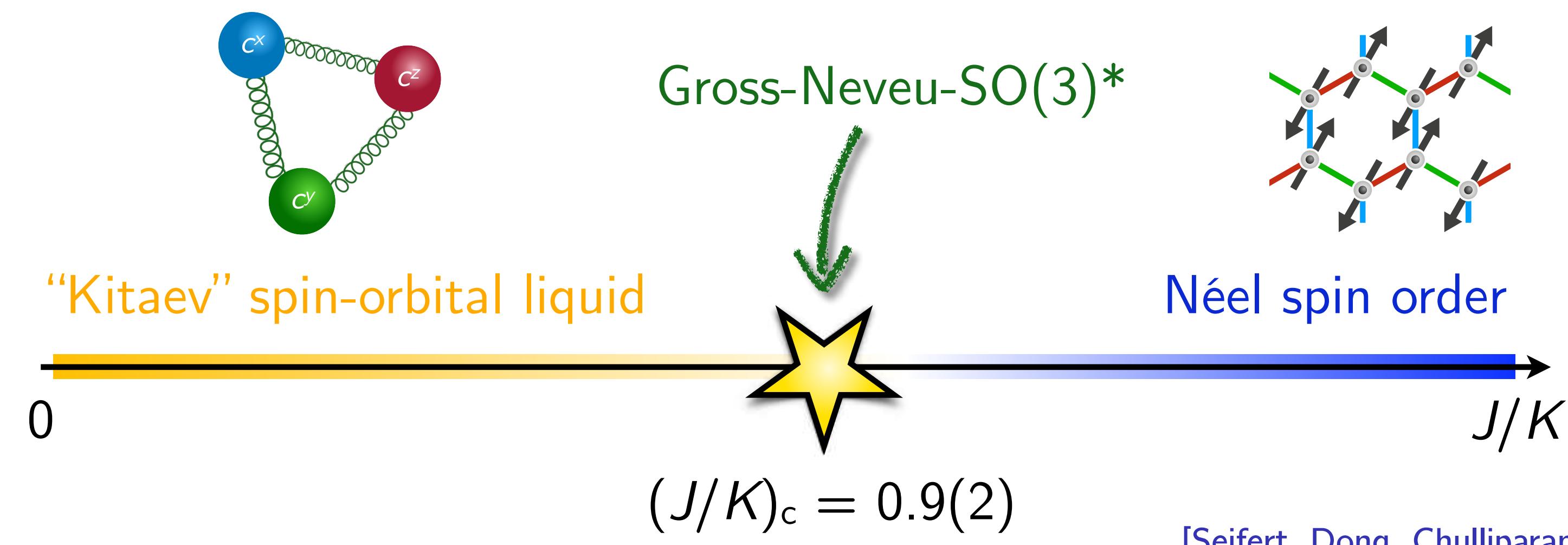


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[Seifert, Dong, Chulliparambil, Vojta, Tu, LJ, arXiv:2009.05051]

# Finite-size spectroscopy: Ising vs Ising\*

Transverse-field Ising:

$$H = -J \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z - h \sum_i \sigma_i^x$$

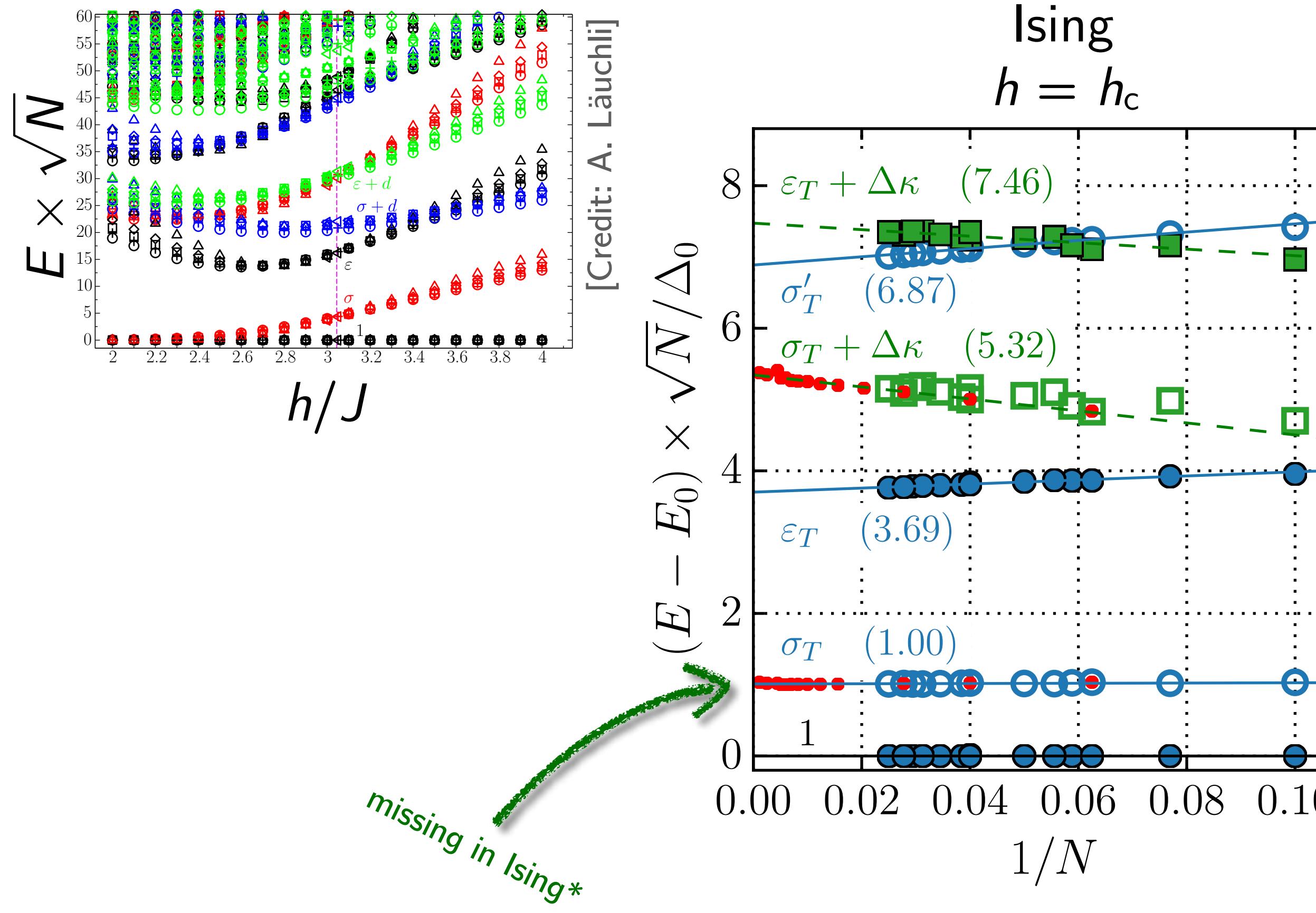
Transverse-field toric code:

$$H = -J \sum_s \prod_{i \in s} \sigma_i^x - J \sum_p \prod_{i \in p} \sigma_i^z - h \sum_i \sigma_i^x$$

# Finite-size spectroscopy: Ising vs Ising\*

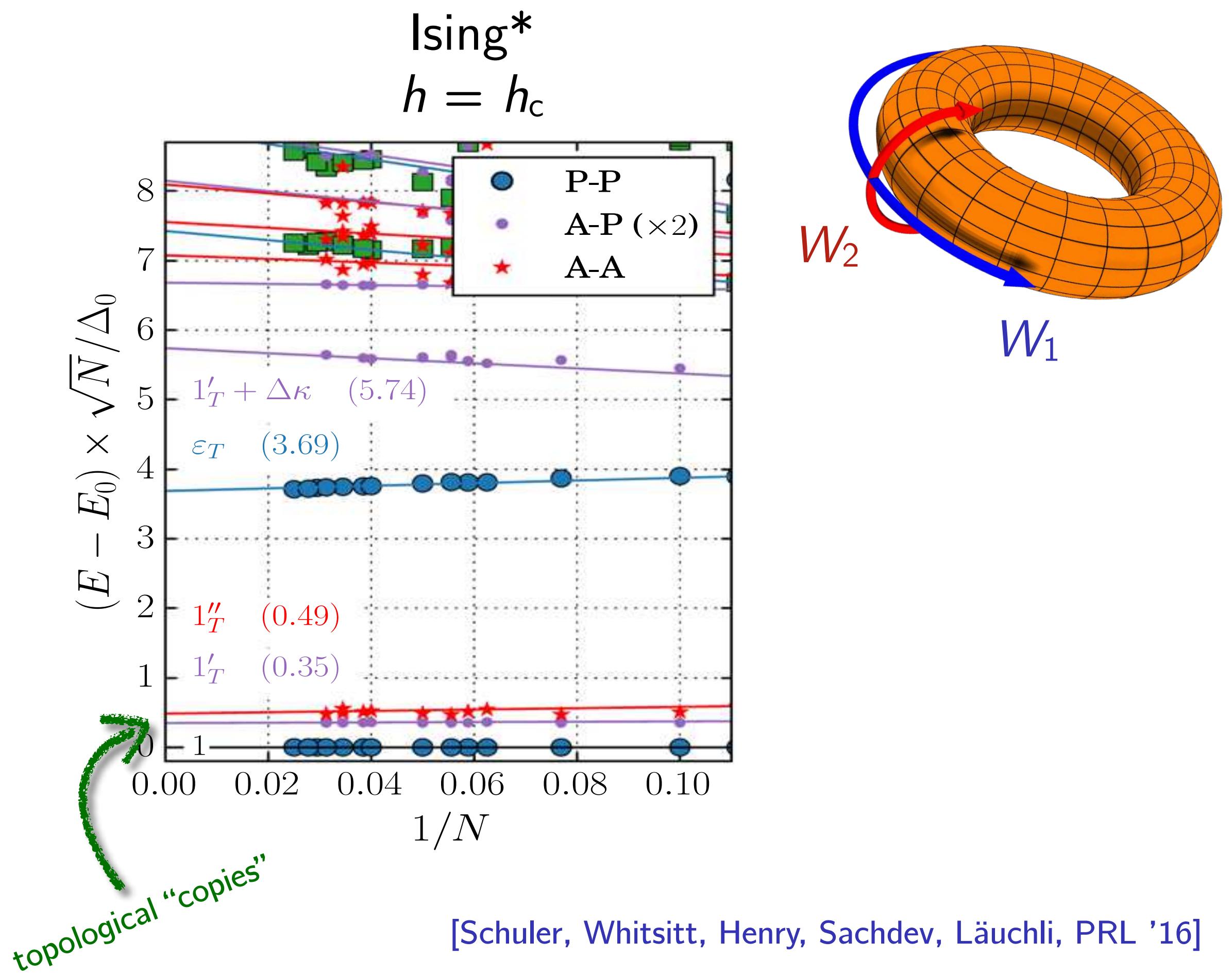
Transverse-field Ising:

$$H = -J \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z - h \sum_i \sigma_i^x$$

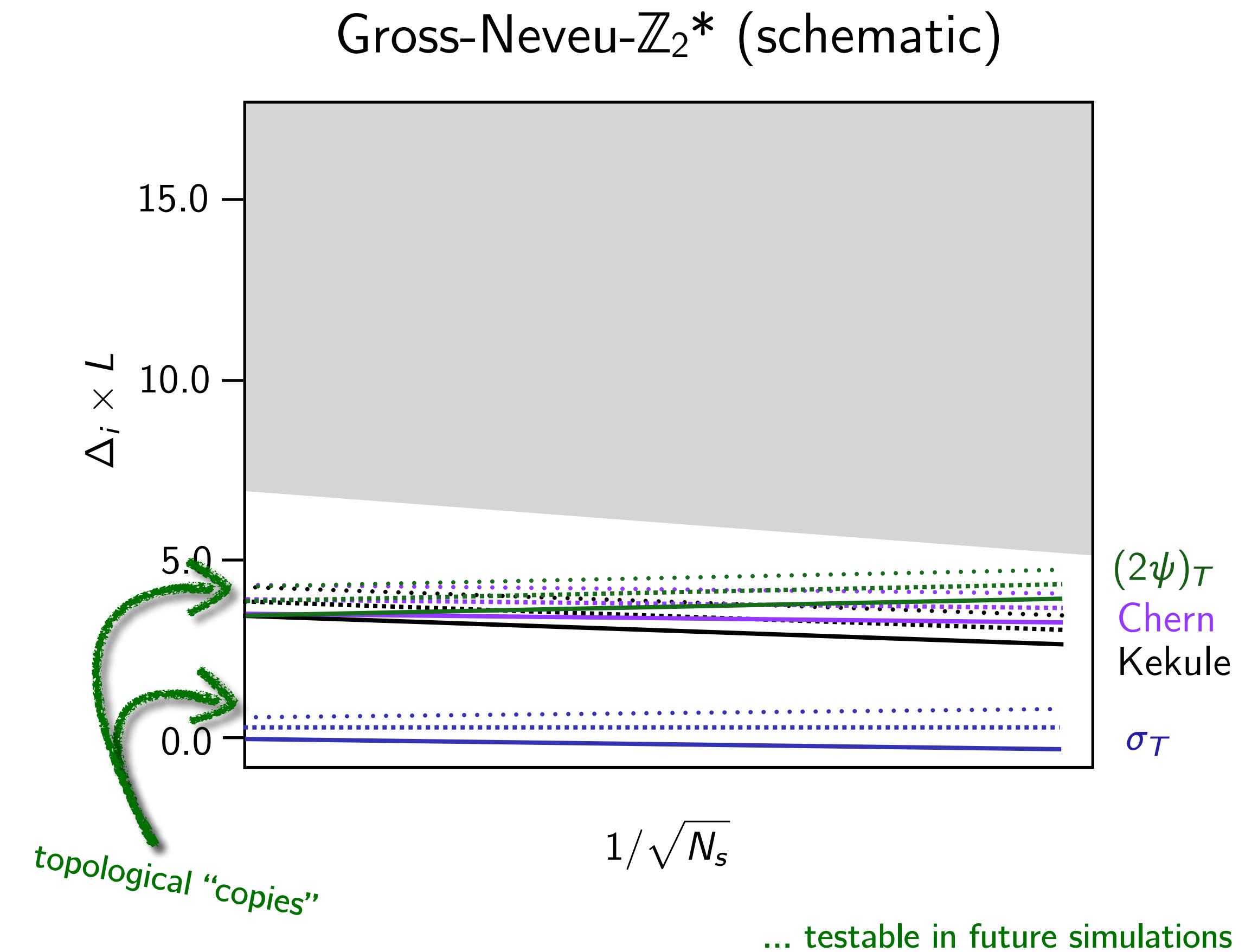
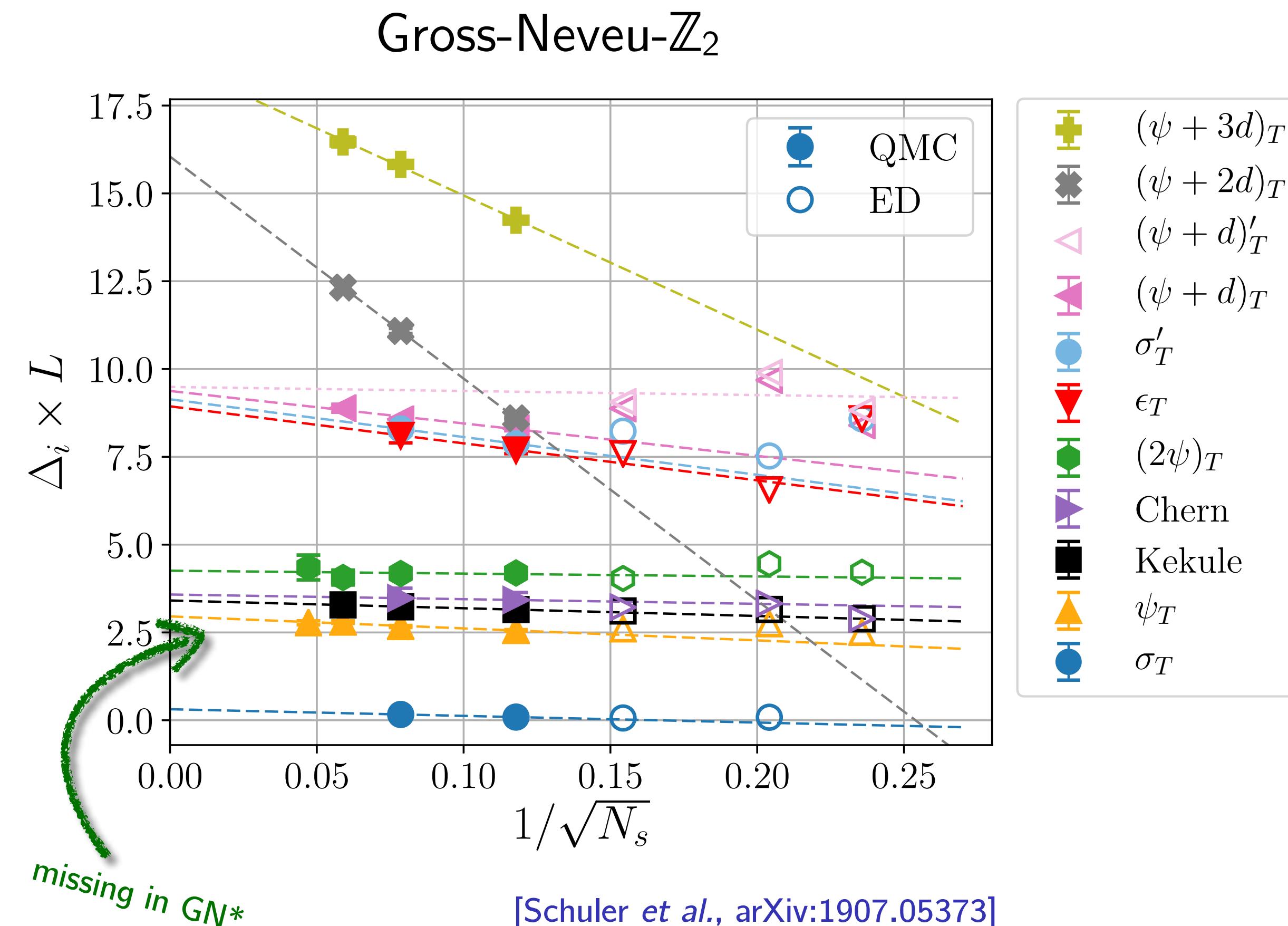


Transverse-field toric code:

$$H = -J \sum_s \prod_{i \in s} \sigma_i^x - J \sum_p \prod_{i \in p} \sigma_i^z - h \sum_i \sigma_i^x$$



# Gross-Neveu vs Gross-Neveu\*



# Outline

- (1) Introduction: “*Beyond-Landau*” quantum criticality
- (2) Spin-1/2: *Field-induced criticality in Kitaev materials*
- (3) Spin-3/2: *Fractionalized fermionic criticality in spin-orbital models*
- (4) Conclusions

# Conclusions

