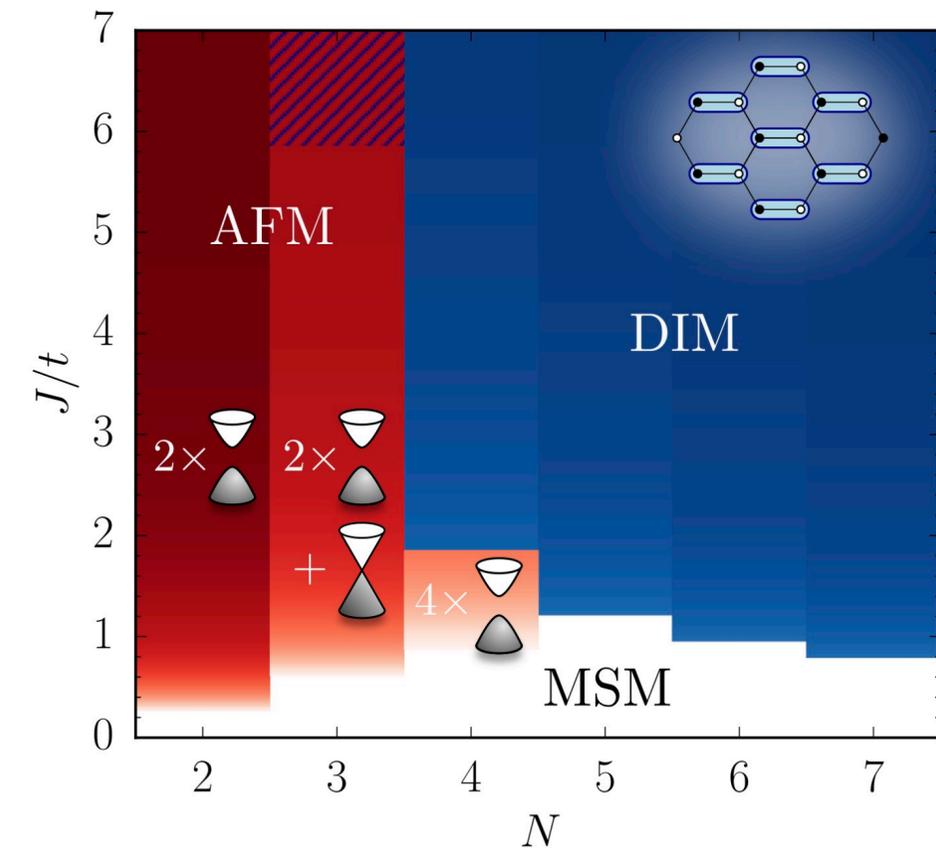
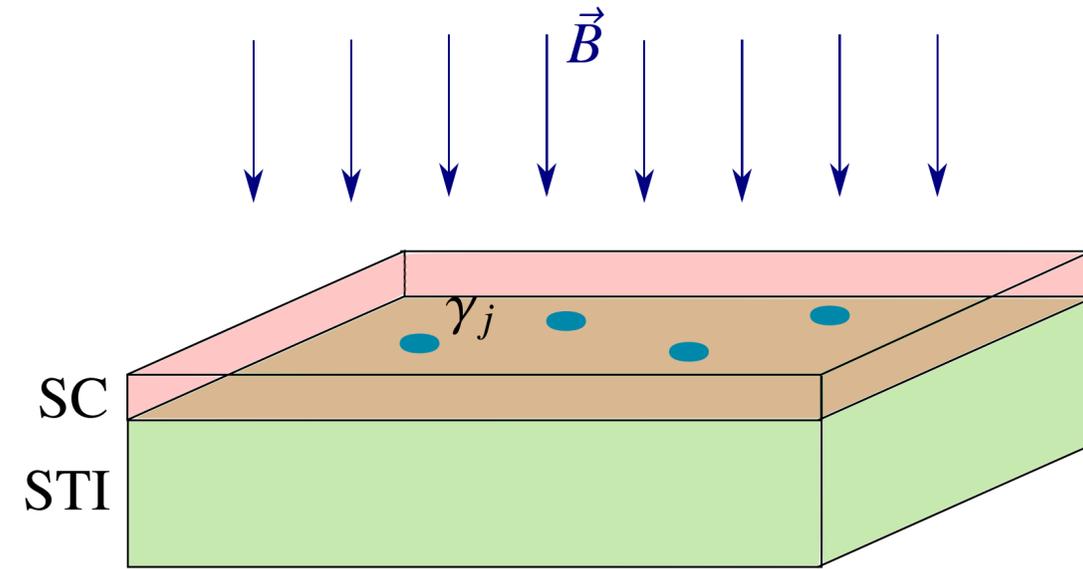


# Interacting Majorana fermions

Lukas Janssen (TUD) & Urban Seifert (UCSB)



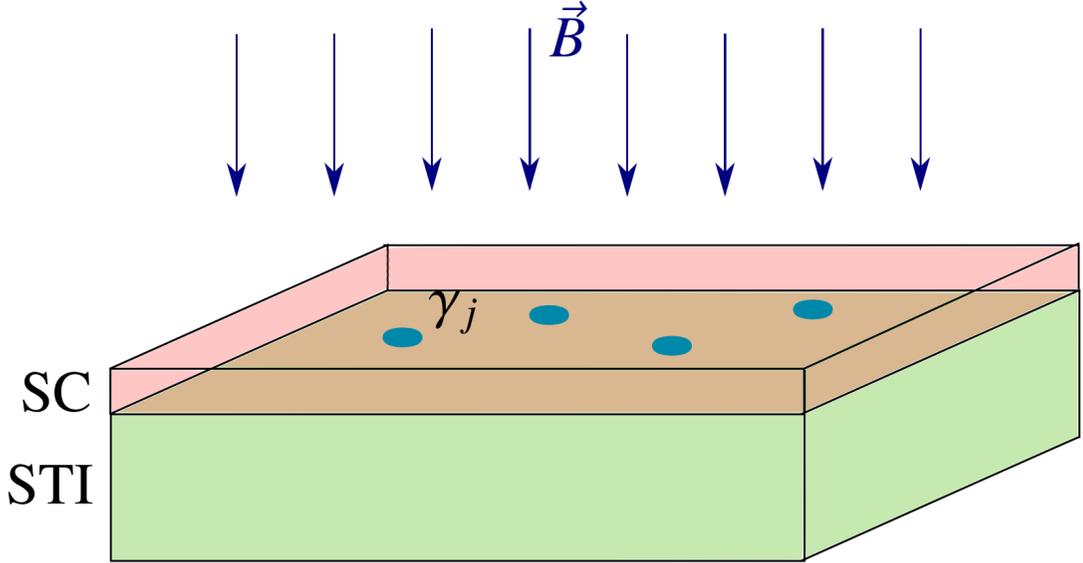
# Motivation #1: Superconductor / topological insulator heterostructures



[Fu & Kane, PRL '08]

[Rahmani & Franz, Rep. Progr. Phys. '19]

# Motivation #1: Superconductor / topological insulator heterostructures



[Fu & Kane, PRL '08]  
[Rahmani & Franz, Rep. Progr. Phys. '19]

Effective model:

tunable by gate voltage

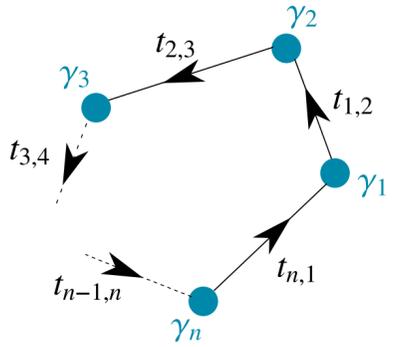
$$H = \sum_{ij} i t_{ij} \gamma_i \gamma_j + \sum_{ijkl} g_{ijkl} \gamma_i \gamma_j \gamma_k \gamma_l + \dots$$

“Majorana-Hubbard models”

Grosfeld-Stern rule:

$$\arg(i^n t_{1,2} \dots t_{n-1,n} t_{n,1}) = \frac{\pi}{2}(n - 2)$$

Square lattice:  $\pi$  flux  
Honeycomb lattice: 0 flux

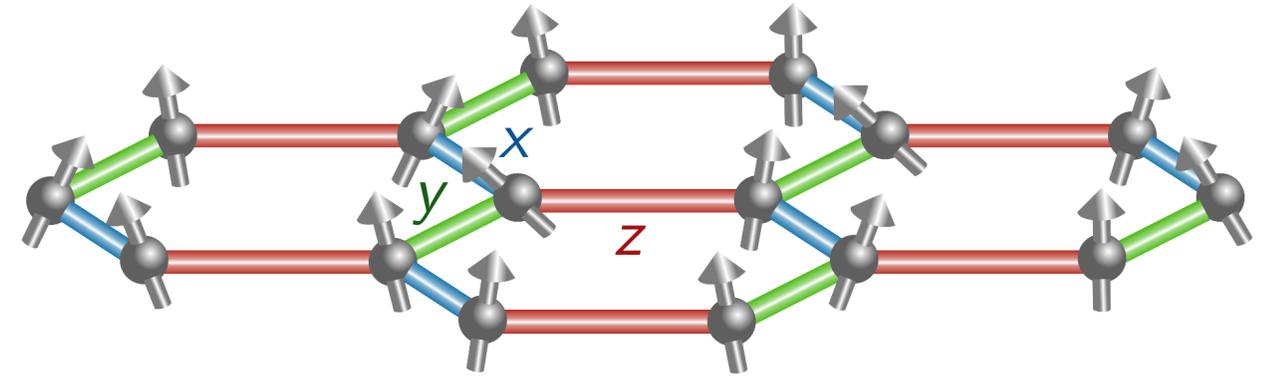


[Grosfeld & Stern, PRB '06]  
[Liu & Franz, PRB '15]  
[Wamer & Affleck, PRB '18]  
[Li & Franz, PRB '18]  
[Tummuru, Nocera, Affleck, PRB '21]

# Motivation #2: Kitaev honeycomb model

Hamiltonian:

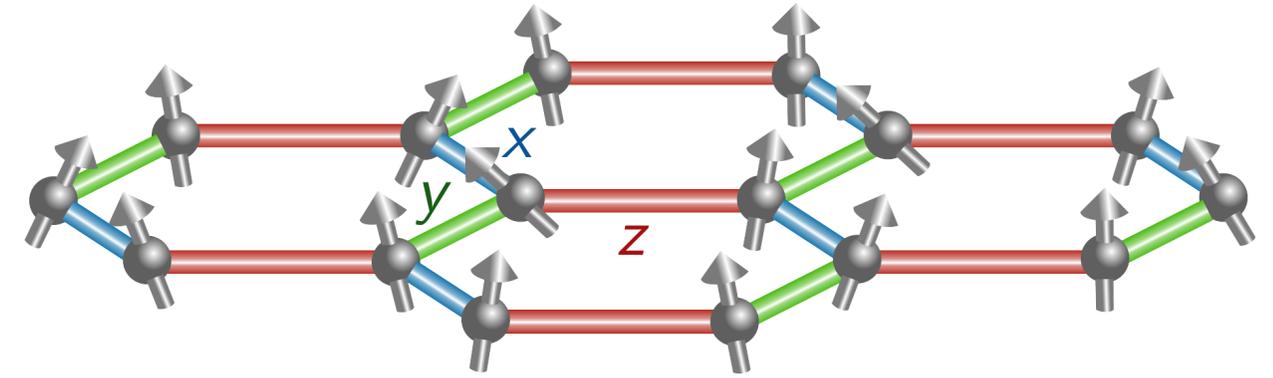
$$\mathcal{H} = K \left( \sum_{\langle ij \rangle_x} \sigma_i^x \sigma_j^x + \sum_{\langle ij \rangle_y} \sigma_i^y \sigma_j^y + \sum_{\langle ij \rangle_z} \sigma_i^z \sigma_j^z \right)$$



# Motivation #2: Kitaev honeycomb model

Hamiltonian:

$$\mathcal{H} = K \left( \sum_{\langle ij \rangle_x} \sigma_i^x \sigma_j^x + \sum_{\langle ij \rangle_y} \sigma_i^y \sigma_j^y + \sum_{\langle ij \rangle_z} \sigma_i^z \sigma_j^z \right)$$

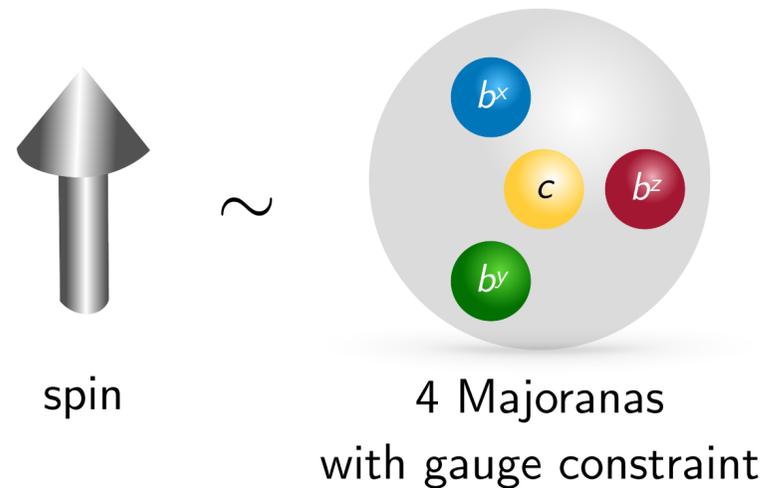


Majorana representation:

$$\sigma^x \sim i b^x c$$

$$\sigma^y \sim i b^y c$$

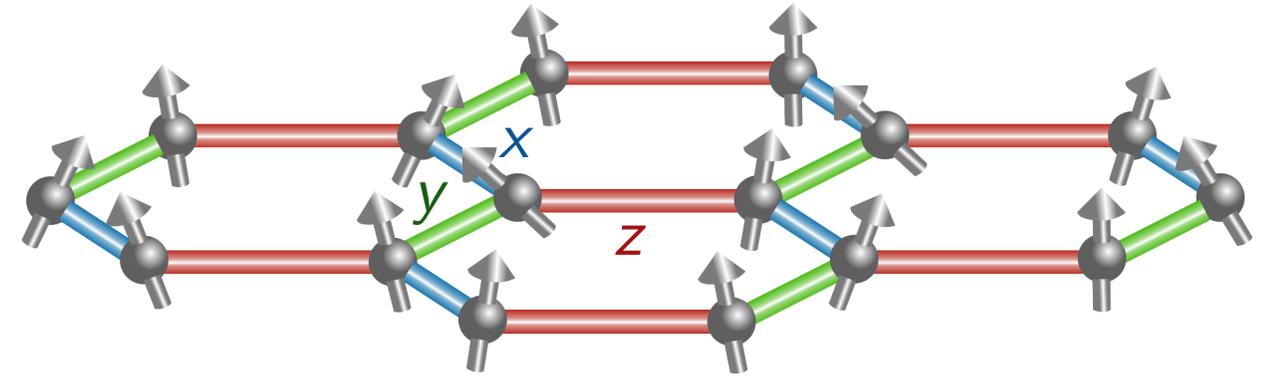
$$\sigma^z \sim i b^z c$$



# Motivation #2: Kitaev honeycomb model

Hamiltonian:

$$\mathcal{H} = K \left( \sum_{\langle ij \rangle_x} \sigma_i^x \sigma_j^x + \sum_{\langle ij \rangle_y} \sigma_i^y \sigma_j^y + \sum_{\langle ij \rangle_z} \sigma_i^z \sigma_j^z \right)$$

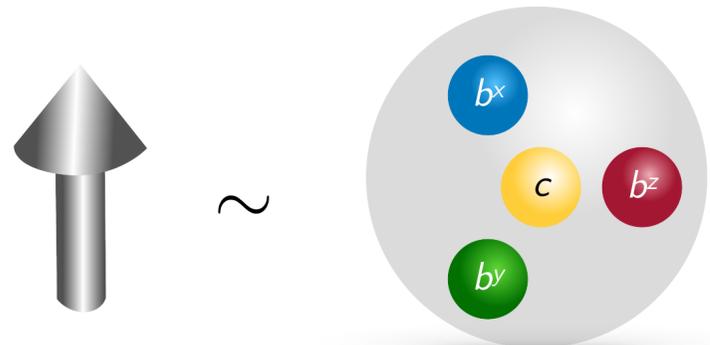


Majorana representation:

$$\sigma^x \sim i b^x c$$

$$\sigma^y \sim i b^y c$$

$$\sigma^z \sim i b^z c$$



spin

4 Majoranas  
with gauge constraint

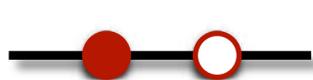
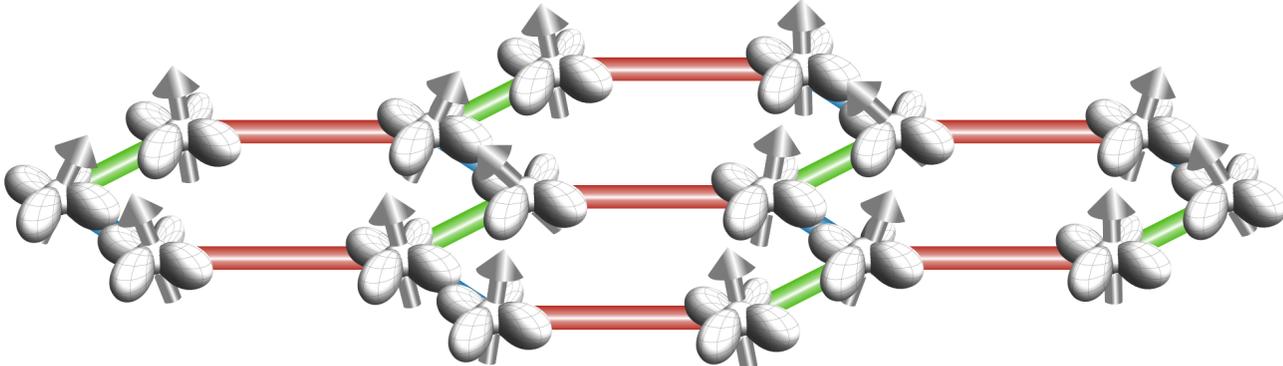
Fractionalization:

$$\mathcal{H} \sim iK \sum_{\langle ij \rangle_\alpha} \underbrace{(i b_i^\alpha b_j^\alpha)}_{\equiv \hat{u}_{ij} = \hat{u}_{ij}^\dagger} c_i c_j$$

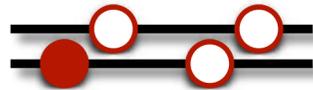
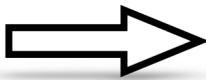
with  $[\hat{u}_{ij}, \tilde{\mathcal{H}}] = 0 \Rightarrow$  static  $\mathbb{Z}_2$  gauge field!

# Beyond Kitaev spin-1/2

Spin-orbital generalization:



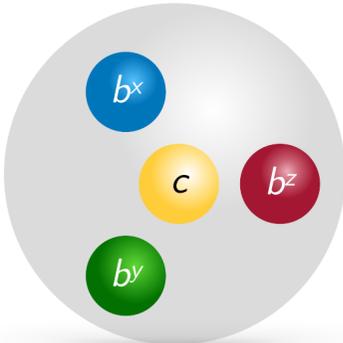
$$\sigma^\alpha \quad 2 \times 2$$



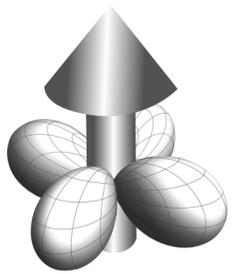
$$\gamma^i = \sigma^\alpha \otimes \tau^\beta \quad 4 \times 4$$



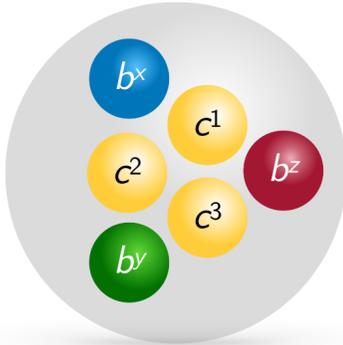
spin



4 Majoranas  
with gauge constraint



spin + orbital



6 Majoranas  
with gauge constraint

[Chulliparambil, *et int.*, LJ, Tu, PRB '20]

# SO(3) Kitaev-Heisenberg spin-orbital model

Hamiltonian:

$$\mathcal{H} = K \sum_{\langle ij \rangle_\alpha} \underbrace{\vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \tau_i^\alpha \tau_j^\alpha}_{\mapsto \hat{u}_{ij} c_i^\top c_j} + J \sum_{\langle ij \rangle} \underbrace{\vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \mathbb{1}_i \mathbb{1}_j}_{\mapsto \frac{1}{4} (c_i^\top \vec{L} c_i) \cdot (c_j^\top \vec{L} c_j)}$$

spin-1 matrices 

with  $[\hat{u}_{ij}, \mathcal{H}] = 0$  still static!

# SO(3) Kitaev-Heisenberg spin-orbital model

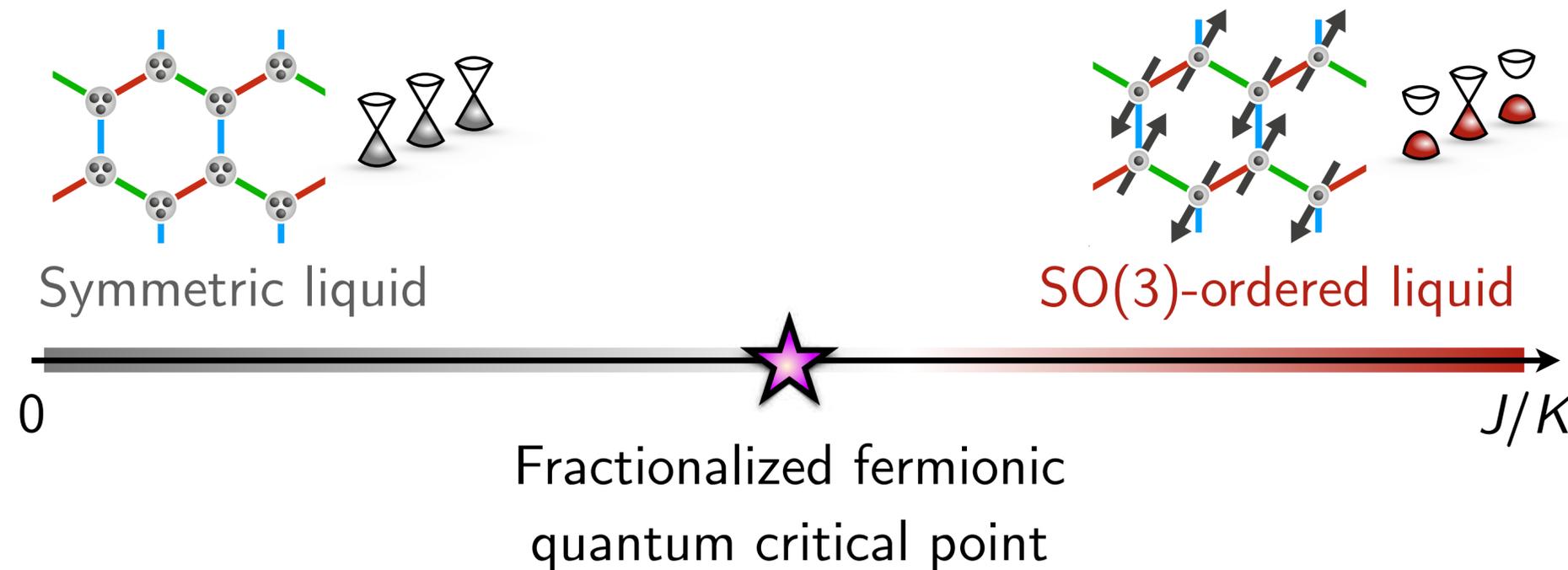
Hamiltonian:

$$\mathcal{H} = K \sum_{\langle ij \rangle_\alpha} \underbrace{\vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \tau_i^\alpha \tau_j^\alpha}_{\mapsto \hat{u}_{ij} c_i^\top c_j} + J \sum_{\langle ij \rangle} \underbrace{\vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \mathbb{1}_i \mathbb{1}_j}_{\mapsto \frac{1}{4} (c_i^\top \vec{L} c_i) \cdot (c_j^\top \vec{L} c_j)}$$

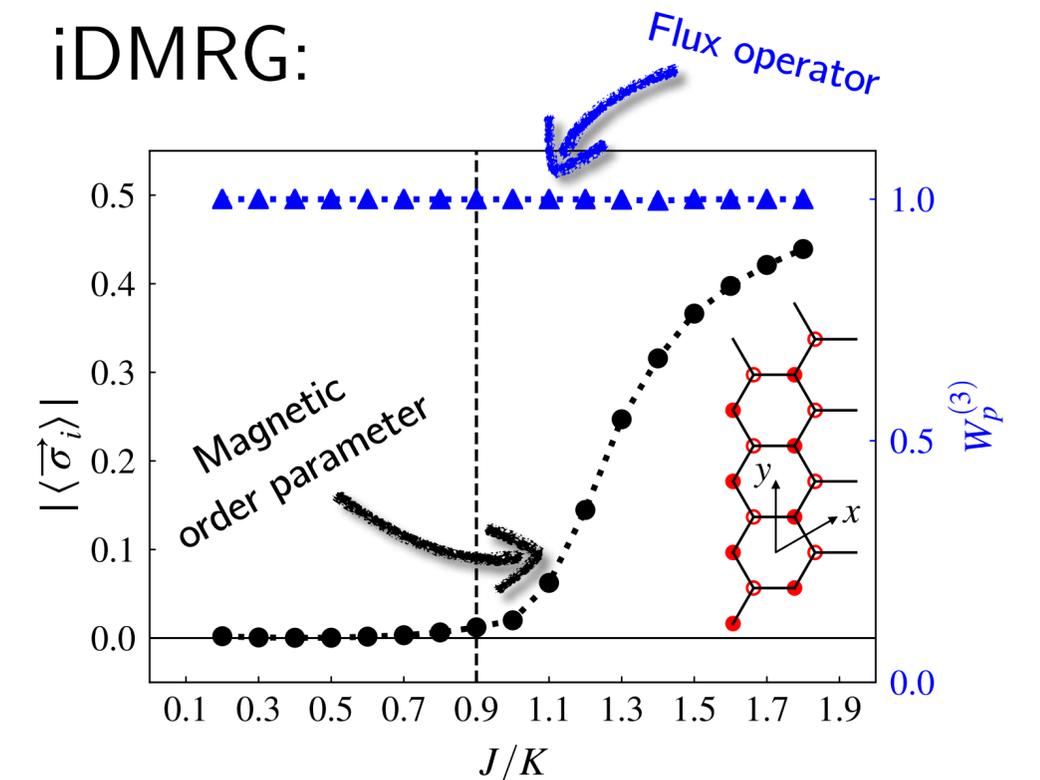
spin-1 matrices

with  $[\hat{u}_{ij}, \mathcal{H}] = 0$  still static!

Phase diagram:



iDMRG:



[Seifert, Dong, et al., LJ, PRL '20]

# SO(N) Majorana-Hubbard models

Hamiltonian:

$$\mathcal{H} = \sum_{\langle ij \rangle} i t_{ij} c_i^\top c_j + J \sum_{\langle ij \rangle} \underbrace{\sum_{a < b} \left( \frac{1}{2} c_i^\top L^{ab} c_i \right) \left( \frac{1}{2} c_j^\top L^{ab} c_j \right)}_{= \frac{1}{2} (c_i^\top c_j)(c_i^\top c_j)}$$

Generators of SO(N)  


with  $c_i \equiv (c_i^1, \dots, c_i^N)^\top$

*“Majorana analog of SU(N)  
Hubbard-Heisenberg model”*

# SO(N) Majorana-Hubbard models

Hamiltonian:

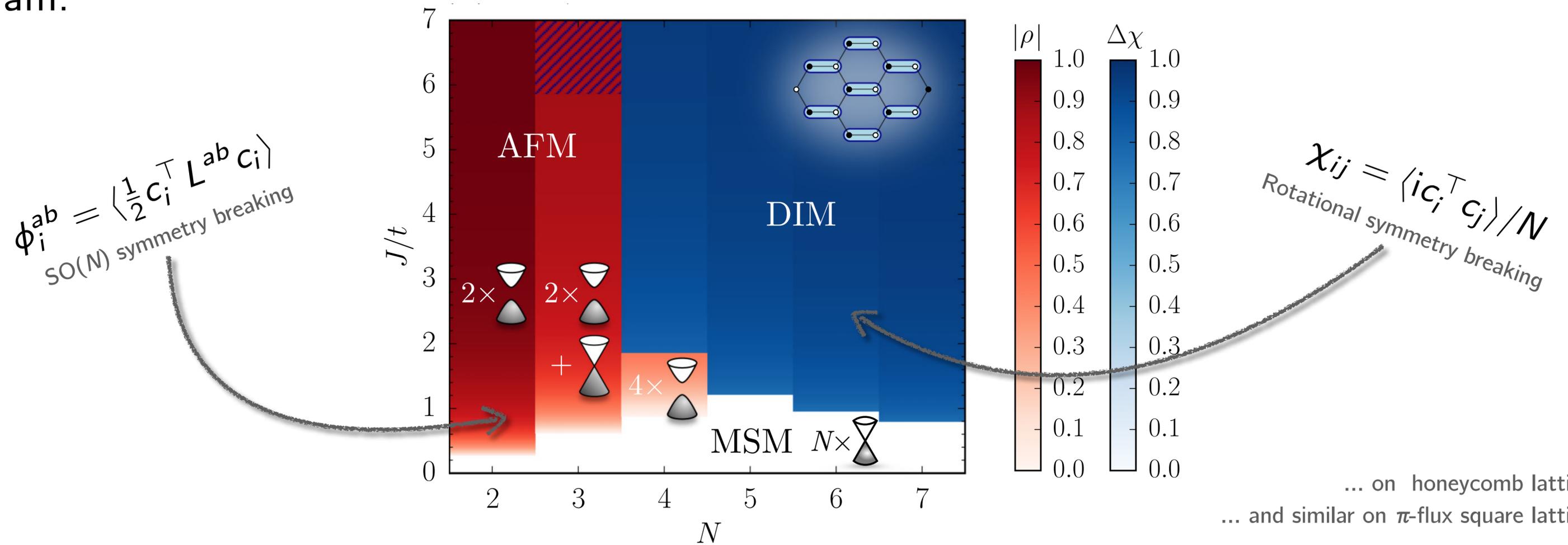
$$\mathcal{H} = \sum_{\langle ij \rangle} it_{ij} c_i^\top c_j + J \underbrace{\sum_{\langle ij \rangle} \sum_{a < b} \left( \frac{1}{2} c_i^\top L^{ab} c_i \right) \left( \frac{1}{2} c_j^\top L^{ab} c_j \right)}_{= \frac{1}{2} (c_i^\top c_j) (c_i^\top c_j)}$$

Generators of SO(N)

with  $c_i \equiv (c_i^1, \dots, c_i^N)^\top$

“Majorana analog of SU(N) Hubbard-Heisenberg model”

Phase diagram:



... on honeycomb lattice  
... and similar on  $\pi$ -flux square lattice

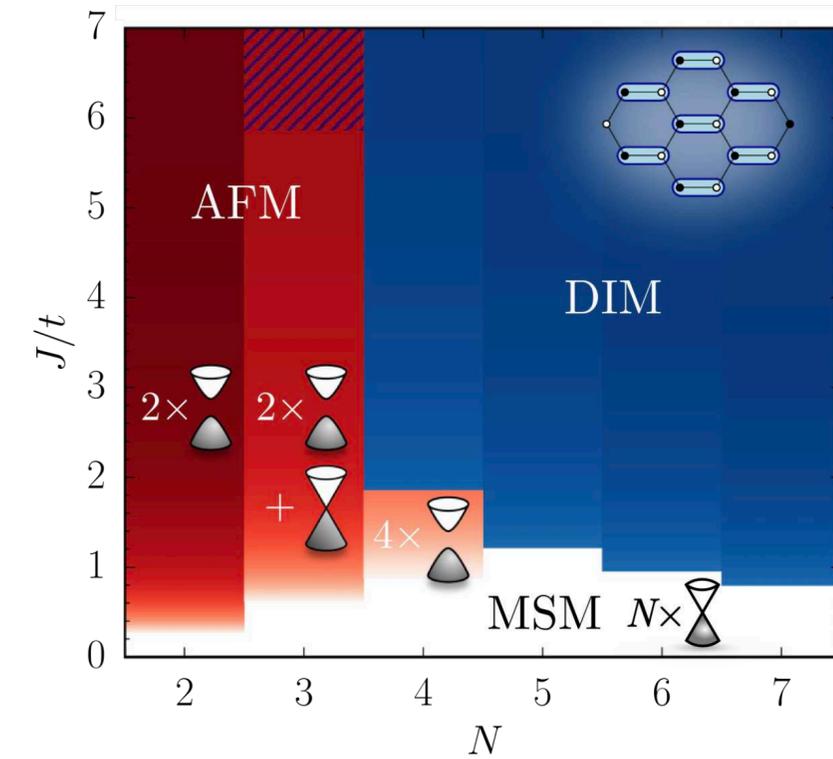
# Quantum phase transitions

Effective model (MSM-to-AFM transition):

$$\mathcal{L} = \bar{\psi}_\alpha \gamma_\mu \partial_\mu \psi_\alpha + \frac{1}{4} \phi^{ab} (r - \partial_\mu^2) \phi^{ab} + \frac{g}{2} \phi^{ab} \bar{\psi}_\alpha (L^{ab})_{\alpha\beta} \psi_\beta$$

$$+ \frac{\lambda_1}{4} (\phi^{ab} \phi^{ab})^2 + \lambda_2 \phi^{ab} \phi^{bc} \phi^{cd} \phi^{da}$$

tensor order parameter



# Quantum phase transitions

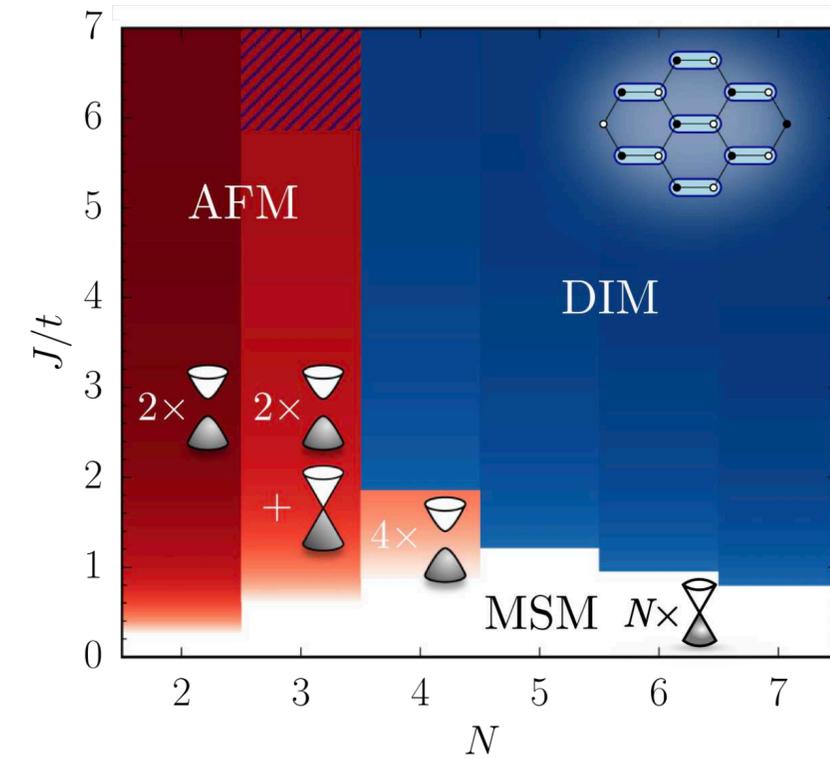
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tensor order parameter

$N \leq 3$ :

$$\frac{\lambda_1}{4} (\phi^{ab} \phi^{ab})^2 + \lambda_2 \phi^{ab} \phi^{bc} \phi^{cd} \phi^{da} = \frac{\lambda_1 + 2\lambda_2}{2} [\text{Tr}(\phi^2)]^2$$



# Quantum phase transitions

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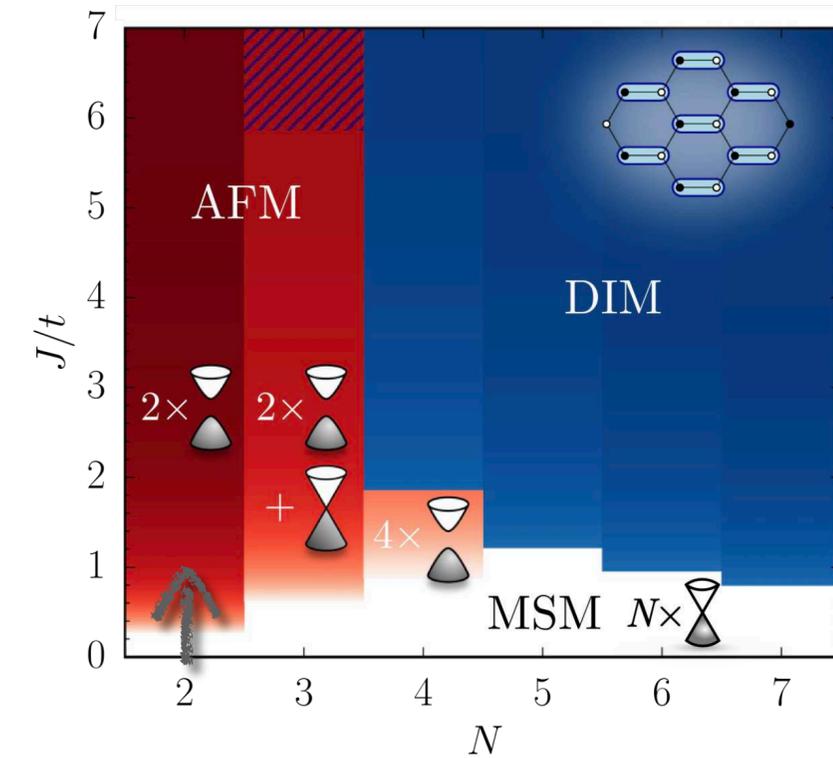
tensor order parameter

$N \leq 3$ :

$$\frac{\lambda_1}{4} (\phi^{ab} \phi^{ab})^2 + \lambda_2 \phi^{ab} \phi^{bc} \phi^{cd} \phi^{da} = \frac{\lambda_1 + 2\lambda_2}{2} [\text{Tr}(\phi^2)]^2$$

$N = 2$ :

$$(\phi^{ab}) = \begin{pmatrix} 0 & \phi \\ -\phi & 0 \end{pmatrix}$$



“Gross-Neveu-Ising”

[Wang, Corboz, Troyer, NJP '14]

[Li, Jiang, Yao, NJP '15]

[Huffman & Chandrasekharan, PRD '17; PRD '20]

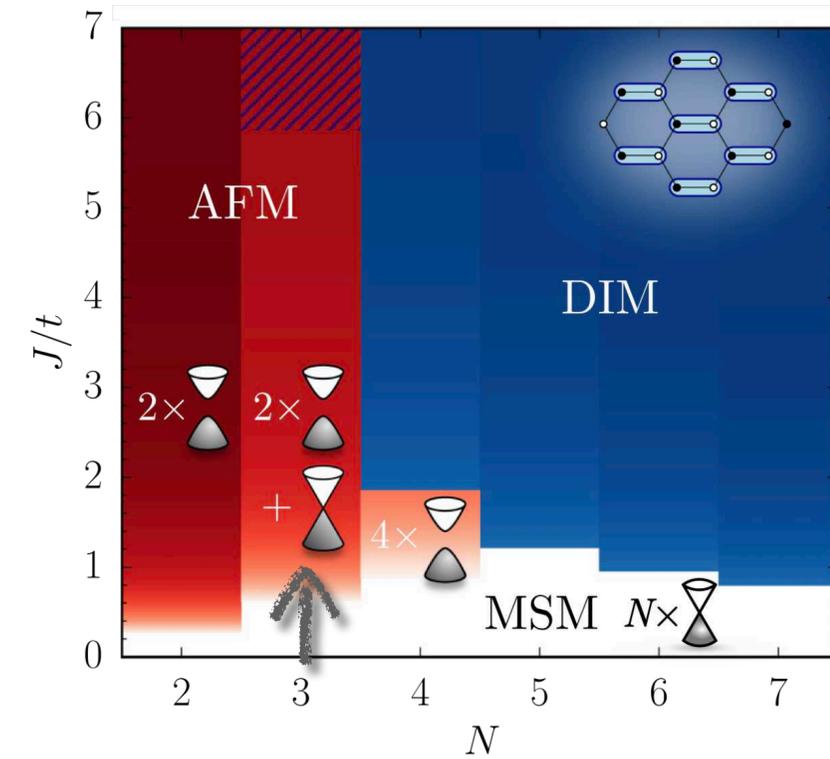
...

# Quantum phase transitions

Effective model (MSM-to-AFM transition):

$$\mathcal{L} = \bar{\psi}_\alpha \gamma_\mu \partial_\mu \psi_\alpha + \frac{1}{4} \phi^{ab} (r - \partial_\mu^2) \phi^{ab} + \frac{g}{2} \phi^{ab} \bar{\psi}_\alpha (L^{ab})_{\alpha\beta} \psi_\beta + \frac{\lambda_1}{4} (\phi^{ab} \phi^{ab})^2 + \lambda_2 \phi^{ab} \phi^{bc} \phi^{cd} \phi^{da}$$

tensor order parameter



$$N \leq 3: \quad \frac{\lambda_1}{4} (\phi^{ab} \phi^{ab})^2 + \lambda_2 \phi^{ab} \phi^{bc} \phi^{cd} \phi^{da} = \frac{\lambda_1 + 2\lambda_2}{2} [\text{Tr}(\phi^2)]^2$$

$$N = 2: \quad (\phi^{ab}) = \begin{pmatrix} 0 & \phi \\ -\phi & 0 \end{pmatrix}$$

$$N = 3: \quad (\phi^{ab}) = \begin{pmatrix} 0 & \phi_3 & \phi_2 \\ -\phi_3 & 0 & \phi_1 \\ -\phi_2 & -\phi_1 & 0 \end{pmatrix}$$

“Gross-Neveu-Ising”

[Wang, Corboz, Troyer, NJP '14]

[Li, Jiang, Yao, NJP '15]

[Huffman & Chandrasekharan, PRD '17; PRD '20]

...

“Gross-Neveu-SO(3)”

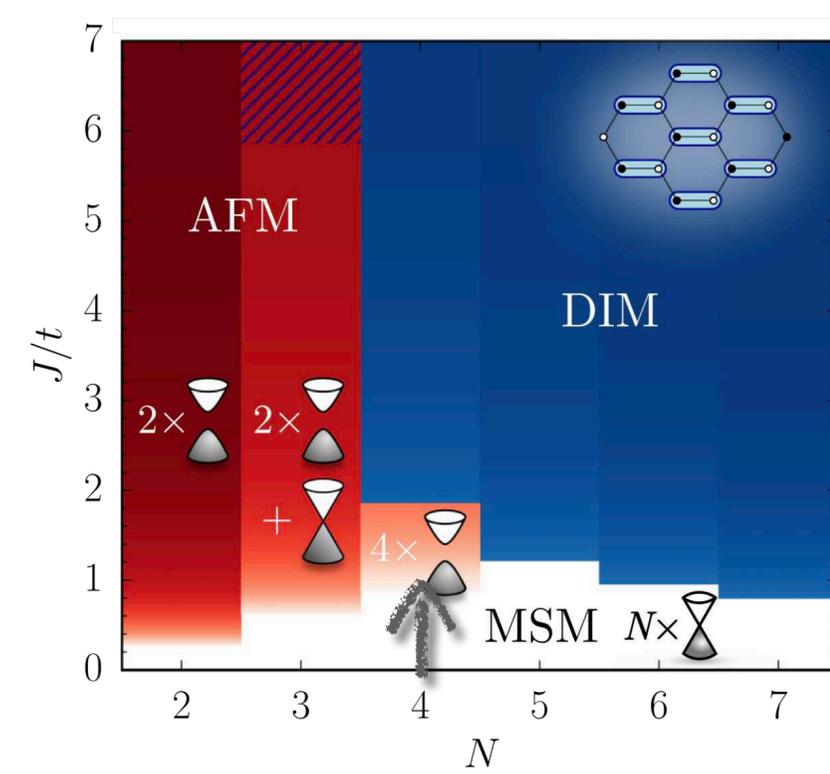
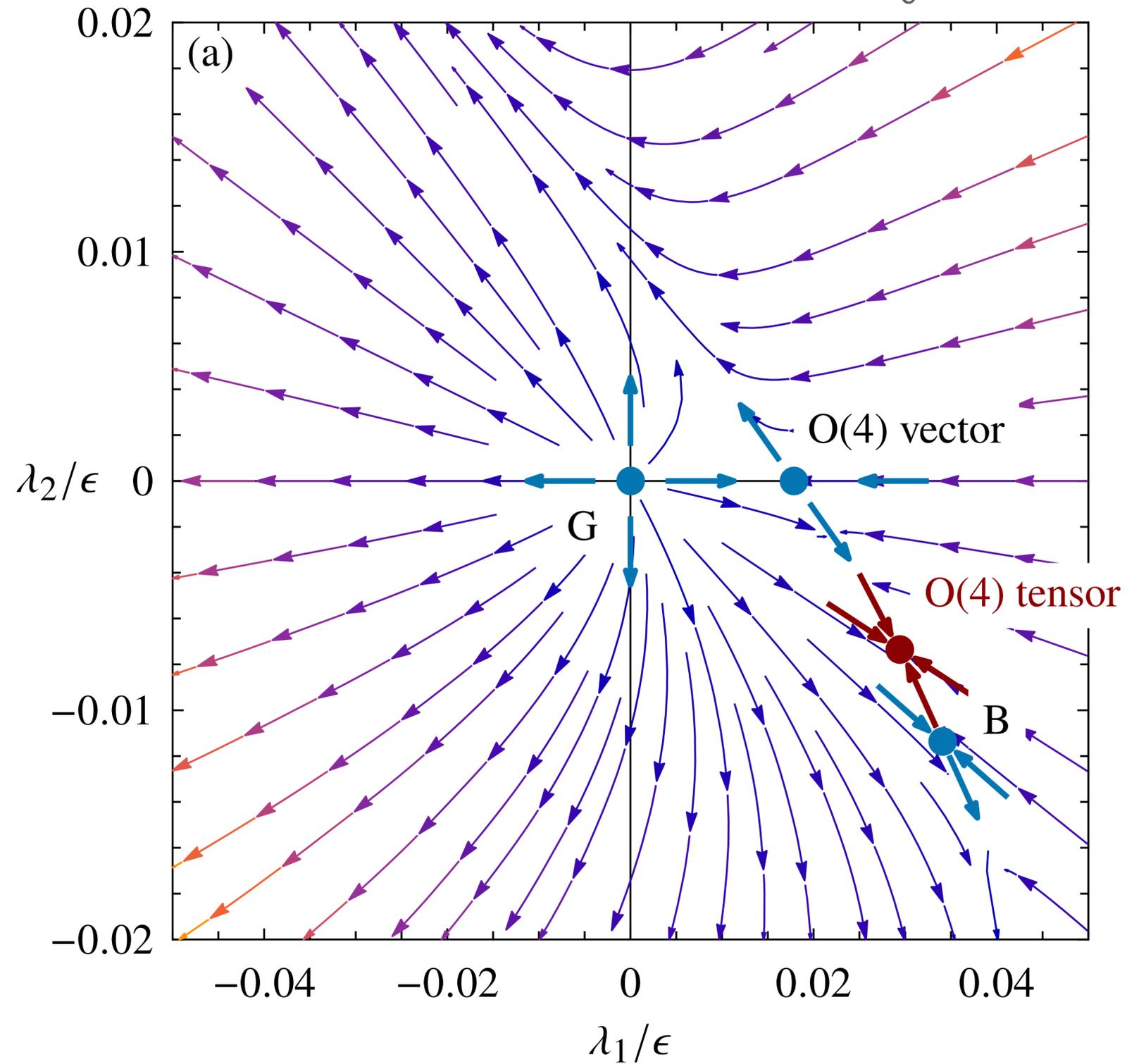
[Seifert, Dong, Chulliparambil, Vojta, Tu, LJ, PRL '20]

[Ray, Ihrig, Gracey, Scherer, LJ, PRB '21]

[Liu, Vojta, Assaad, LJ, PRL '22 (Editors' Suggestion)]

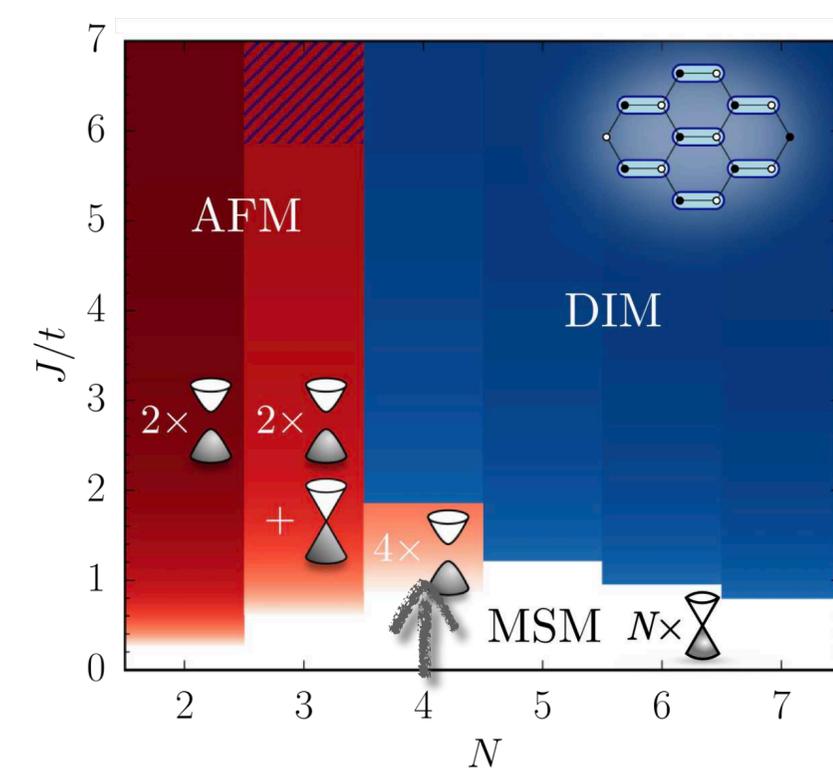
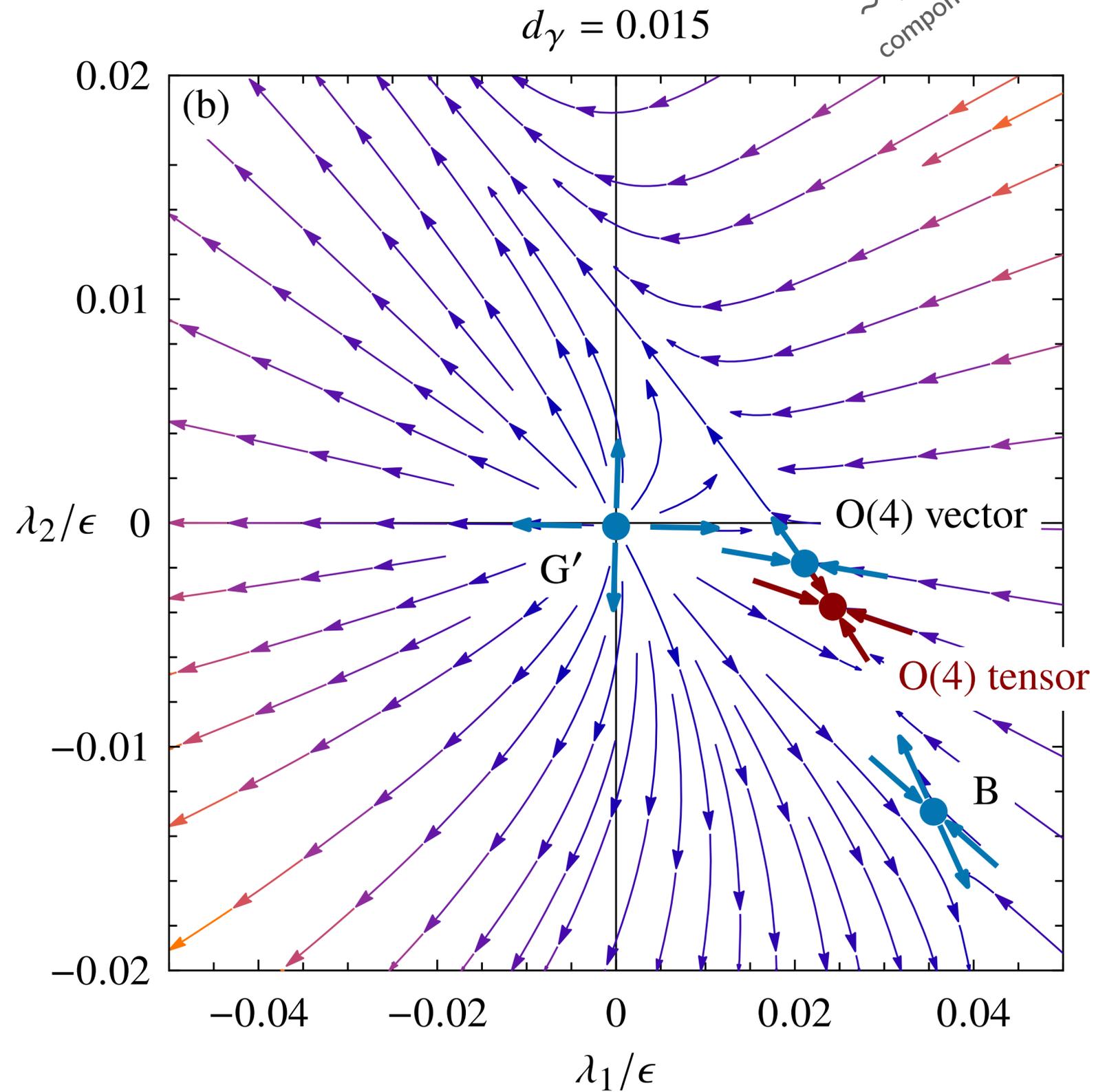
# Renormalization group flow: $N = 4$

Bosons only:



# Renormalization group flow: $N = 4$

Bosons  
+ "few" fermions:

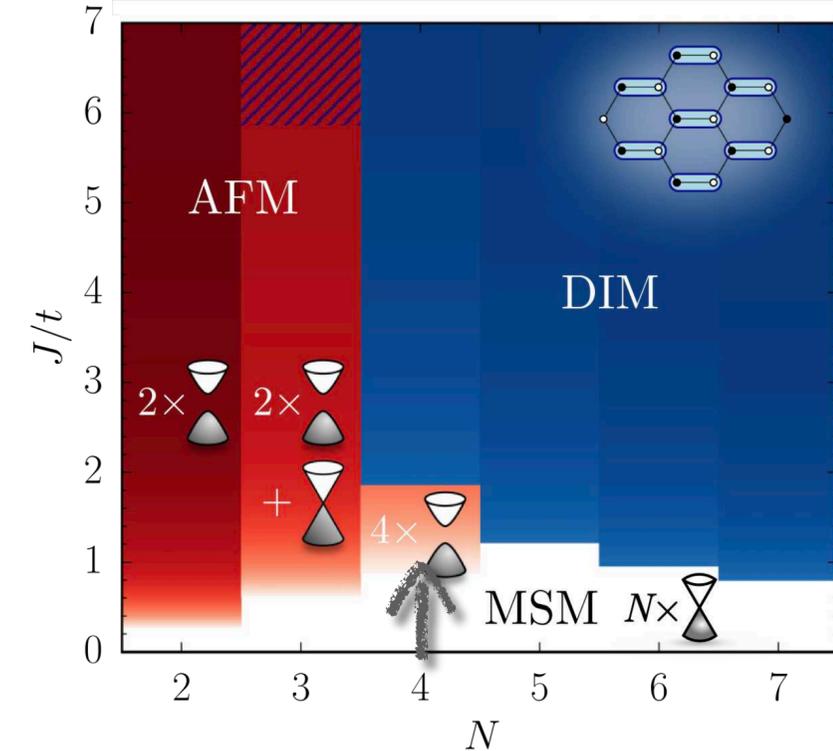
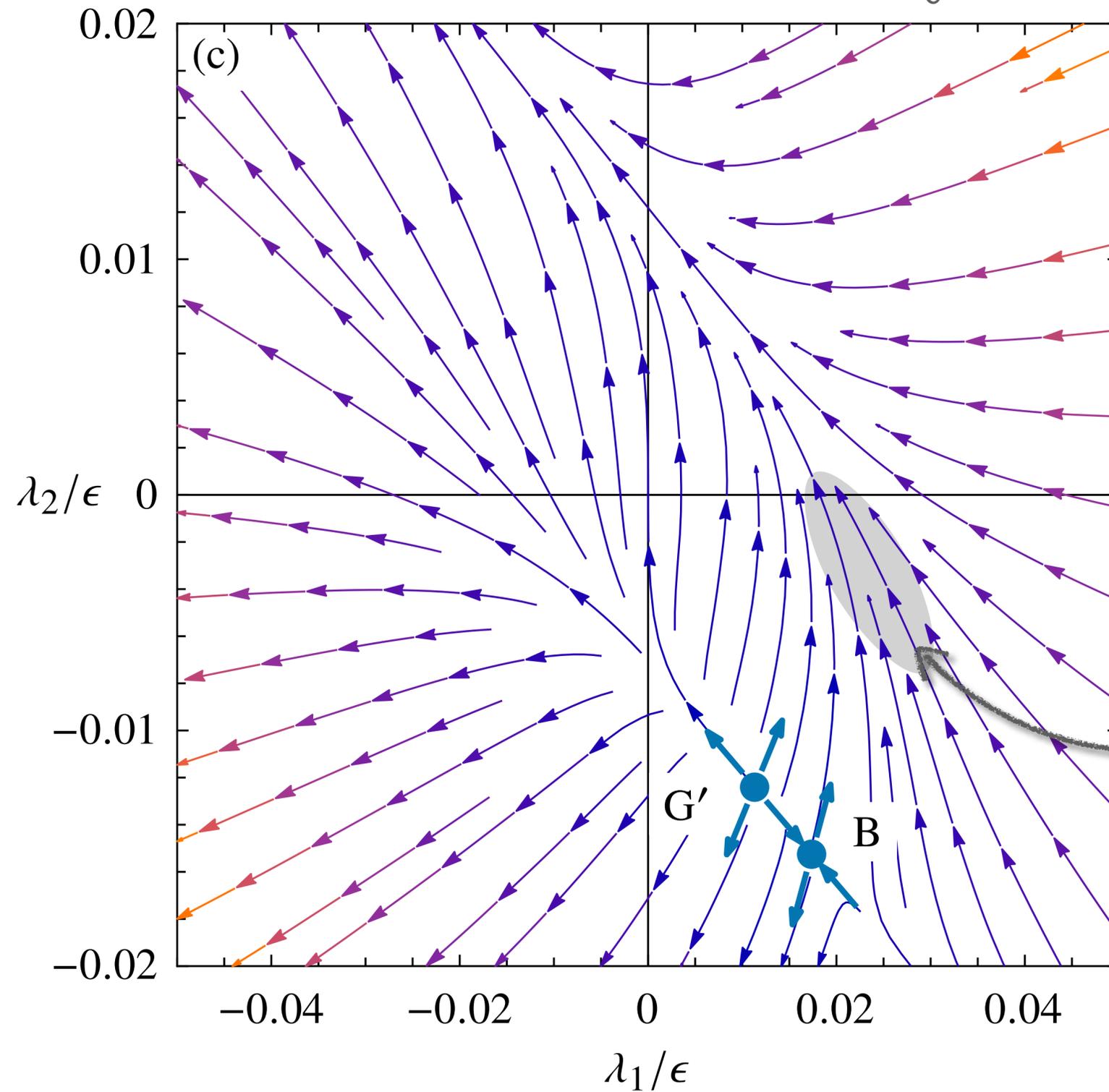


# Renormalization group flow: $N = 4$

Bosons  
+ "few more" fermions:

$$d_\gamma = 1.450$$

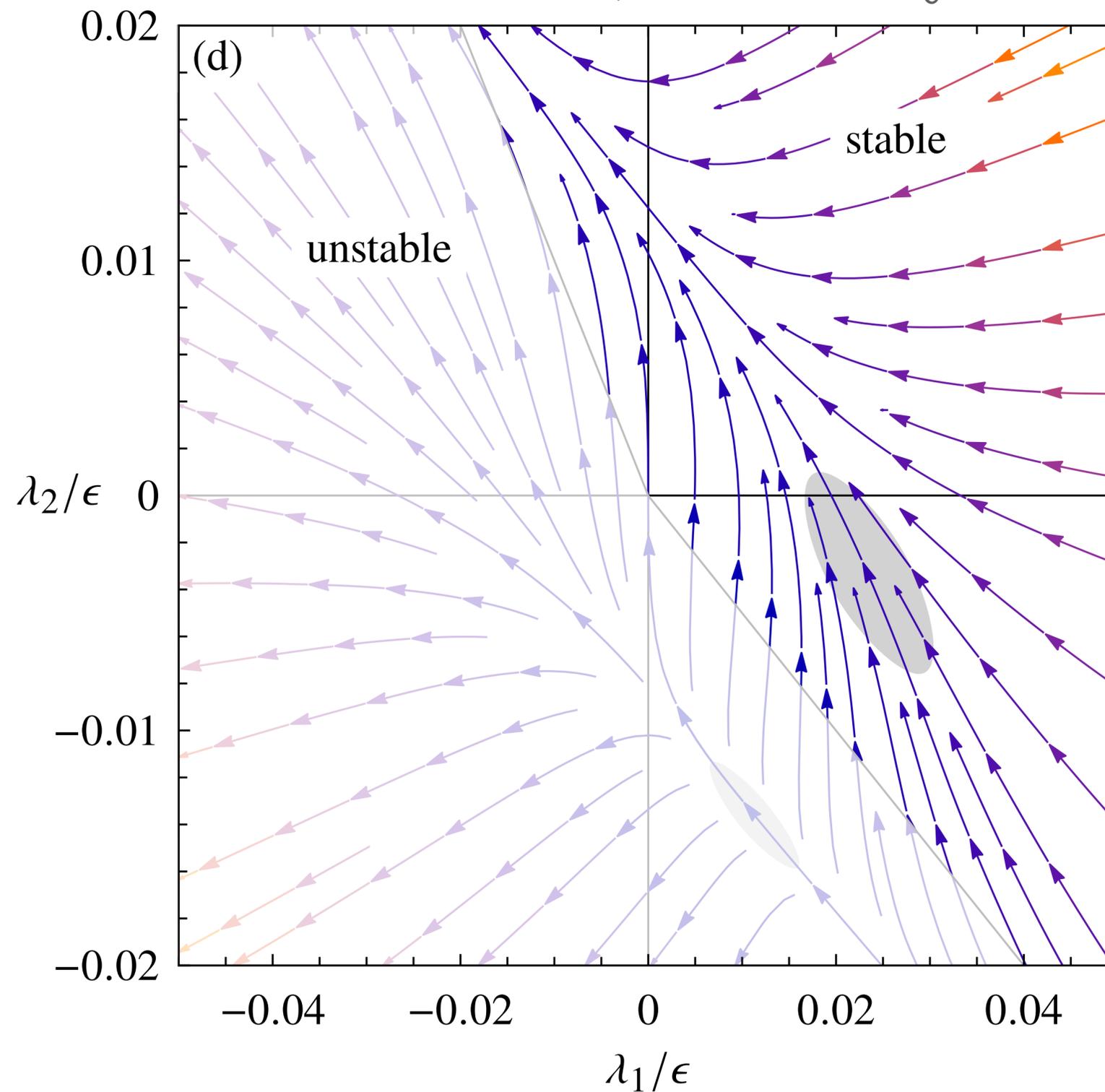
$\sim$  #(spinor components)



Fixed-point annihilation!

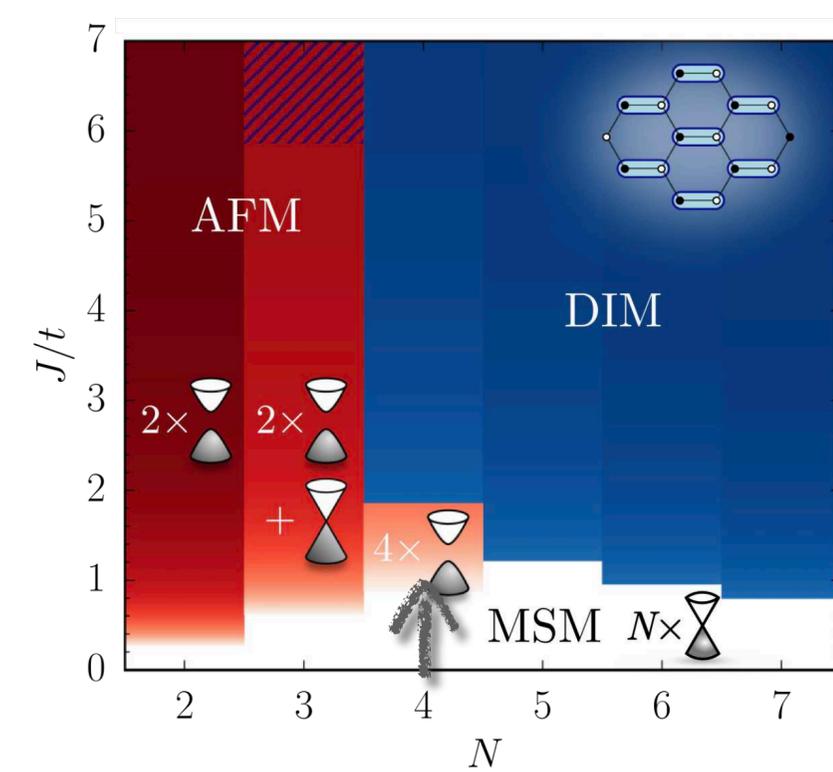
# Renormalization group flow: $N = 4$

Bosons  
+ "all" fermions:



$$d_\gamma = 2$$

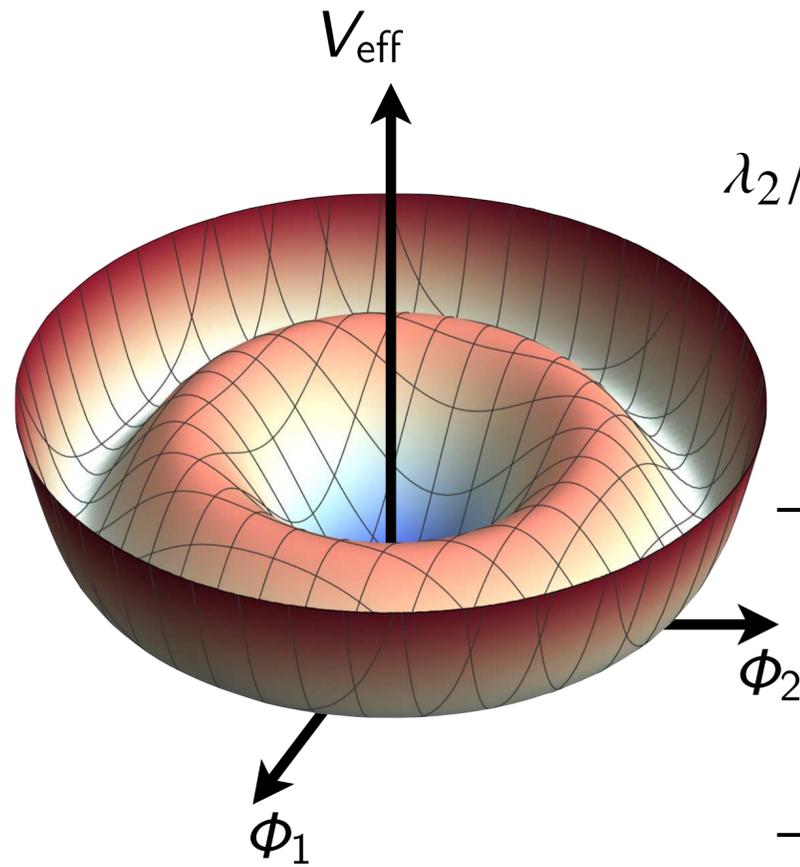
$\sim$  #(spinor components)



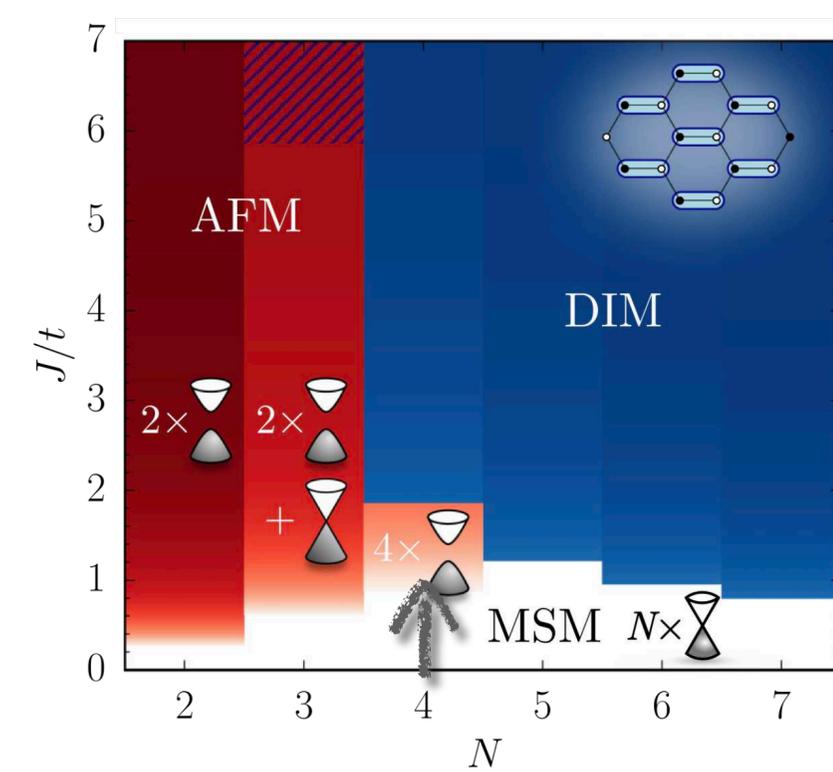
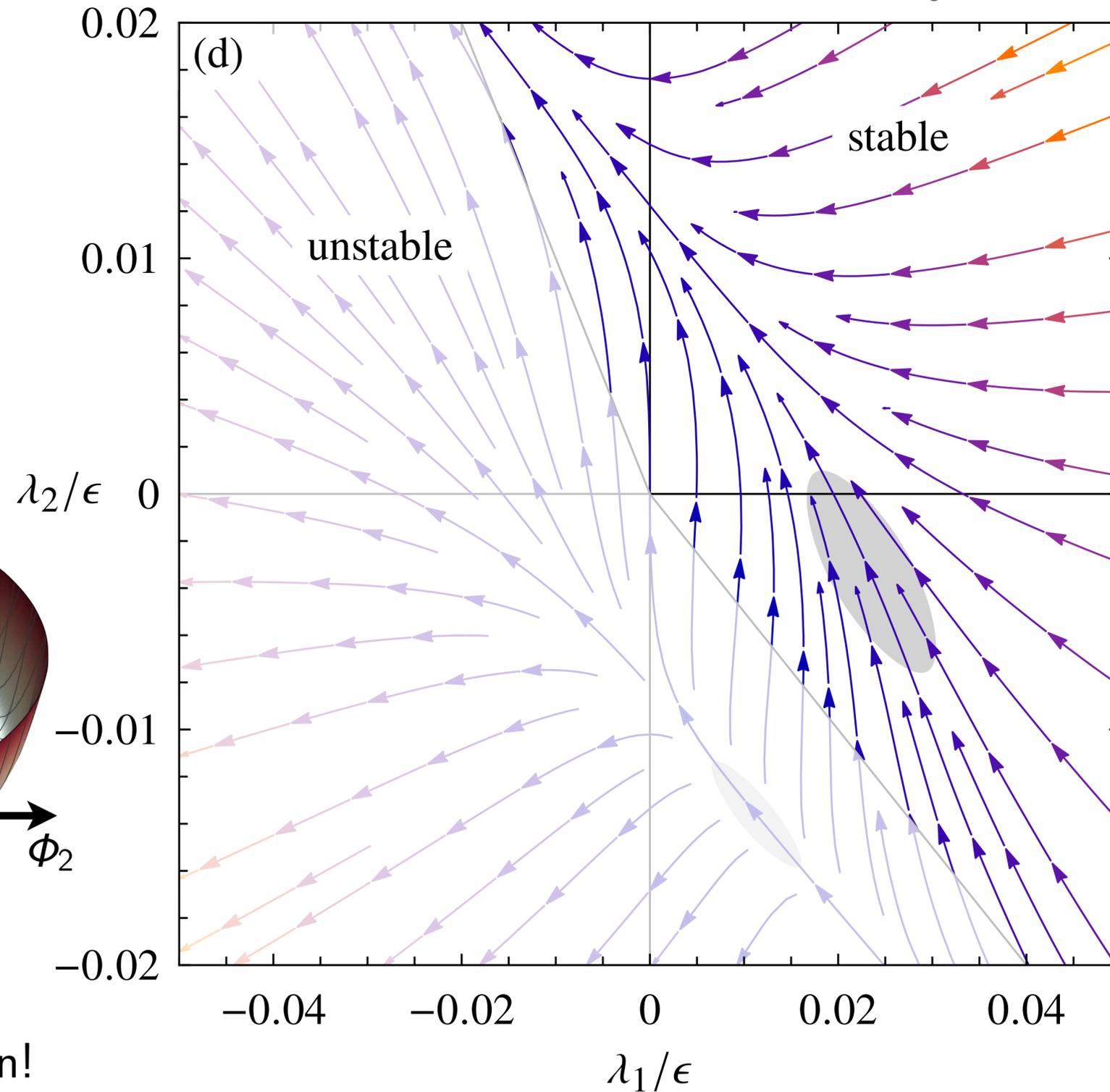
Only runaway flow!

# Renormalization group flow: $N = 4$

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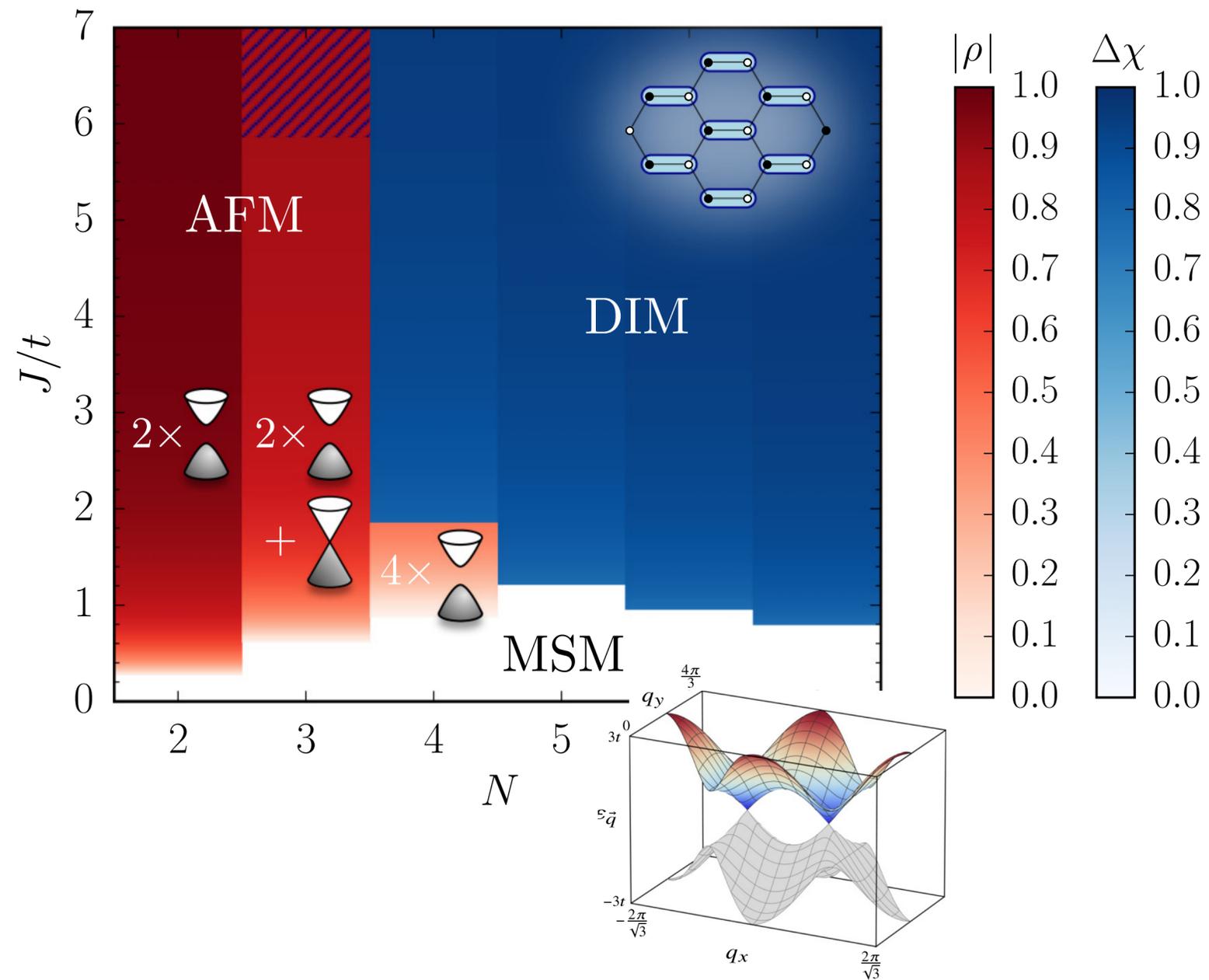
Weak first-order transition!



Only runaway flow!

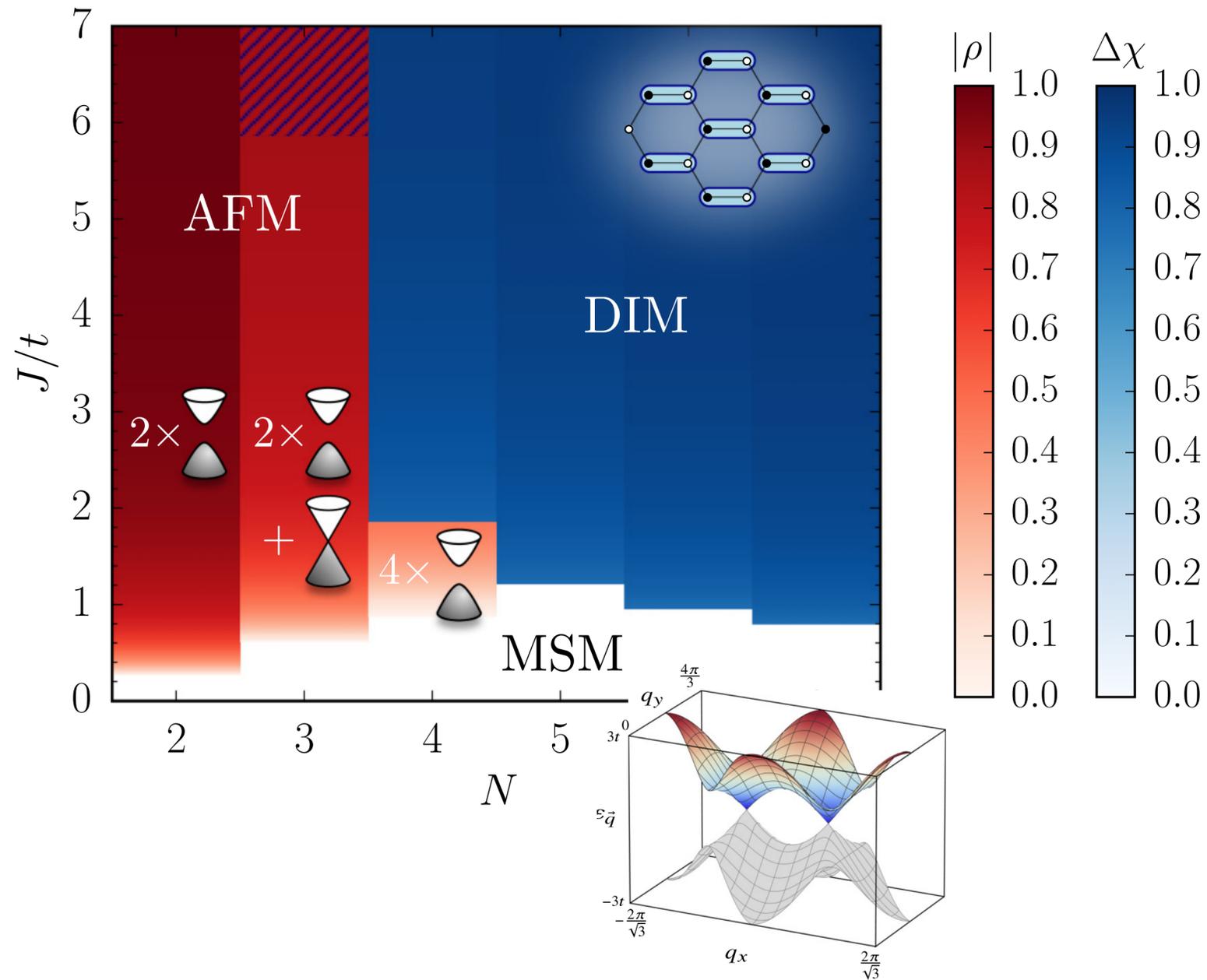
# Conclusions

SO( $N$ ) Majorana-Hubbard models



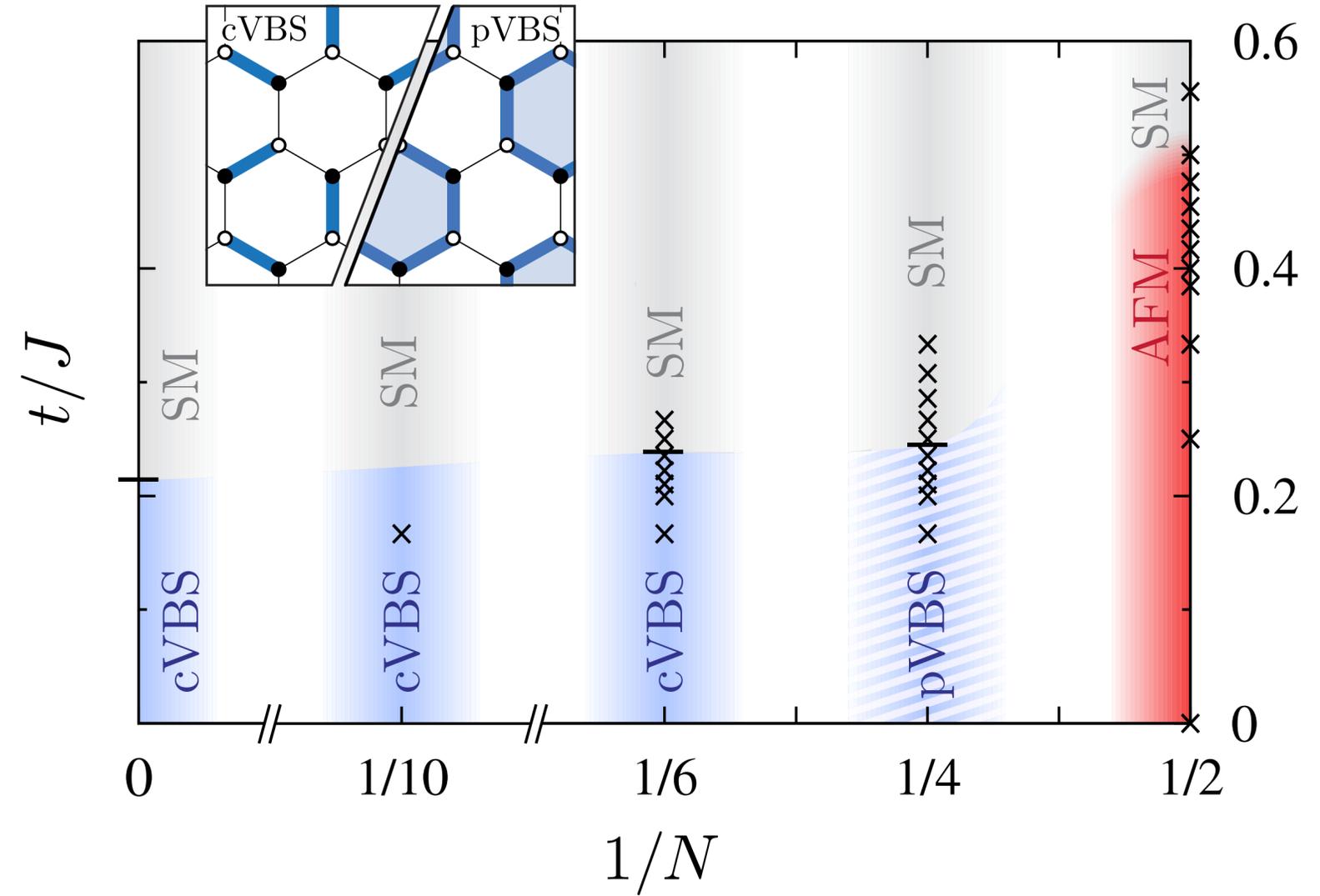
# Conclusions

SO( $N$ ) Majorana-Hubbard models



[LJ & Seifert, PRB '22]

SU( $N$ ) Hubbard-Heisenberg models



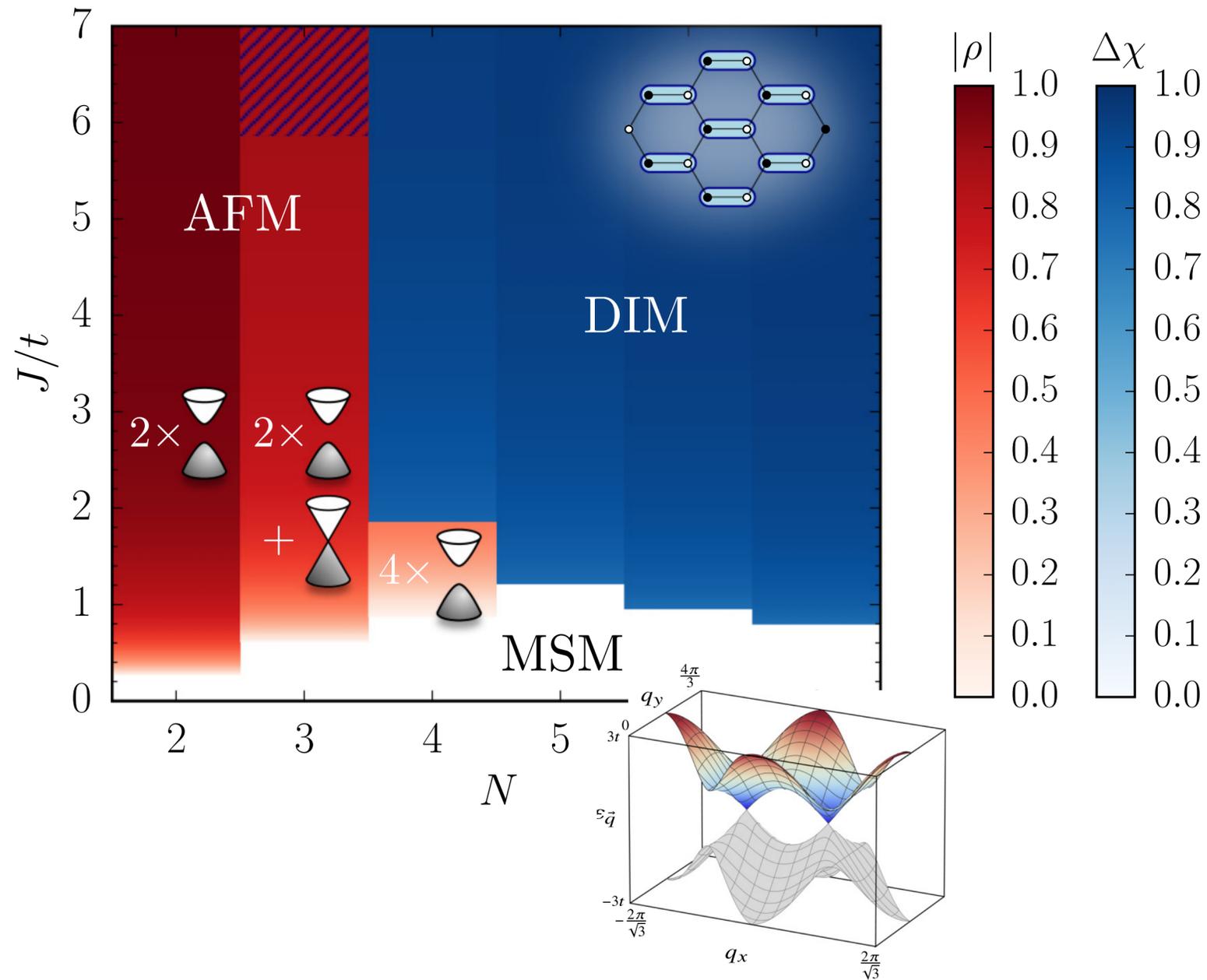
[Lang, Meng, Muramatsu, Wessel, Assaad, PRL '13]

[Affleck & Marston, PRB '88]

[Read & Sachdev, NPB '89]

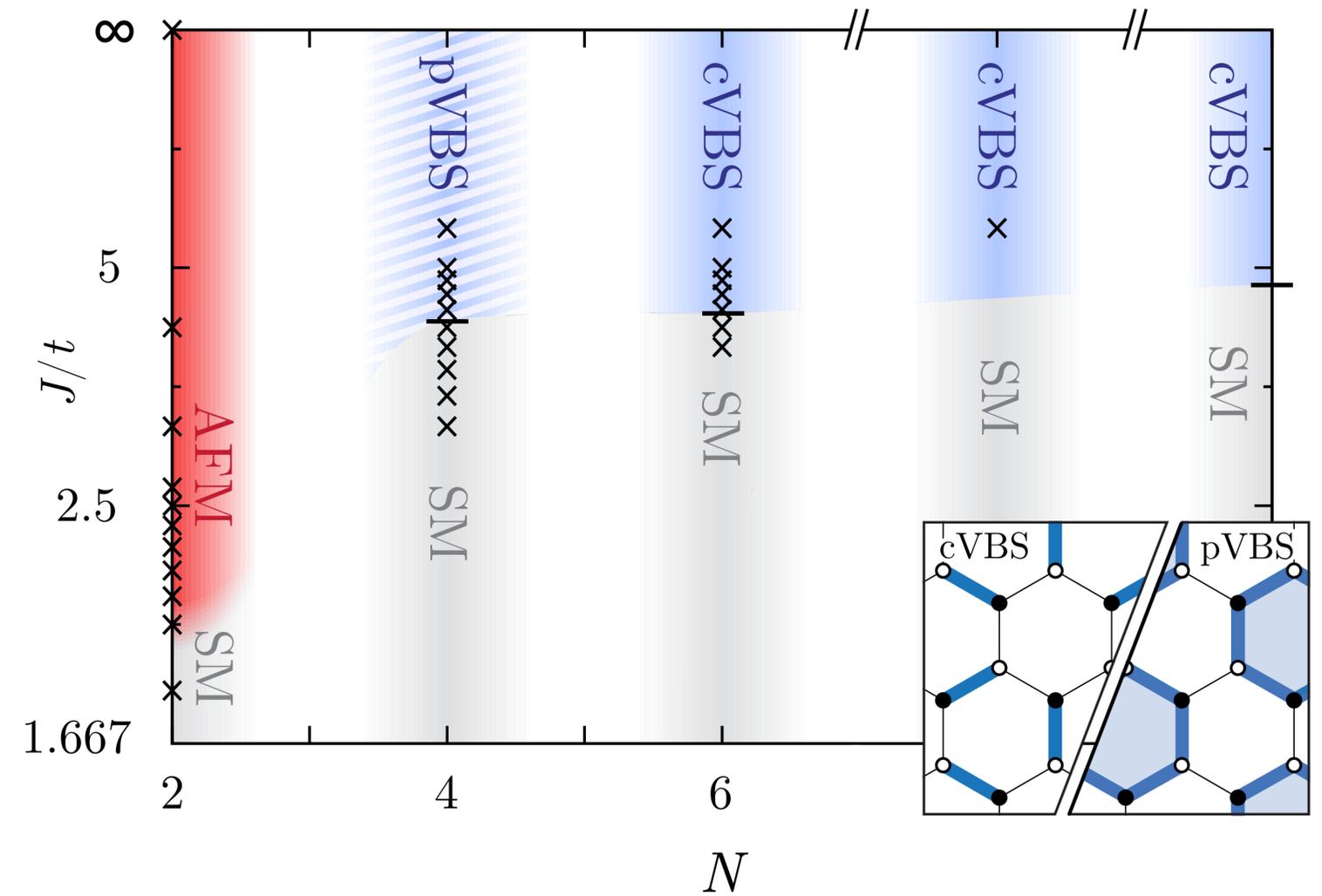
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